

# Non-Universal Gaugino Masses, Dark Matter and the LHC

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- Gaugino Sector of the MSSM (2 Slides)
- Dark Matter Hints (8 Slides)
- High-Energy Theoretical Motivation (2 Slides)
- LHC Phenomenology (4 Slides)

# Quick Review

- Gauginos part of vector supermultiplets:  $A_a = \{\lambda_a, (A_\mu)_a, D_a\}$ ,  $a = 1, 2, 3$

Names	spin 1/2	spin 1	$SU(3)_C, SU(2)_L, U(1)_Y$
gluino, gluon	$\tilde{g}$	$g$	$(\mathbf{8}, \mathbf{1}, 0)$
winos, W bosons	$\tilde{W}^\pm \tilde{W}^0$	$W^\pm W^0$	$(\mathbf{1}, \mathbf{3}, 0)$
bino, B boson	$\tilde{B}^0$	$B^0$	$(\mathbf{1}, \mathbf{1}, 0)$

- Supersymmetry breaking independent of EWSB  
Thus in SUSY limit we have massless gauginos up to EWSB effects

- Soft SUSY-breaking gaugino masses:  $\mathcal{L}_{\text{soft}} \ni -\frac{1}{2}M_a \lambda_a \lambda_a + \text{c.c.}$

- Gaugino masses run independently at one loop

$$\frac{dM_a}{dt} = \frac{1}{8\pi^2} b_a g_a^2 M_a, \quad b_a = -(3C_a - \sum_i C_a^i) \Rightarrow \{b_1, b_2, b_3\} = \left\{ \frac{33}{5}, 1, -3 \right\}$$

- Three ratios  $M_a/g_a^2$  therefore constant (up to two loop effects)

$$\frac{M_1}{g_1^2} \simeq \frac{M_2}{g_2^2} \simeq \frac{M_3}{g_3^2} \rightarrow M_3 : M_2 : M_1 \simeq 6 : 2 : 1 \quad \text{at EW scale}$$



# Gaugino Masses – EM Neutral Sector

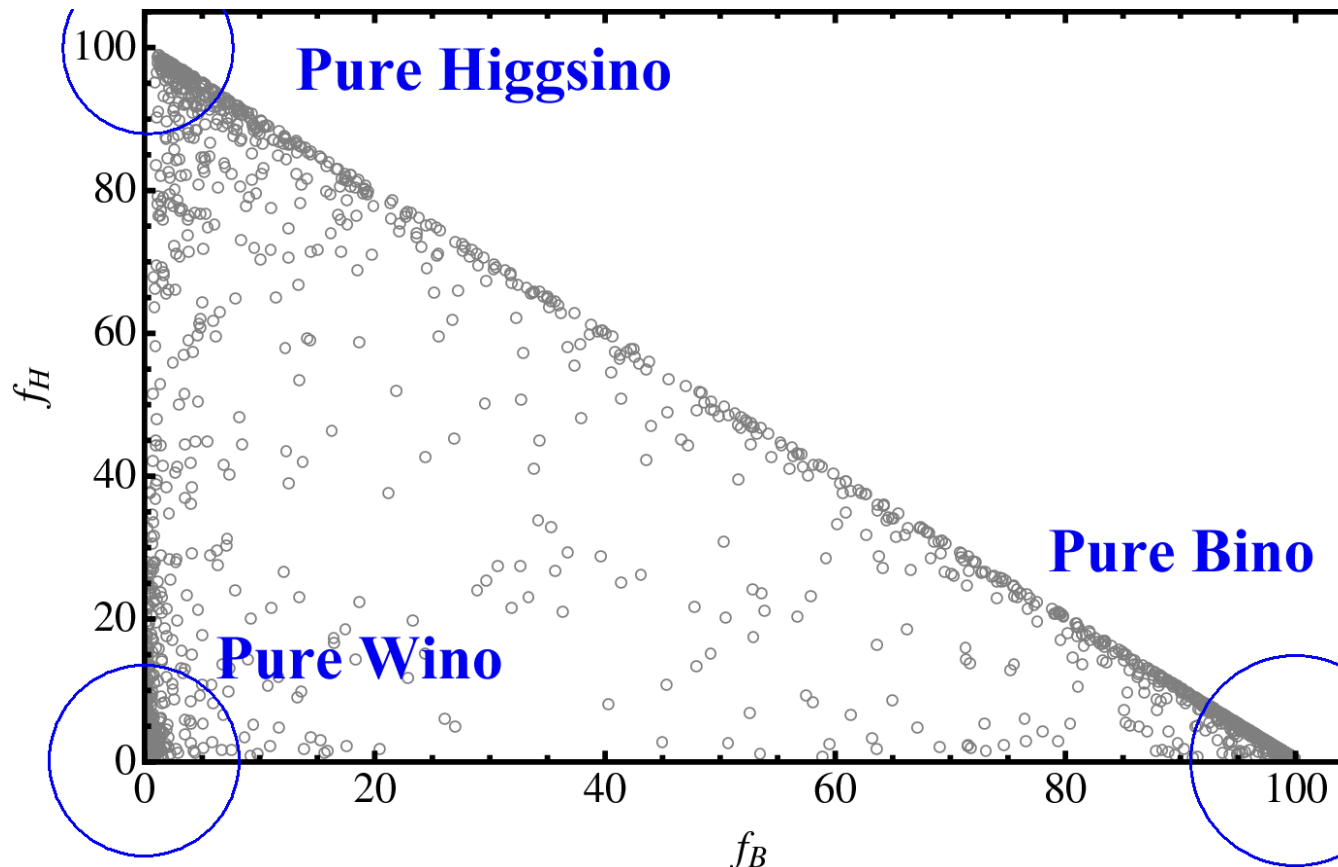
⇒ Can model possibilities via  $M_a = m_{1/2}(1 + \delta_a)$

Arkani-Hamed, Delgado, Giudice, NPB 741 (2006) 108

- $\delta_1 = \delta_2 = \delta_3$  produces bino-like LSP:  $\tilde{N}_1 \sim \tilde{B}$ ;  $\tilde{N}_2 \sim \tilde{W}^0$
- $\{\delta_1 = 0, \delta_2 < 0\}$  produces wino-like LSP
- $\{\delta_2 > 0, \delta_3 < 0\}$ ,  $|\delta_3| < |\delta_2|$  produces Higgsino-like LSP via RGEs + EWSB

$$M_Z^2 = 5.9M_3^2 - 1.8\mu^2 + 0.4m_0^2 - 0.4M_2^2 + \dots$$

Kane, Lykken, BDN, Wang, PLB 551 (2003) 146



# Dark Matter

# Dark Matter Signals – the Earliest SUSY Signature

- Assumption: lightest neutralino is stable LSP  $\Rightarrow$  dark matter

Goldberg, PRL 50 (1983) 1419

- Prediction: annihilation into photons, positrons, anti-protons, neutrinos

Silk & Srednicki, PRL 53 (1984) 624

- ★ Photons & neutrinos “point” back to source: high density areas such as galactic center or center of sun/earth
- ★ Charged particles must be propagated from origin to earth numerically
- ★ Both depend on the *halo profile*  $\rho_\chi(r)$  assumed for the dark matter candidate, but to varying degrees

- Begin with positrons:

