

Unitarity and Dark Matter in the Private Higgs Model

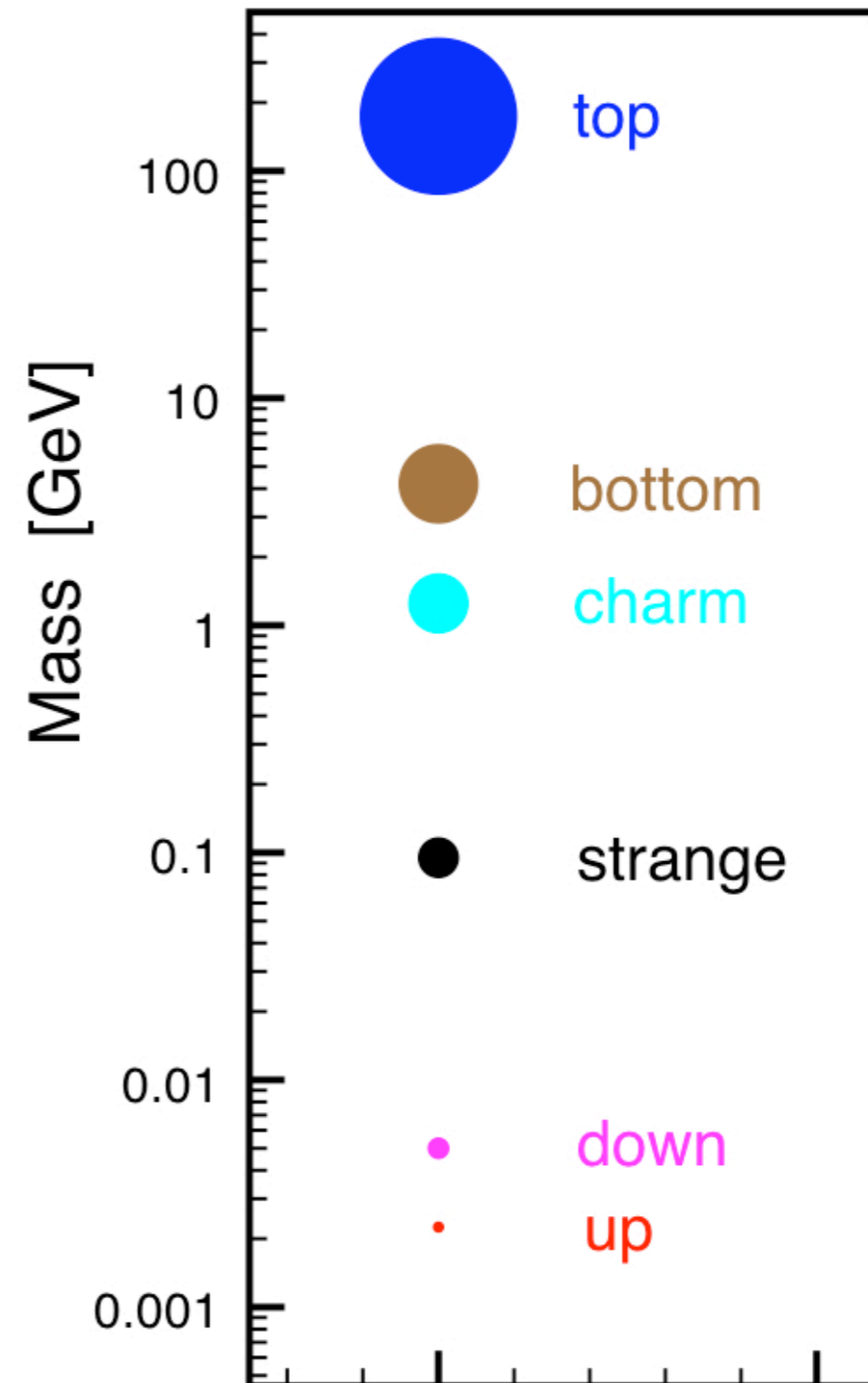
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- **Motivation**
- **Survey of the Model**
- **Perturbative Unitarity and the PH model**
- **Private Higgs Dark Matter (PHDM)**
- **Indirect Detection of PHDM**



Motivation

- “Gi-normous” hierarchy in the fermion mass spectrum
- Quark sector: Top quark is 10^5 times heavier than up quark! (top special?)
- Only one Higgs doublet (e.g., SM):
mass hierarchy \Leftrightarrow Yukawa hierarchy
- Possible explanations?
 - SUSY: Yukawa unification in $S0(10)$ GUTs
 - ExDim: “location, location, location, ...”
- Something simpler?



