# An Unusual Two-Higgs Doublet Model from Warped Space

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## Outline The Elevator Talk

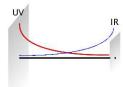
### Overview

- $\bullet$   $t\bar{t}$  condensation in RS
- Two Higgs Doublet Model
- **3** Move  $Q_{3L}$  away from IR brane
- Ease EWPO constraints.

## Randall Sundrum Models The Bare Essentials

RS Models are 5D theories with a non-trivial warped geometry.

$$\mathrm{d}s^2 = \mathrm{e}^{-k|y|} \mathrm{d}x^\mu \mathrm{d}x_\mu - \mathrm{d}y^2.$$



- The SM states are the zero modes of 5D fields.
- The fermions and KK modes are localised in the XD.

### Electroweak Constraints and $Z ightarrow bar{b}$

Low energy theory: integrate out KK modes.

 $\Longrightarrow$  Large effective operators for  $t_R$ ,  $Q_{3L}$ .

#### Relevant LEP constraints:

- T-parameter;
- $\bullet$   $Z \rightarrow b\bar{b}$ .

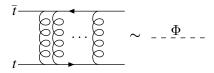
### Solve these problems by:

- Gauging  $SU(2)_R$ ;
- Adding a  $P_{LR}$  symmetry.

Can also ease constraints by moving  $Q_{3L}$  away from IR brane.

## Composite Higgses in RS

KK gluon mediates strong top-sector interactions.



Extra Higgs  $\implies$  extra contribution to  $m_{top}$   $\implies$  move  $Q_{3L}$  away from IR brane.

#### Earlier ideas:

- Using exotic fermions; or
- Using high (30 TeV) KK scale; and
- No fundamental Higgs.

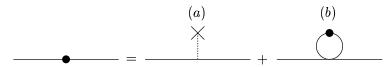


## The Gap Equation NJL Models

Integrate out KK gluon:



Generate mass term non-perturbatively:

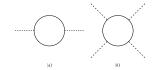


Rewrite NJL four-fermion term using auxillary scalar:

$$\frac{g}{M_{KK}^2} \left( \bar{\Psi} \Psi \right)^2 = M_{KK}^2 \Phi^2 + g \Phi \bar{\Psi} \Psi.$$

## Physics Below the KK Scale RGEs

Renormalisation of fields below KK gluon mass:



#### Generates:

- Kinetic term for Φ;
- ② Sizeable kinetic mixing between H,  $\Phi$ ;
- Mass mixing between H, Φ;
- All possible quartic terms.

Model quite predictive:  $M_{KK}$ ,  $m_0$ ,  $\mu$ ,  $\lambda_t$ ,  $g_t$ ,  $\lambda_0$ .

### Diagonalising the Kinetic Sector

#### Need to:

- Remove kinetic mixing;
- Bring kinetic terms to canonical normalisation.

Accomplish these goals with the field redefinition:

$$H = \hat{H}$$

$$\Phi = -\frac{\lambda_t}{g_t}\hat{H} + \frac{1}{g_t\sqrt{\epsilon}}\hat{\Phi}$$

Leads to very simple Lagrangian:

$$\begin{split} V(\hat{H},\hat{\Phi}) &= M_{hh}^2 \hat{H}^\dagger \hat{H} - M_{h\Phi}^2 \left( \hat{H}^\dagger \hat{\Phi} + \hat{\Phi}^\dagger \hat{H} \right) + M_{\Phi\Phi}^2 \hat{\Phi}^\dagger \hat{\Phi} \\ &\quad + \frac{1}{2} \lambda_0 (\hat{H}^\dagger \hat{H})^2 + \frac{1}{\epsilon} (\hat{\Phi}^\dagger \hat{\Phi})^2. \end{split}$$

## Matching to the Standard Model

Low Energy Boundary Conditions

In our model, both Higgses acquire a vev.

 $\Longrightarrow$  CP,  $U(1)_{em}$  are automatically conserved.

### Matching to SM:

- **1** Match EWSB:  $v_{ew}^2 = v_H^2 + v_{\Phi}^2$ .
- ② Match top quark mass:  $m_t = \frac{v\cos\beta}{\sqrt{2\epsilon}}$ .

This determines both vevs!



## A Qualitative Analysis The Decoupling Limit

Neglecting mass mixing, scalar potential has symmetry

$$SU(2)_{\Phi L} \times SU(2)_{\Phi R} \times SU(2)_{HL} \times SU(2)_{HL}$$
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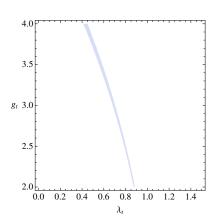
Vevs break this to  $SU(2)_{\Phi D} \times SU(2)_{HD}$ ; mass mixing to  $SU(2)_V$ .  $\Longrightarrow$  Implies degeneracy of  $H^{\pm}$ ,  $A^0$  ( $SU(2)_V$  triplet).

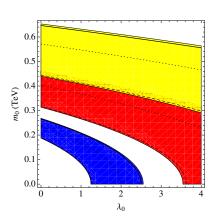
Model lives in decoupling limit:  $M_H^2 \sim M_{KK}^2 \gg v_{ew}$ .

⇒ Lighter scalar very SM-like;

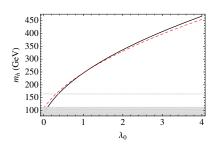
⇒ Small mixing among neutral scalars.

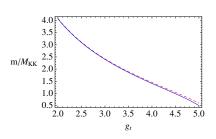
### Numerical Results I 2HDM Lagrangian Parameters



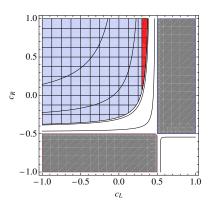


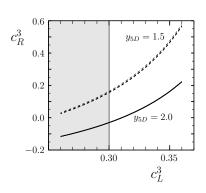
## Numerical Results II Spectrum





### Numerical Results III RS Lagrangian Parameters





## Flavour-Changing Neutral Currents Or, How I Learnt to Stop Worrying and Love My Model

Won't this model lead to tree-level FCNCs?

Yes.

Yukawa sector:

$$\mathcal{L}_{Y} = -\frac{\sqrt{2}\mathcal{M}_{ij}^{d}}{v\sin\beta} \, \overline{Q_{Li}} \, d_{jR} \, \hat{H} - \frac{\sqrt{2}\mathcal{M}_{ij}^{u}}{v\sin\beta} \, \overline{Q_{iL}} \, u_{jR} \, \widetilde{\hat{H}} + \frac{1}{\sqrt{\epsilon}} \, \overline{Q_{3L}} \, t_{R} \left( \widetilde{\hat{\Phi}} - \frac{\widetilde{\hat{H}}}{\tan\beta} \right) + h.c.$$

### Suppression factors:

- Light Higgs:  $\sin(\alpha \beta) \approx 10^{-3}$ ;
- Heavy Higgs:  $M_{KK}^{-1}$ ;
- Small Mixing Angles:  $\overline{Q_{3L}}$ ,  $t_R$  mostly top.

## Summary & Conclusions Time For Tea!

- Formation of a composite scalar doublet is quasi-generic in RS models.
- The resultant 2HDM eases some of the constraints.
- Flavour-changing top decays expected; other FCNCs probably small.