Cirelli et al., NPB 800 (2008) 204; 813 (2009) 1

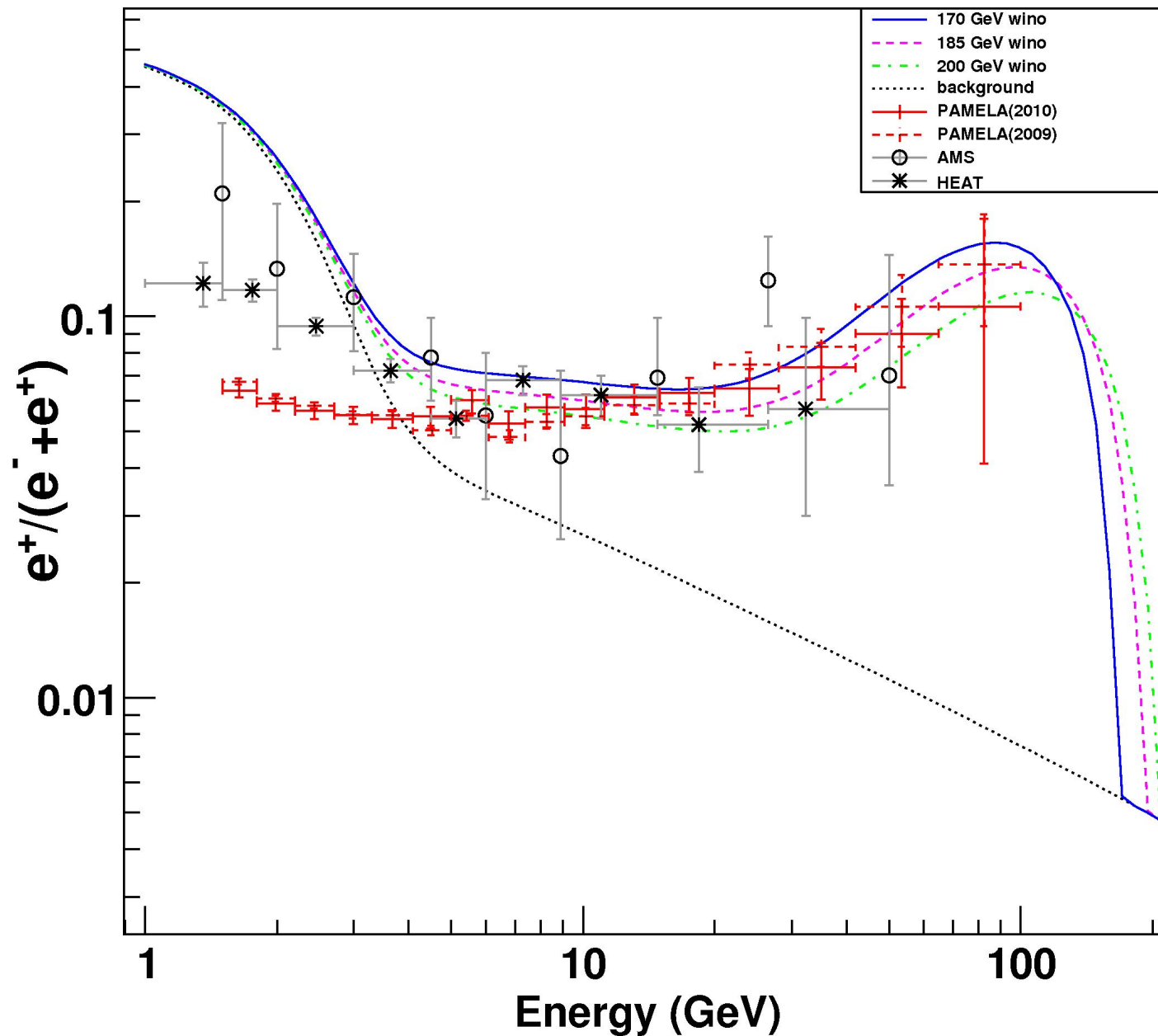
$$\Phi_{\bar{e}}(E) \simeq \frac{\tau_E B_{\bar{e}} c}{8\pi b(E)} \frac{\rho_\chi^2(r = R_0)}{m_{\tilde{N}_1}^2} F(E), \quad b(E) = 1 \text{ GeV} \left( \frac{E}{1 \text{ GeV}} \right)^2$$

$$F(E) = \int_E^{M_{\tilde{N}_1}} dE' \sum_k \langle \sigma v \rangle_{\text{halo}}^k \frac{dN_{\bar{e}}^k}{dE'} \cdot \mathcal{I}(E, E')$$

- ★  $B_{\bar{e}}$  = boost factor,  $\tau_E = \tau \times 10^{16}$  sec is the diffusion time scale and  $\mathcal{I}(E, E')$  is the halo function
- ★ For SUSY models, most important final state is usually  $k = W^+W^-$

# SUSY Fits to Positron Flux Measurements

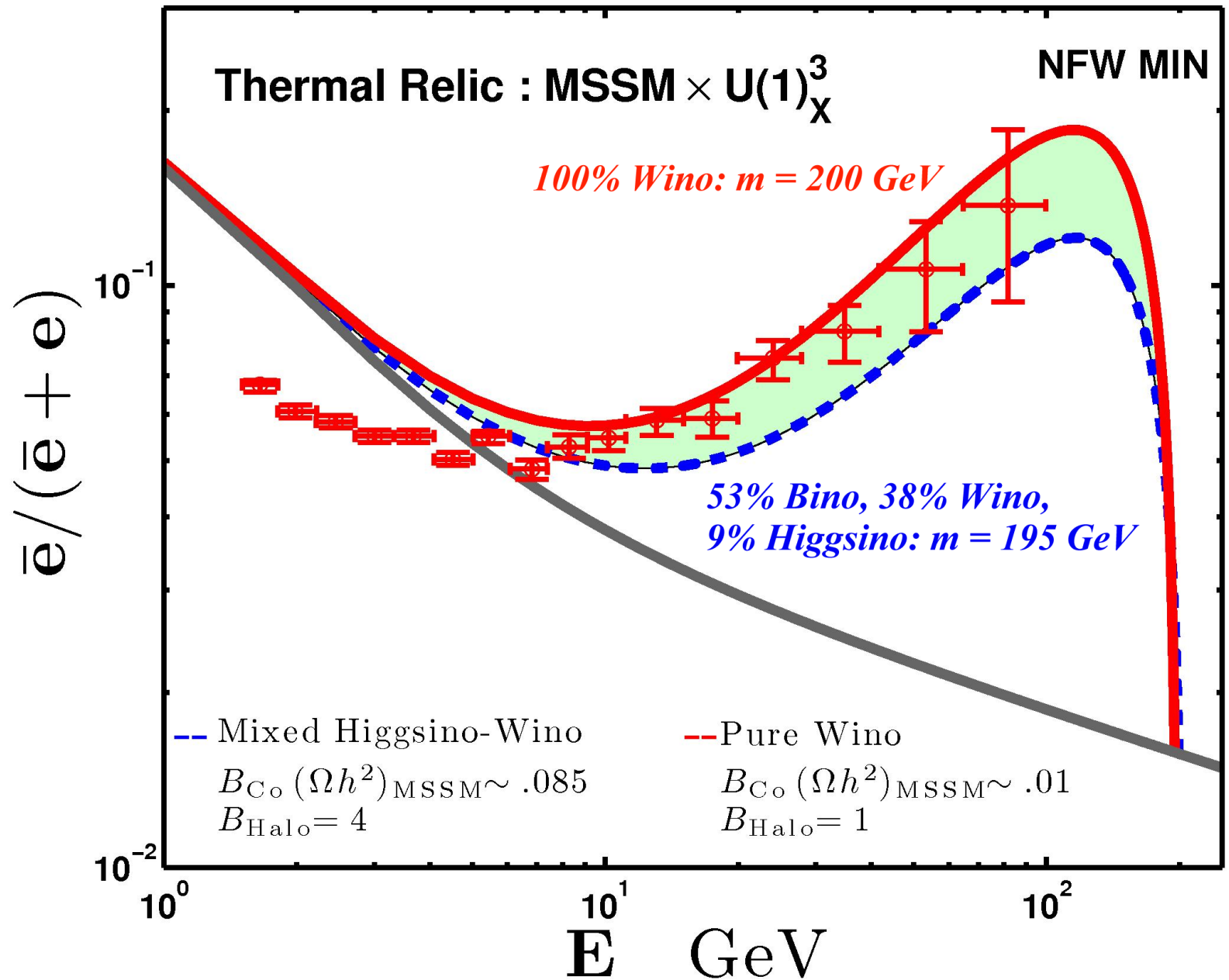
PAMELA Collaboration, arXiv:1001.3522



⇒ Best fits require  $\langle\sigma v\rangle_{WW} \simeq 2 \times 10^{-24} \text{ cm}^3/\text{s}$  and prefer NFW “min” profile