The Private Higgs (PH) Model

(Porto and Zee, arXiv:0712.0448)

- The Main Idea:
 - Introduce one “private” Higgs doublet per quark (“Democratic Higgs”?)
 - Construct SSB pattern s.t. *all* Yukawa couplings of order one
- PH doublets (φ_q) have same $SU(2) \times U(1)$ quantum numbers as SM Higgs
- Also, introduce a gauge singlet scalar S (for reasons given below)
- Impose set of six separate discrete symmetries K_q where:

$$U_R \rightarrow -U_R \quad (D_R \rightarrow -D_R) \quad , \quad \varphi_q \rightarrow -\varphi_q \quad , \quad S \rightarrow -S$$

- Lagrangian:

$$\mathcal{L} = \mathcal{L}_{SM-H} - \sum_q (y_D^{PH} \bar{Q}_L \phi_D D_R + y_U^{PH} \bar{Q}_L \tilde{\phi}_U U_R) \\ + \partial_\mu S \partial^\mu S + \sum_q [(D_\mu \phi_q)^\dagger D^\mu \phi_q - V(S, \phi_q)] ,$$

The PH Potential and EWSB

$$V(S, \phi_q) = \frac{\lambda_S}{4} \left(S^2 - \frac{v_d^2}{2} \right)^2 + \sum_q \left(\frac{1}{2} M_{\phi_q}^2 \phi_q^\dagger \phi_q + \lambda_q (\phi_q^\dagger \phi_q)^2 - g_{sq} S^2 \phi_q^\dagger \phi_q \right) \\ - \sum_{q \neq q'} \left(\frac{\gamma_{qq'}}{\sqrt{2}} v_s S \phi_q^\dagger \phi_{q'} + a_{qq'} \phi_q^\dagger \phi_{q'} \phi_q^\dagger \phi_{q'} + b_{qq'} \phi_q^\dagger \phi_q \phi_{q'}^\dagger \phi_{q'} + c_{qq'} \phi_q^\dagger \phi_{q'} \phi_{q'}^\dagger \phi_q \right) + h.c.$$

- Top PH plays the role of the SM Higgs (i.e., responsible for m_W and m_Z)
- Use vev of S and g_{st} coupling to drive EWSB:

$$\frac{1}{2} M_{\phi_t}^2 - g_{st} v_s^2 \equiv \mu_t^2 < 0$$

- Non-top PH fields acquire vev's in slightly different manner
- Use vev's of S and top PH... along with cubic term $\gamma_{qq'}$:

$$\frac{1}{2} M_{\phi_q}^2 \phi_q^\dagger \phi_q - \frac{\gamma_{qt}}{\sqrt{2}} \frac{v_h v_s^2}{2} \phi_q$$



$$\langle \phi_q \rangle = \frac{\gamma_{qt}}{\sqrt{2}} \frac{v_h v_s^2}{M_{\phi_q}^2} \equiv v_q$$

PH-enomenology

- Lighter the quark... heavier its PH partner (“up” PH $\sim 10^2 - 10^3$ TeV)
- To interest of LHC:
 - Two light scalars: h^0 and K^0
 - Heavy Scalar H^0 , charged scalar H^\pm and pseudoscalar A_b
- H^0 , H^\pm and pseudoscalar A_b all have masses $\sim 1-2$ TeV and $O(1)$ Yukawa couplings to bottom quarks
- For small mixing between top PH and S (angle = β):
 - h^0 has same properties as SM Higgs
 - K^0 provides a good candidate for Dark Matter (DM)!