Feldman, Kane, Lu, BDN, arXiv:1002.2430

# SUSY Fits to Positron Flux Measurements

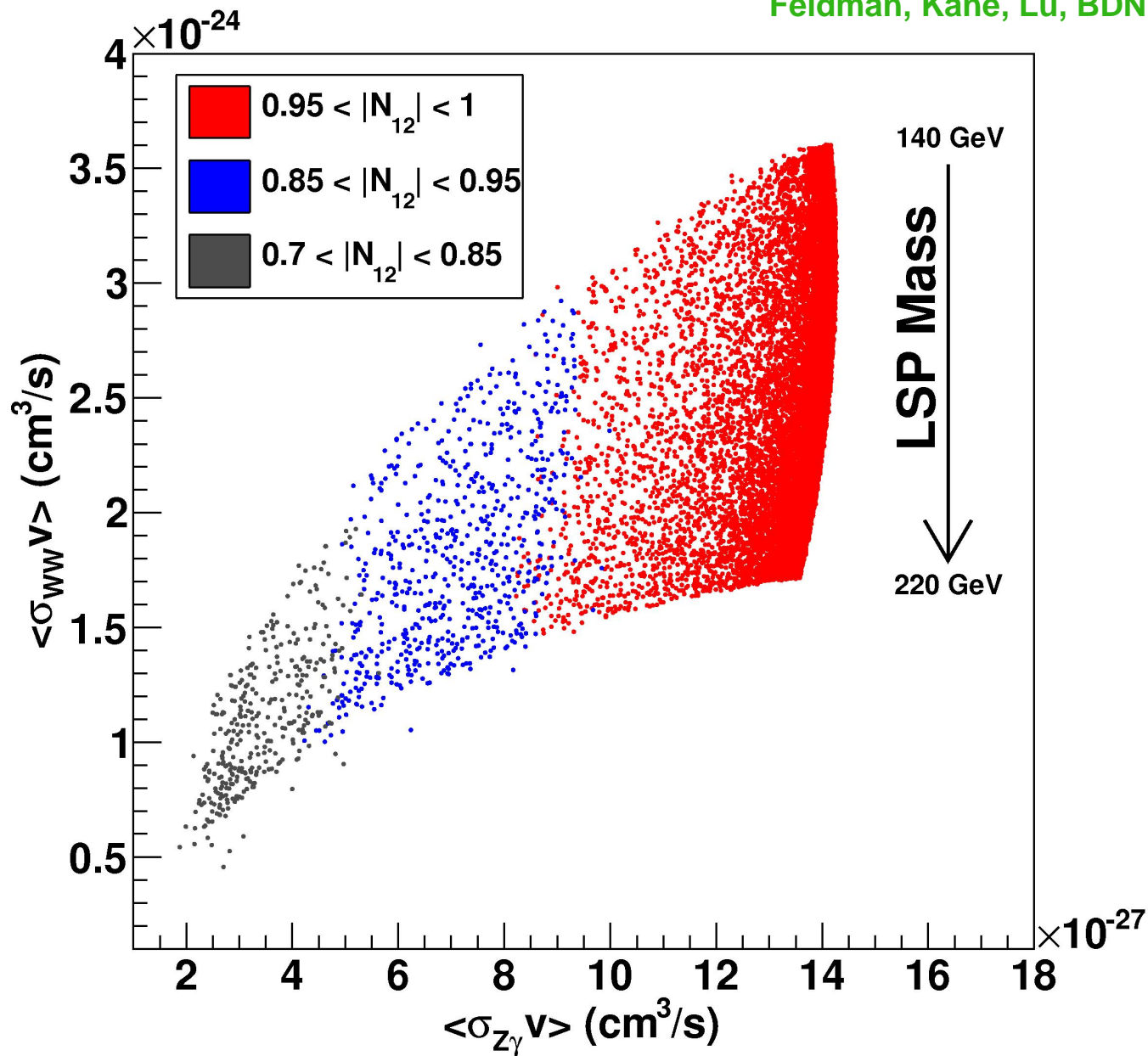


⇒ Pure wino not necessary – but must compensate with  $B_{\bar{e}}$  (here  $B_{HALO}$ )

Feldman, Liu, Nath, BDN, arXiv:0907.5392

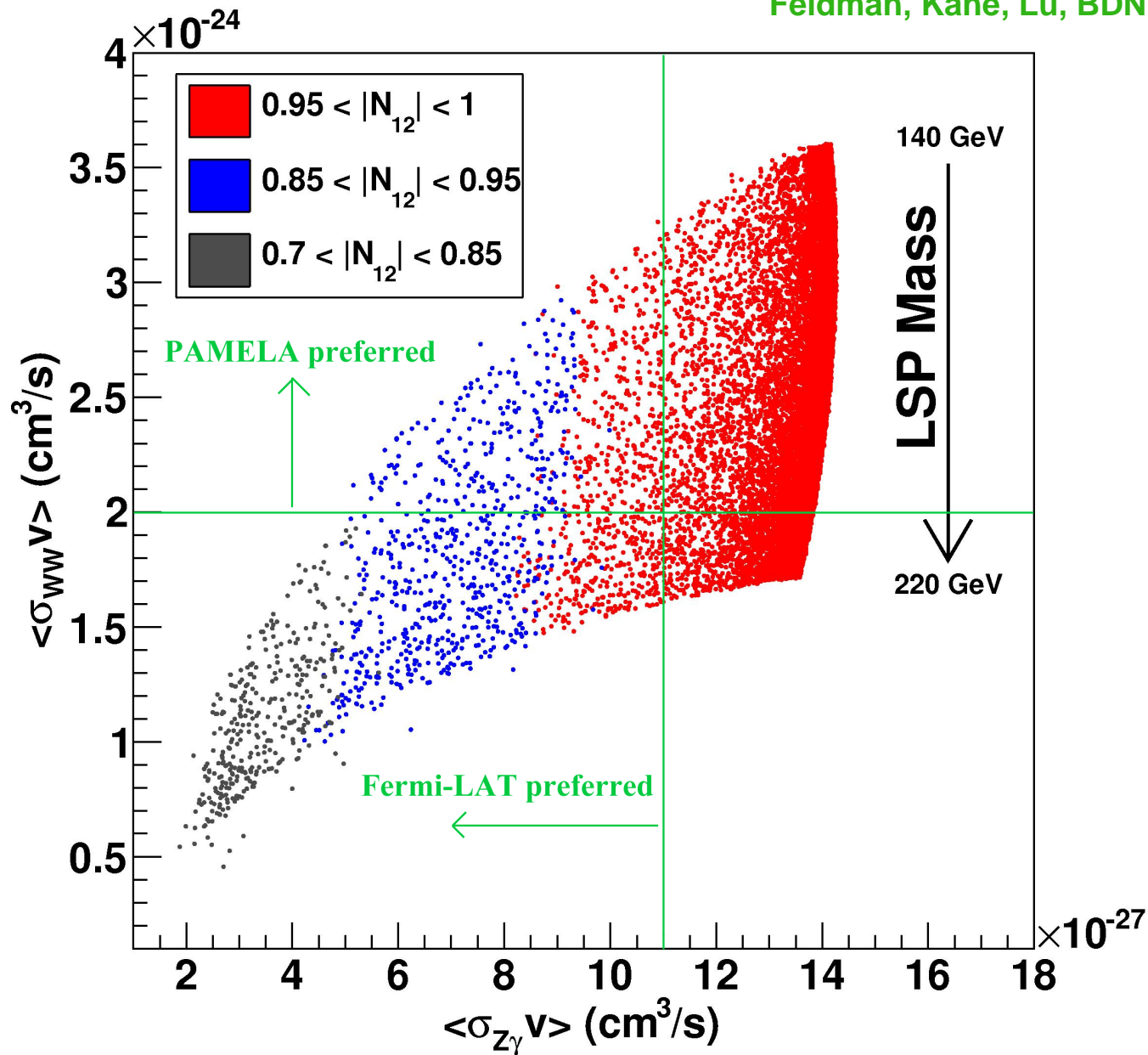
# Photons versus Positrons

Feldman, Kane, Lu, BDN, arXiv:1002.2430



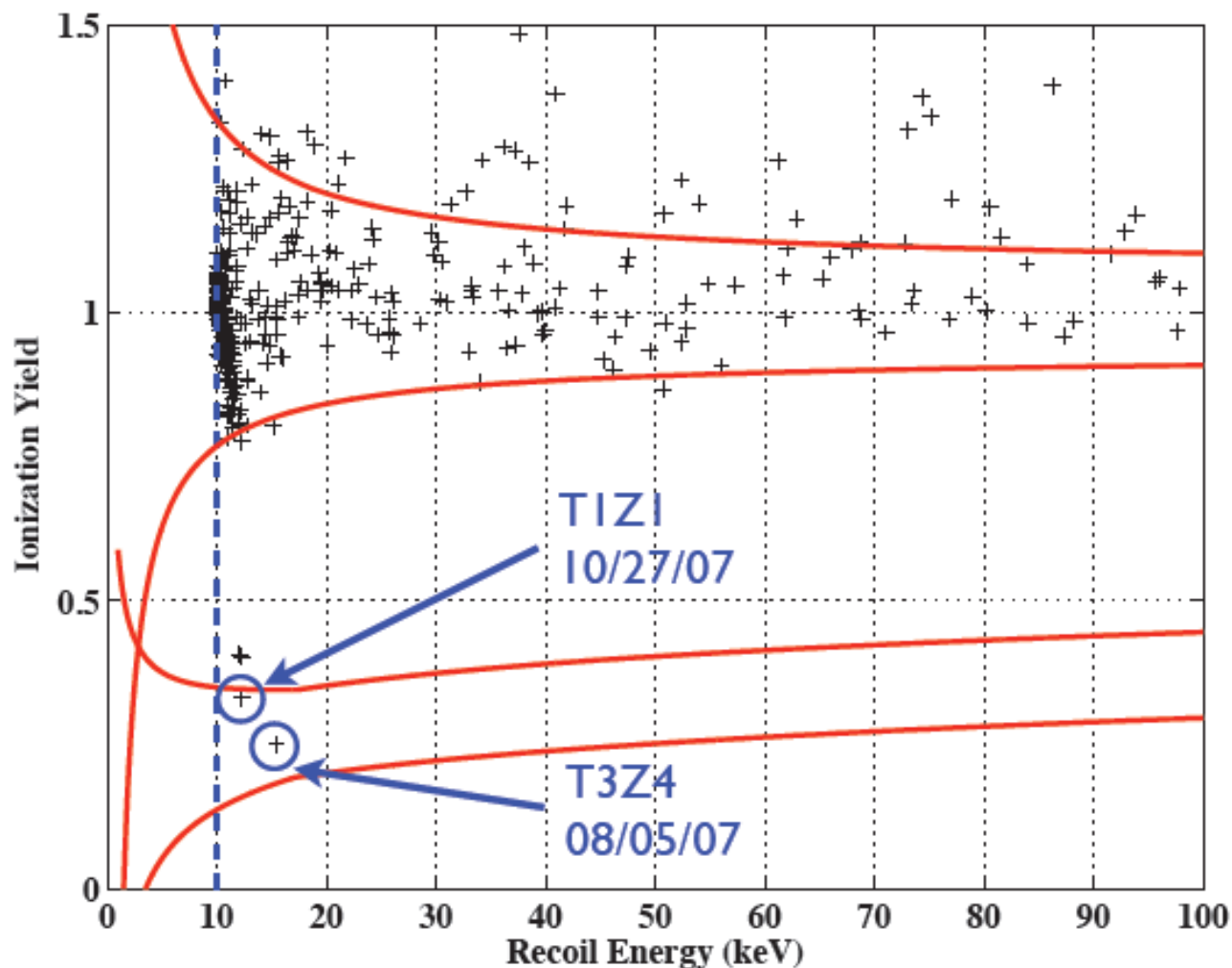
# Photons versus Positrons

Feldman, Kane, Lu, BDN, arXiv:1002.2430



# Direct Detection Experiments: CDMS II

⇒ December 2009 data release for 14 Ge detectors by CDMS-II Collaboration



CDMS II Collaboration, Science 327 (2010) 1620

- Two events in signal region with (revised) background estimate of  $0.8 \pm 0.1$  (stat)  $\pm 0.2$  (sys) events
- Implies an interaction cross-section  $\sigma_{\chi p}^{\text{SI}} \sim 10^{-44} \text{ cm}^2 = 1 \times 10^{-8} \text{ pb}$



# Fitting to CDMS II

⇒ Differential recoil rate at direct detection experiments given by

$$\frac{dR}{dE} = \sum_i c_i \frac{\rho_\chi \sigma_{\chi i}^{\text{SI}} |F_i(q_i)|^2}{2m_\chi \mu_{i\chi}^2} \int_{v_{\min}}^{\infty} \frac{f(\vec{v}, t)}{v} d^3v,$$

with  $F_i(q_i)$  being a nuclear form factor for  $i$ -th target nucleus

- Calculation of integrated event rate depends on experimental configuration

$$R = \int_{E_{\min}}^{E_{\max}} \frac{dR}{dE} dE; \quad (\text{Germanium}) : 10 \text{ keV} \leq E_{\text{recoil}} \leq 100 \text{ keV}$$

Altunkaynak, Holmes and BDN, arXiv:0804.2899

Point	A	B	C	D	E
$m_{\chi_1^0}$ (GeV)	138	190	175	112	230
$\delta_2$	0.65	0.62	-0.6	0.82	-0.47
$\delta_3$	-0.35	-0.3	-0.3	-0.35	-0.3
B%	3.0%	70.2%	0.3%	5.4%	40.9%
W%	0.4%	0.4%	95.8%	0.5%	53.0%
H%	96.6%	29.4%	3.9%	94.1%	6.1%
$\sigma_{\chi p}^{\text{SI}} \times 10^{45}$ (cm <sup>2</sup> )	11.9	44.4	41.3	35.3	74.8
$N_{\text{Ge}}$ (184 kg-days)	0.51	1.36	1.30	1.65	1.90

Holmes and BDN, arXiv:0912.4507

⇒ **Wino-like LSP preferred, but probably not 100% wino**

- Pure wino better for PAMELA (no boost factor) but tension with anti-protons and photons without help from halo model and/or diffusion parameters
- Higgsino or Bino component of 5-10% (at least) needed to avoid photon and anti-proton constraints – but need  $\mathcal{O}(5)$  boost factors to get PAMELA
- If CDMS-II is seeing a signal, will need even more substantial Higgsino component for large enough cross section

**All scenarios (probably) require non-thermal relic production mechanisms**

# High-Scale Theoretical Motivation

# What Can Cause Non-Universalities?

In supergravity, gaugino masses have a very simple form:

$$m_{\lambda_a} = \sum_n \frac{g_a^2}{2} \frac{F^n}{M_{\text{PL}}} \text{Re}[\partial_n f_a] ; \quad f_a = f_a(Z^n)$$

where  $f_a$  are gauge kinetic functions which depend on SM gauge singlets  $Z^n$

⇒ *So what are some mechanisms for producing non-universal gaugino masses?*

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1. Grand Unified Theories
2. Independent Gauge Kinetic Functions
3. Loop Effects

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3. Loop Effects

⇒ An example of the last item is the mirage pattern of gaugino masses

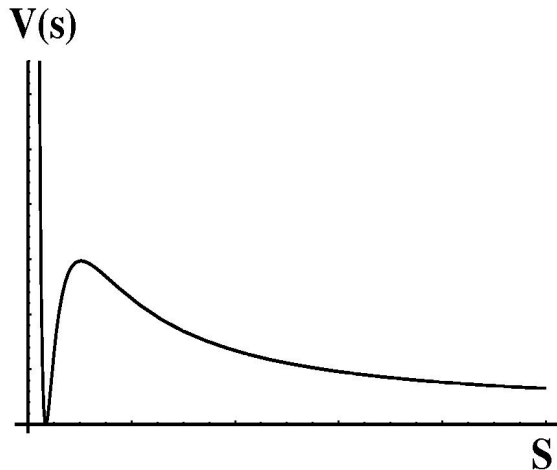
$$M_1 : M_2 : M_3 \simeq (1 + 0.66\alpha) : (2 + 0.2\alpha) : (6 - 1.8\alpha)$$