Perturbative Unitarity in the PH Model

- Requiring perturbative unitarity provides important constraints (e.g., in the SM, limits on m_h and/or scattering energy... see Lee, Quigg and Thacker '77)
- In models with extended scalar sectors, $hh \rightarrow hh$ probes self-interactions
- Analysis performed in terms of “partial waves”, e.g. the J=0 partial wave:

$$a_0 = \frac{1}{16\pi s\beta^2} \int_{-s\beta^2}^0 \mathcal{A}(h^0 h^0 \rightarrow h^0 h^0) dt$$

- In the PH model:



$$a_0 = -\frac{1}{32\pi} \left\{ g_{h^0 h^0 h^0 h^0} + \frac{g_{h^0 h^0 h^0}^2}{s - m_{h^0}^2} + \frac{g_{h^0 h^0 K^0}^2}{s - m_{K^0}^2} + \frac{g_{h^0 h^0 H^0}^2}{s - m_{H^0}^2} - \frac{2}{s\beta^2} \left[g_{h^0 h^0 h^0}^2 \ln \left(\frac{s}{m_{h^0}^2} - 3 \right) \right. \right. \\ \left. \left. + g_{h^0 h^0 K^0}^2 \ln \left(\frac{s - 4m_{h^0}^2 + m_{K^0}^2}{m_{K^0}^2} \right) + g_{h^0 h^0 H^0}^2 \ln \left(\frac{s - 4m_{h^0}^2 + m_{H^0}^2}{m_{H^0}^2} \right) \right] \right\}.$$

Bounds on SM-like Higgs Mass

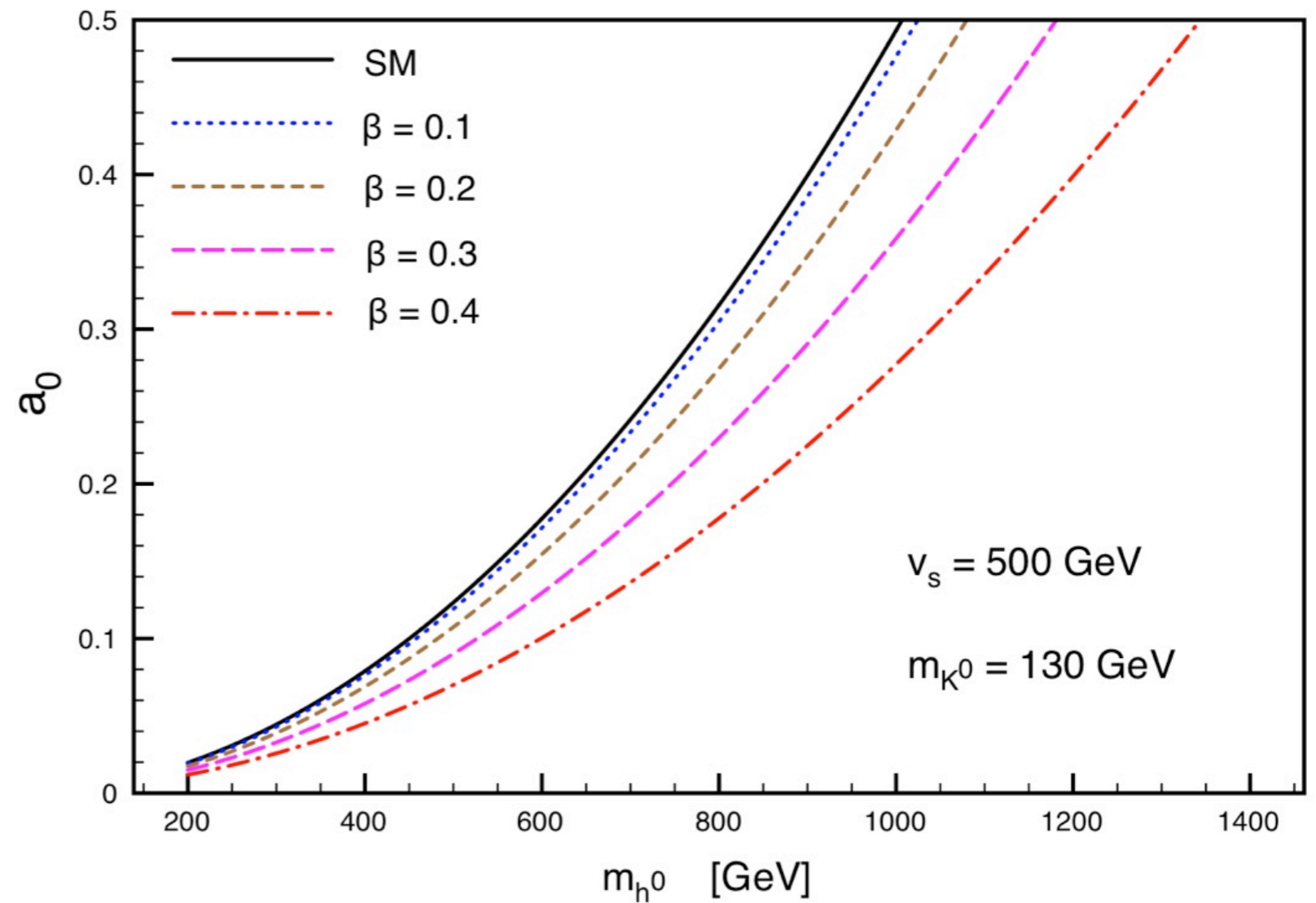
- To extract maximum Higgs mass consistent with PU, take $s \gg m^2$ limit:

$$\lim_{s \rightarrow \infty} a_0 = -\frac{1}{32\pi} g_{h^0 h^0 h^0 h^0} .$$

- SM Limit ($\beta \rightarrow 0$):

$$m_h \lesssim 1 \text{ TeV}$$

- PH = *softens* bounds
- Similar results for $W_L W_L, Z_L Z_L$, etc.
- Results nearly independent of m_K

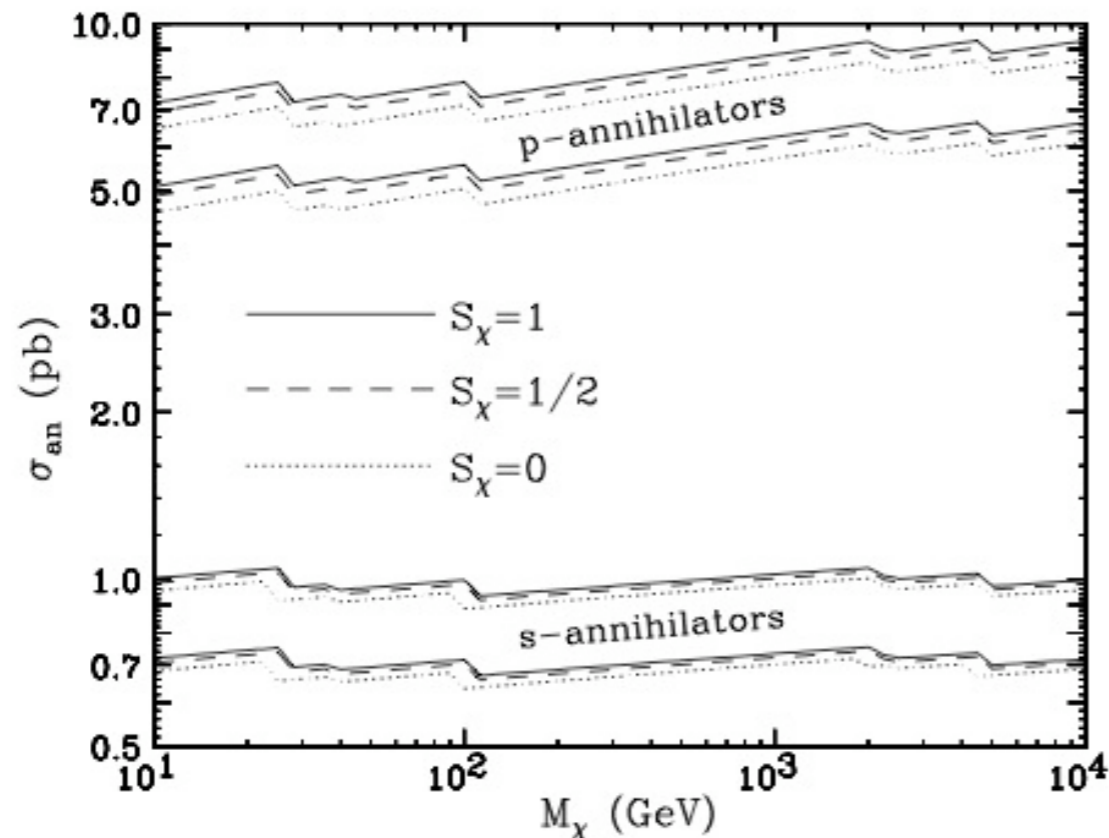


Private Higgs Dark Matter (PHDM)