The mirage pattern (competition between tree and anomaly-mediated contributions to soft masses) appears in a number of phenomenologically successful string constructions:

- Kähler stabilized heterotic string models Binetruiy, Gaillard, Wu, **NPB 481** (1996) 109  
Gaillard and BDN, **IJMP A22** (2007) 1451
- Type-IIB flux compactifications with anti- $D_3$  branes Kachru, Kallosh, Linde, Trivedi, **PRD 68** (2003) 046005  
Choi, Falkowski, Nilles, Olechowski, **NPB 718** (2005) 113
- $M$ -theory compactified on fluxless  $G_2$  manifolds Acharya, Kane, et al.,  
**PRL 97**(2006) 191601  
**PRD 76** (2007) 126010  
**PRD 78** (2008) 065038

⇒ Common features:

- Single modulus stabilized by gaugino condensation
- Kähler potential for this modulus substantially altered from tree-level value
- Tuning of cosmological constant ( $\langle V \rangle$ ) to zero by adjusting parameters



# LHC Implications



- For each point studied 100,000 events generated with PYTHIA + PGS4 with the level 1 trigger only
  - ⇒ Typically this is about  $5 \text{ fb}^{-1}$  of signal
- A single SM sample was generated, including  $5 \text{ fb}^{-1}$  of top, bottom, dijets and gauge boson production (both single and double production)
  - ⇒ This background sample was suitably weighted to be included with each of our “signal” samples
- Initial object-level cuts to keep an object in the event record

Object	Minimum $p_T$	Minimum $ \eta $
Photon	20 GeV	2.0
Electron	20 GeV	2.0
Muon	20 GeV	2.0
Tau	20 GeV	2.4
Jet	50 GeV	3.0

⇒ After object-level cuts we impose event-level cuts – an example:

- $\cancel{E}_T > 150 \text{ GeV}$
- Transverse sphericity  $S_T > 0.1$
- $H_T = \cancel{E}_T + \sum_{\text{Jets}} p_T^{\text{jet}} > 600 \text{ GeV}$  (400 GeV for events with 2 or more leptons)

# Benchmark Models I: PAMELA Examples

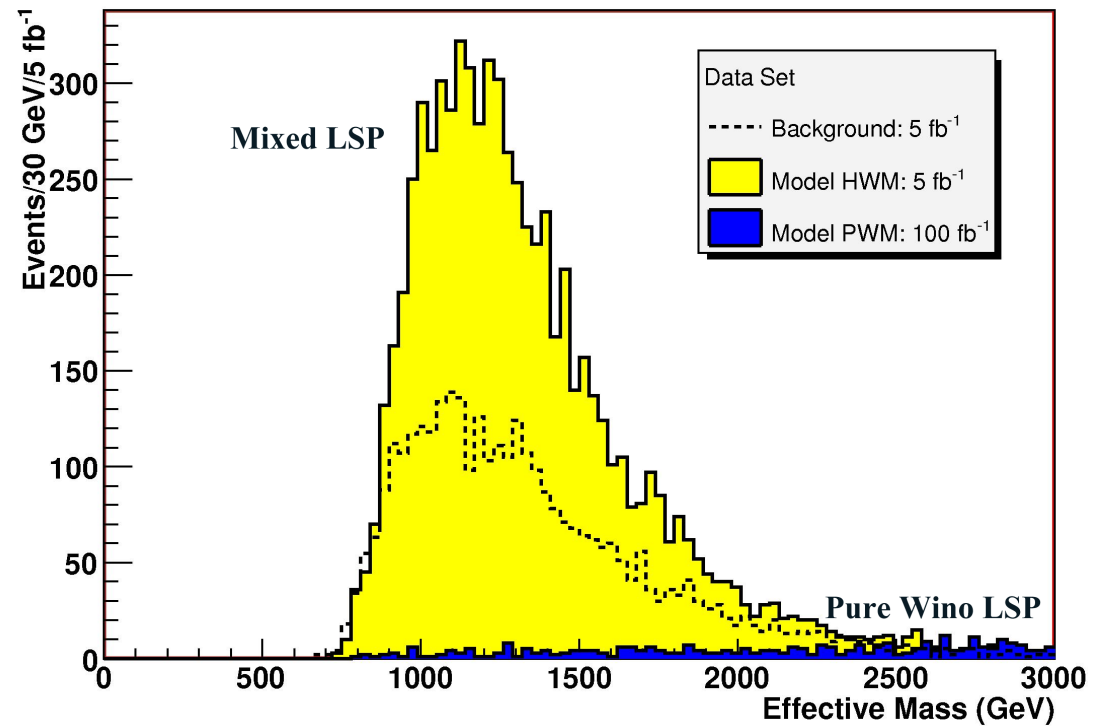
Feldman, Liu, Nath, BDN, reference

Mass	Mixed LSP	Pure Wino LSP
$m_{\tilde{N}_1}$	198.9	195.2
$m_{\tilde{N}_2}$	217.0	357.0
$m_{\tilde{N}_3}$	429.9	1025
$m_{\tilde{N}_4}$	451.3	1029
$m_{\tilde{C}_1}$	208.8	195.5
$m_{\tilde{C}_2}$	448.6	1036
$m_{\tilde{t}_1}$	648.5	1516
$m_{\tilde{t}_2}$	866.8	1749
$m_{\tilde{b}_1}$	841.4	1729
$m_{\tilde{b}_2}$	970.2	1902
$m_{\tilde{\tau}_1}$	817.7	1011
$m_{\tilde{\tau}_2}$	822.8	1041
$m_{\tilde{g}}$	707.1	1929

Signature	Mixed LSP		Pure Wino LSP	
	Events	$S/\sqrt{B}$	Events	$S/\sqrt{B}$
Multijets	8766	183.74	50	1.05
Lepton + jets	2450	32.25	26	0.34
OS dileptons + jets	110	6.39	4	0.23
SS dileptons + jets	60	11.77	0	NA
Trileptons + jets	14	2.47	0	NA

- Big impact of gluino mass in number of multijet events
- Small mass gaps significantly reduce number of leptonic events

$M_{\text{eff}}$  Distribution



10 fb<sup>-1</sup> at  $\sqrt{s} = 14$  TeV

# Benchmark Models II: CDMS-II Examples

Holmes and BDN, arXiv:0912.4507

Point	C	D	E
$\delta_2$	-0.6	0.82	-0.47
$\delta_3$	-0.3	-0.35	-0.3
B%	0.3%	5.4%	40.9%
W%	95.8%	0.5%	53.0%
H%	3.9%	94.1%	6.1%
$m_{\tilde{N}_1}$	175	112	230
$m_{\tilde{N}_2}$	235	130	239
$m_{\tilde{N}_3}$	505	252	504
$m_{\tilde{N}_4}$	513	846	515
$m_{\tilde{C}_1}$	175	123	234
$m_{\tilde{C}_2}$	514	846	515
$m_{\tilde{g}}$	952	890	951
$m_{\tilde{t}_1}$	719	544	709
$m_{\tilde{t}_2}$	862	964	865
$m_{\tilde{b}_1}$	809	766	812
$m_{\tilde{b}_2}$	874	943	871
$m_{\tilde{\tau}_1}$	344	338	352
$m_{\tilde{\tau}_2}$	414	752	424
$m_h$	113	114	113
$\sigma_{\text{SUSY}}^{7 \text{ TeV}}$ (pb)	1.2	2.7	0.4
$\sigma_{\text{SUSY}}^{10 \text{ TeV}}$ (pb)	2.5	5.1	1.3
$\sigma_{\text{SUSY}}^{14 \text{ TeV}}$ (pb)	5.7	10.0	3.7

- All models can produce signals at CDMS II – C & E can fit PAMELA data as well
- Signal simulated:  $1 \text{ fb}^{-1}$  at  $\sqrt{s} = 14 \text{ TeV}$
- Again, healthy multijets but disappearance of leptonic events

Numbers of Events

Point	C	D	E
Multijets	402	436	298
$1l + \text{jets}$	202	310	111
OS $2l + \text{jets}$	12	45	7
SS $2l + \text{jets}$	6	16	3
$3l + \text{jets}$	4	6	1

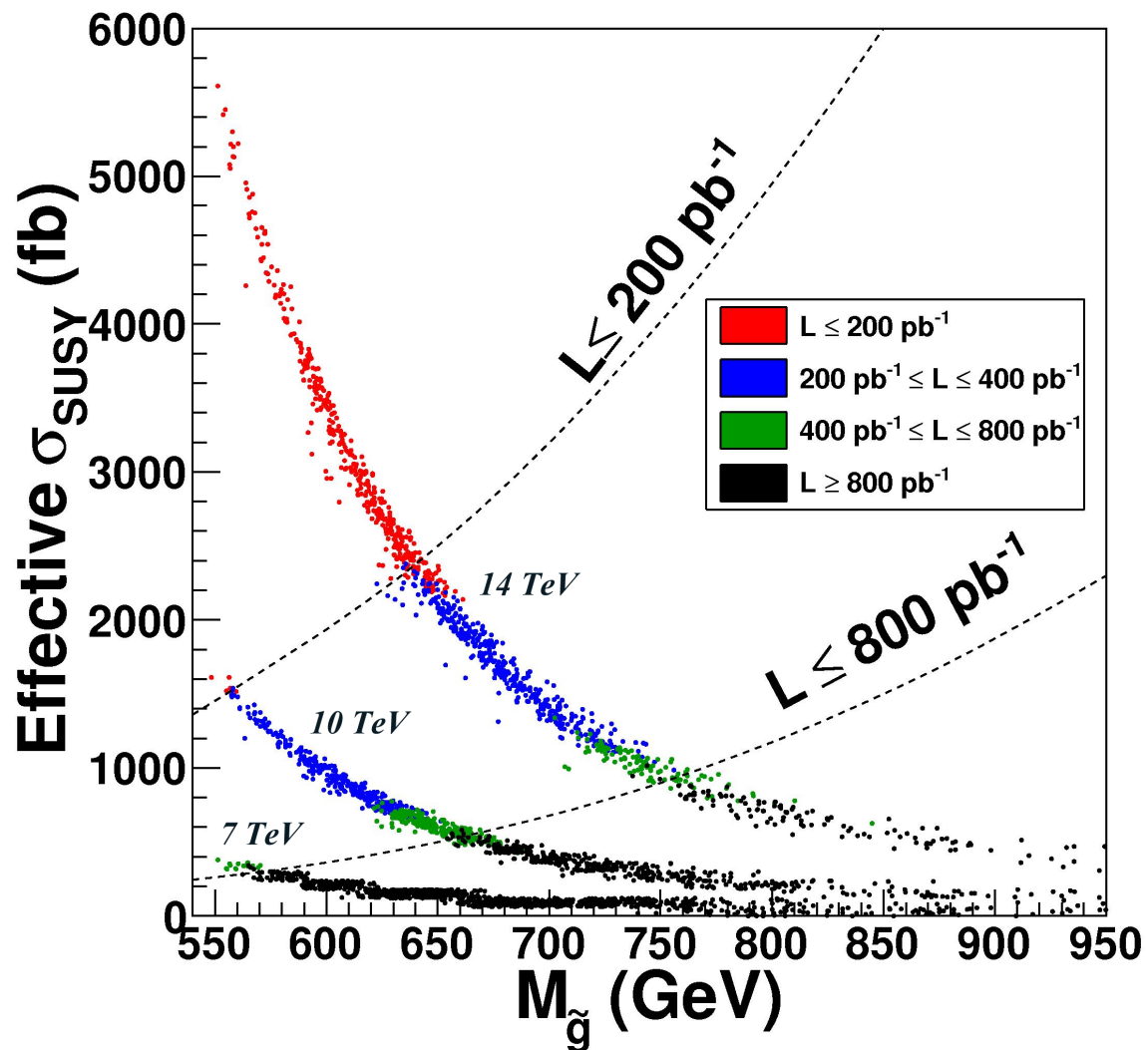
Significance  $S/\sqrt{B}$

Point	C	D	E
Multijets	26.9	29.1	19.9
$1l + \text{jets}$	8.2	12.5	4.5
OS $2l + \text{jets}$	2.0	7.4	1.2
SS $2l + \text{jets}$	2.3	6.0	1.1
$3l + \text{jets}$	1.6	2.5	0.4

# General PAMELA-consistent Models

⇒ General rule: Discovery of DM-motivated models needs a light gluino

Feldman, Kane, Lu, BDN, arXiv: 1002.2430



- High wino-content (for PAMELA) implies small mass gap between  $\tilde{C}_1/\tilde{N}_2$  and LSP
- Result: major reduction in expected leptonic SUSY signatures
- Increasing Higgsino content to match CDMS (and photon data) requires a light gluino
- Result: multijet signals may be our only handle

⇒ We will need to learn how to do more with less!

Must look for new signatures targeted to non-universalities in gaugino sector

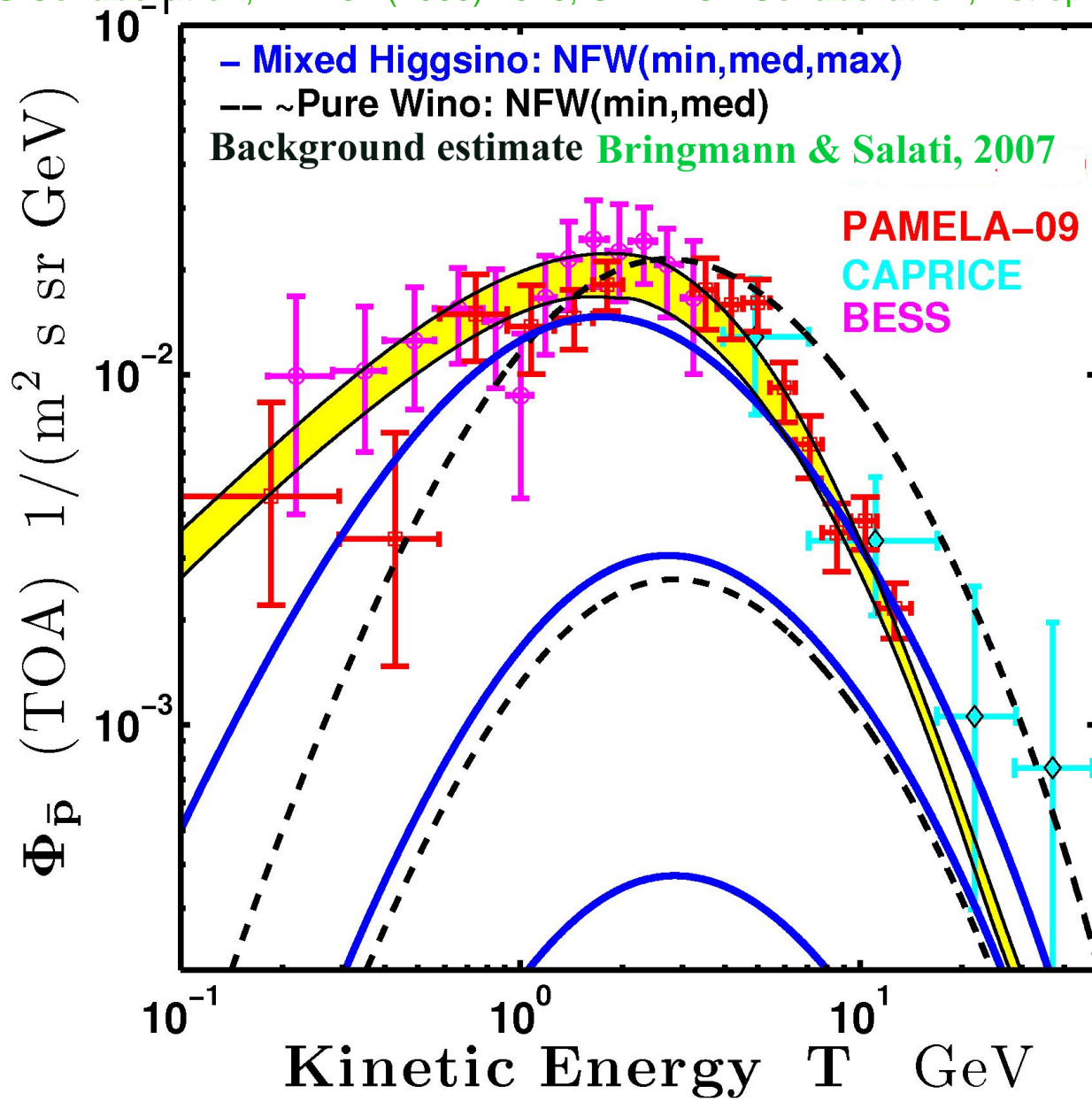
- Dark matter hints strongly disfavor pure Bino LSP (i.e. mSUGRA)
- PAMELA needs wino predominance; CDMS/photons want strong Higgsino admixture
- Such models find a natural home in many (all?) semi-realistic string constructions
- Likely that mass gaps between  $\tilde{C}_1/\tilde{N}_2$  and LSP small, so leptonic signatures a bust
- Will need to learn to do more with jet-based signatures and hope the gluino is lighter than in mSUGRA models