- PH provides a good candidate for a WIMP
- Scalar DM:
“Gauge Singlet Scalar DM” (Zee et al., Davoudiasl et al., etc.)
and “Inert Doublet Model” (Barbieri et al., etc.)
- Any form of DM is constrained by WMAP measurements:



$$\Omega_{DM} h^2 = 0.111 \pm 0.018$$

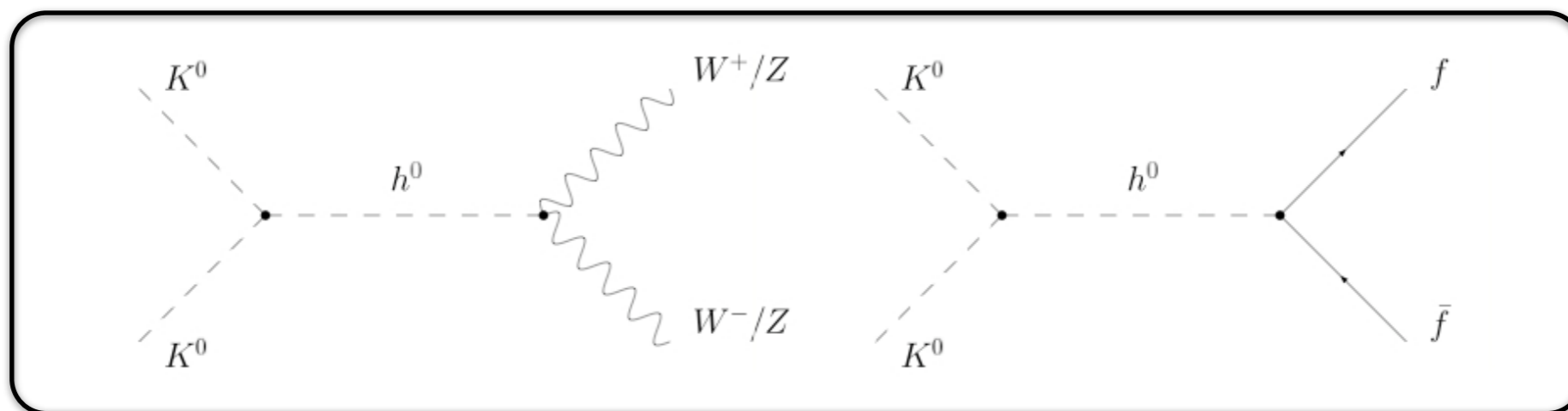


- Birkedal et al. (PRD70, 077701 (2004)): Limits on DM abundance translate into limits on annihilation cross section.
- For scalar “s-annihilator”:

$$\sigma_{an} = 0.85 \pm 0.15 \text{ pb}$$

Constraining PHDM

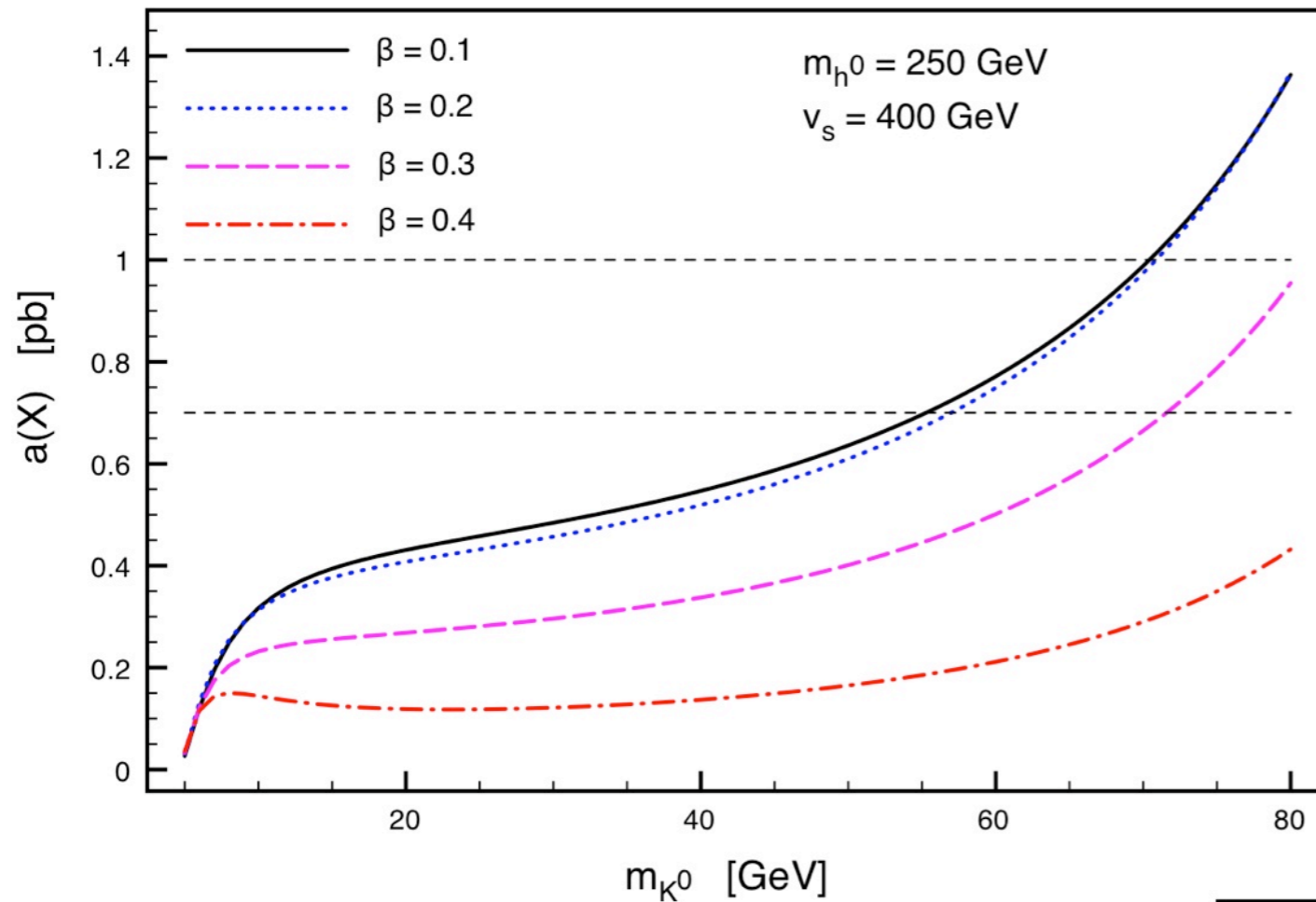
- Focus on small values of mixing between top PH and S (to avoid copious decays into SM particles)
- Then, K^0 kept in equilibrium with cosmic fluid via:



- Note: K^0 s-channel exchange suppressed by β^2 , while t-channel diagrams are suppressed by β^4 .
- Consider two scenarios:
 - “Light K^0 ” scenario: $m_K < m_W$ such that annihilation into b 's dominate
 - “Heavy K^0 ” scenario: $m_W < m_K < 2m_W$ s.t. annihilation into WW, ZZ pairs dominates

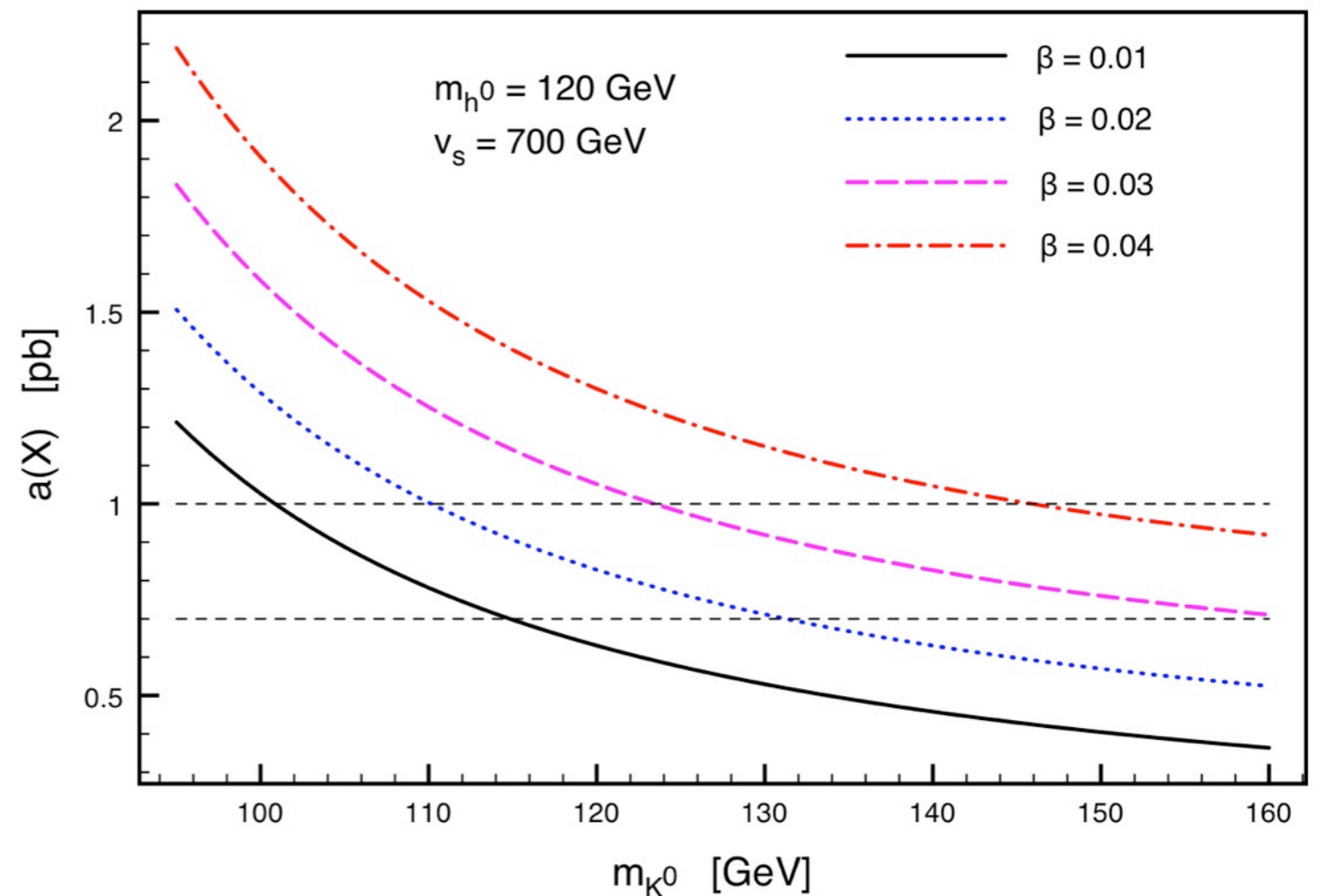
“Light K^0 ” Scenario:

- One “representative” point in parameter space... not a “full scan”
- Require $v_b \sim m_b$
- WMAP requires smaller mixing and masses $\sim 50 - 75$ GeV range



“Heavy K^0 ” Scenario:

- In general, easier to incorporate larger mass range
- Much smaller mixing angles
- Larger values of S 's v_{ev}