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  - Will need to learn to do more with jet-based signatures and hope the gluino is lighter than in mSUGRA models
- ⇒ *Gaugino sector is truly a window on the high-energy world: we may be on the verge of revolutionary discoveries!*

**Back-Up Slides**

# What About Anti-protons?

BESS Collaboration, PRL 84 (2000) 1078; CAPRICE Collaboration, Astrophys. J. 561 (2001) 787



⇒ Greater tension for pure wino LSP; OK for NFW “min” and “med” halo profiles



⇒ Halo profiles especially important in this situation

- Annihilation rates scale like the *square* of the density
- We observe the entire line-of-sight to the galactic center – therefore need to know the halo *profile*  $\rho_\chi(r)$
- Many possible profiles suggested in literature; each can be summarized by one parameter  $\bar{J}(\Delta\Omega)$

$$\bar{J}(\Delta\Omega) \equiv \frac{1}{\Delta\Omega} \int_{\Delta\Omega} d\Omega' J(\psi'); \quad J(\psi) = \frac{1}{8.5 \text{ kpc}} \int_{\text{l. o. s.}} ds(\psi) \left( \frac{\rho_\chi(r)}{0.3 \text{ GeV/cm}^3} \right)^2$$

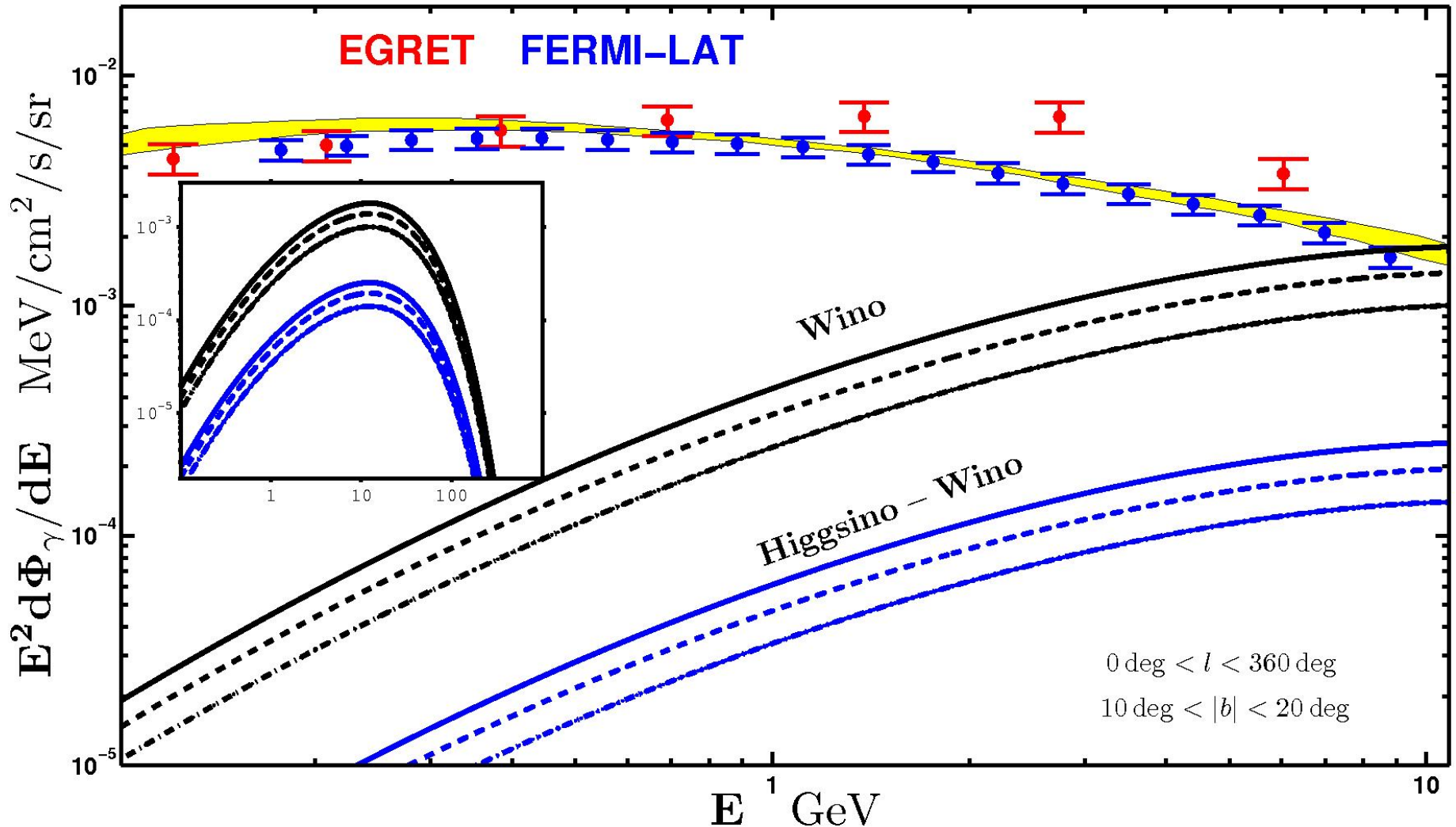
⇒ Two types of signal: continuous spectrum and mono-energetic lines

$$\frac{d\Phi_\gamma}{dE_\gamma} = 0.94 \times 10^{-13} \sum_i \frac{dN_\gamma^i}{dE_\gamma} \left( \frac{\langle \sigma_i v \rangle}{10^{-29} \text{ cm}^3 \text{ s}^{-1}} \right) \left( \frac{100 \text{ GeV}}{m_\chi} \right)^2 \bar{J}(\Delta\Omega) \Delta\Omega$$

- Typical sensitivities require  $\Phi_{\min} \sim 10^{-10}$  photons/cm<sup>2</sup>/sec
- Plain vanilla NFW profile gives  $\bar{J}(10^{-5} \text{ sr}) = 1.3 \times 10^4$
- Much less for isothermal core-type profiles

# Continuous Spectrum from Galactic Center

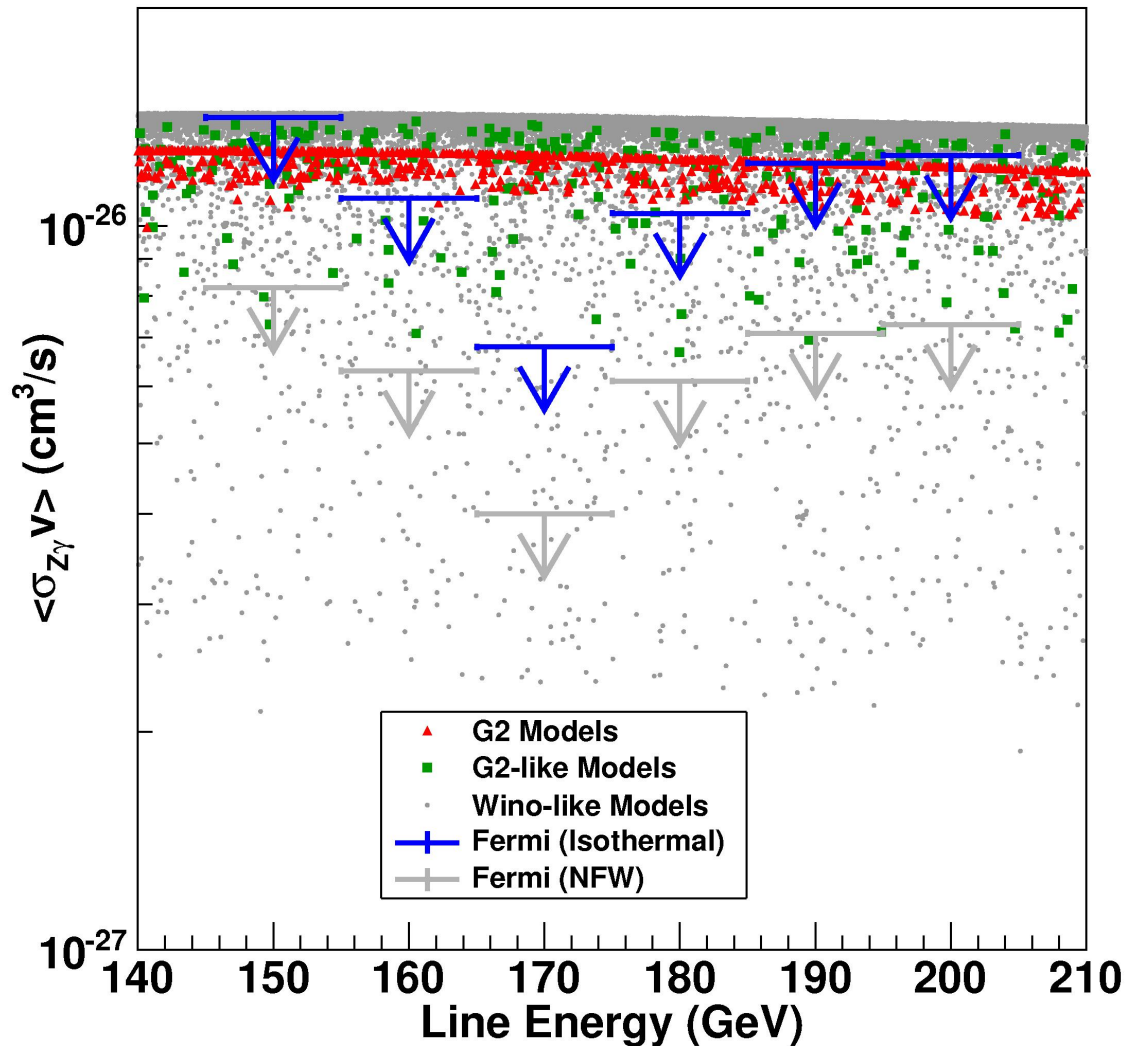
EGRET Collaboration, *Astrophys. J.* **481** (1997) 205  
Fermi-LAT Collaboration, arXiv: 0907.0294



- ⇒ Not much constraint on any profile from galactic center
- Profiles here: Einasto, NFW, isothermal (top to bottom)
- ⇒ More substantial constraints on pure-wino case coming from dwarf galaxies?

# Monochromatic Signals

⇒ Monochromatic gamma ray signals a “smoking gun” for dark matter



- Loop-induced diagrams provide annihilation into  $\gamma\gamma$  and  $\gamma Z$  final states
- Monoenergetic signals with  $E_{\gamma\gamma} = m_\chi$  and  $E_{\gamma Z} = m_\chi - M_Z^2/4m_\chi$
- Easy to pick out over background, but branching fractions reduce rate by factors of  $10^3 - 10^4$
- Pure-wino models capable of getting PAMELA correct in trouble!

Fermi-LAT Collaboration, arXiv: 1001.4531

## 3. Loop Effects

- Gauge coupling automatic when single modulus controls all gauge couplings
- Example: heterotic string models with  $f_a = S$  (gauge coupling relation...)
- Non-universalities now arise only at the loop level

$$\mathcal{L} \sim \int d^2\theta f_a (W^\alpha W_\alpha)_a \rightarrow \int d^2\theta \left( S + \frac{1}{16\pi^2} X_a \right) (W^\alpha W_\alpha)_a$$

- If  $\langle F^X \rangle \sim 16\pi^2 \langle F^S \rangle$  non-universalities are  $\mathcal{O}(1)$  in gaugino sector

# Testing for the Mirage Pattern

- ⇒ Our goal is to ask how well we can determine  $\alpha$  at the LHC using only **actual observations**
- Most importantly, can we demonstrate  $\alpha \neq 0$ ?
  - Want to do this independent of any particular model
  - Not going to assume reconstruction any sparticle masses
- ⇒ Basic idea: use an ensemble of signatures wisely chosen to perform a fit of Monte Carlo to “data”
- We break the problem into a “base model” specified by the parameters

$$\left\{ \begin{array}{c} \tan \beta, m_{H_u}^2, m_{H_d}^2 \\ M_3, A_t, A_b, A_\tau \\ m_{Q_{1,2}}, m_{U_{1,2}}, m_{D_{1,2}}, m_{L_{1,2}}, m_{E_{1,2}} \\ m_{Q_3}, m_{U_3}, m_{D_3}, m_{L_3}, m_{E_3} \end{array} \right\}$$

and a value of  $\alpha$  which determines the three gaugino masses  
(with overall scale set by  $M_3$ )

- ⇒ Choose a random “base model” and construct “alpha-line” based off this point
- Each line:  $-0.5 \leq \alpha \leq 1.0$  for the parameter  $\alpha$  in steps of  $\Delta\alpha = 0.05$

# Signature List C

- Here we allow as much as 30% correlation between any two signatures

	Description	Min Value	Max Value
Counting Signatures			
1	$N_\ell$ [ $\geq 1$ leptons, $\leq 4$ jets]		
2	$N_{\ell+\ell^-}$ [ $M_{\text{inv}}^{\ell^+\ell^-} = M_Z \pm 5$ GeV]		
3	$N_B$ [ $\geq 2$ B-jets]		
[0 leptons, $\leq 4$ jets]			
4	$M_{\text{eff}}^{\text{any}}$	1000 GeV	End
5	$M_{\text{inv}}^{\text{jets}}$	750 GeV	End
6	$\cancel{E}_T$	500 GeV	End
[0 leptons, $\geq 5$ jets]			
7	$M_{\text{eff}}^{\text{any}}$	1250 GeV	3500 GeV
8	$r_{\text{jet}}$ [3 jets $> 200$ GeV]	0.25	1.0
9	$p_T$ (4th Hardest Jet)	125 GeV	End
10	$\cancel{E}_T/M_{\text{eff}}^{\text{any}}$	0.0	0.25
[ $\geq 1$ leptons, $\geq 5$ jets]			
11	$\cancel{E}_T/M_{\text{eff}}^{\text{any}}$	0.0	0.25
12	$p_T$ (Hardest Lepton)	150 GeV	End
13	$p_T$ (4th Hardest Jet)	125 GeV	End
14	$\cancel{E}_T + M_{\text{eff}}^{\text{jets}}$	1250 GeV	End

## Signature "List" C

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3	$N_B$ [ $\geq 2$ B-jets]		
[0 leptons, $\leq 4$ jets]			
4	$M_{\text{eff}}^{\text{any}}$	1000 GeV	End
5	$M_{\text{inv}}^{\text{jets}}$	750 GeV	End
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[ $\geq 1$ leptons, $\geq 5$ jets]			
11	$\cancel{E}_T/M_{\text{eff}}^{\text{any}}$	0.0	0.25
12	$p_T$ (Hardest Lepton)	150 GeV	End
13	$p_T$ (4th Hardest Jet)	125 GeV	End
14	$\cancel{E}_T + M_{\text{eff}}^{\text{jets}}$	1250 GeV	End

## Signature “List” C

- Some signatures designed to detect changes in the softness of decay produces in cascade decays

- Particularly effective is the ratio  $r_{\text{jet}} \equiv \frac{p_T^{\text{jet3}} + p_T^{\text{jet4}}}{p_T^{\text{jet1}} + p_T^{\text{jet2}}}$