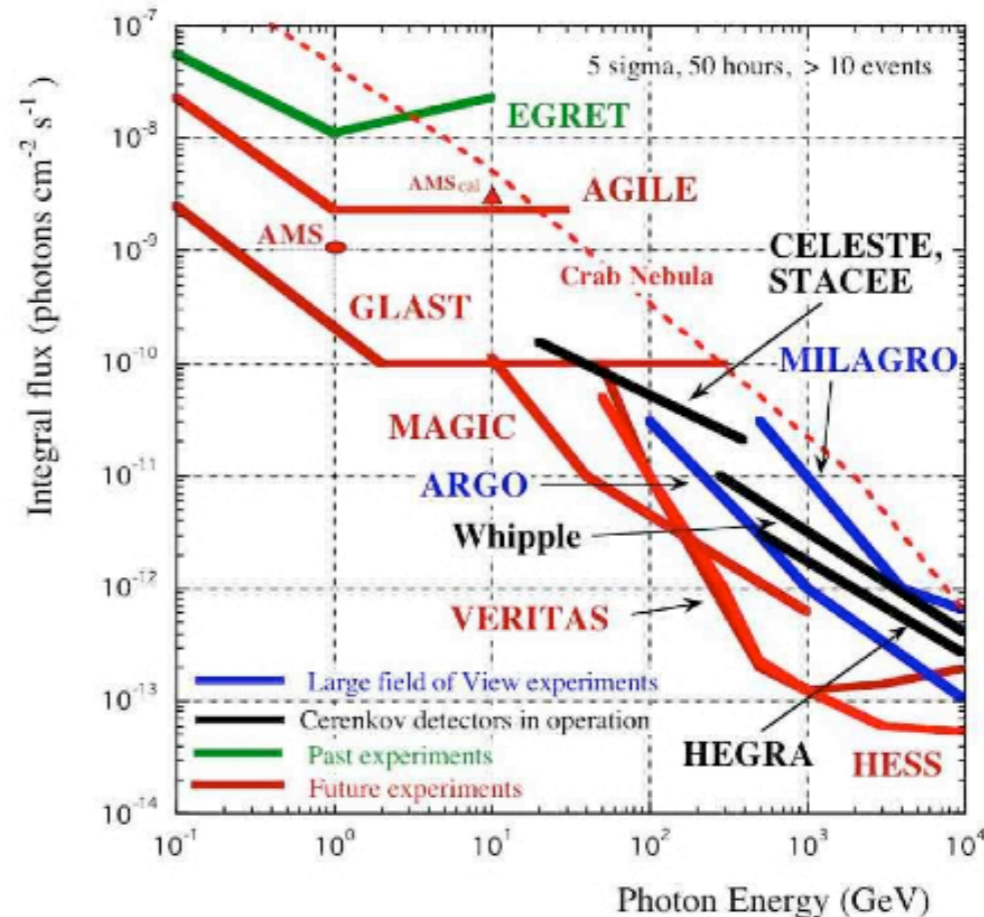


Indirect Detection of DM

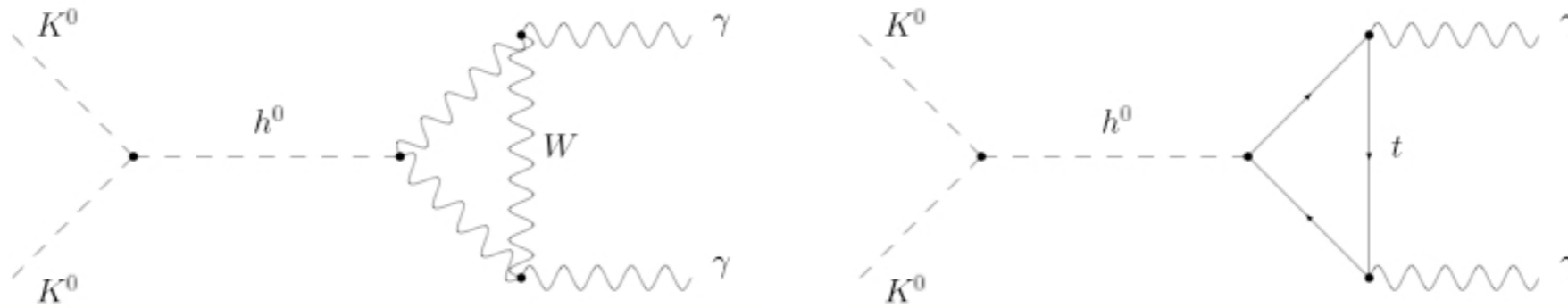
- Annihilation c.s. for DM \sim velocity-independent in non-relativistic regime
- DM collected in the galactic halo \rightarrow anomalous cosmic rays

γ rays \rightarrow information on DM

- Experiments:
 - Ground-based (atmospheric cerenkov telescopes): VERITAS, HESS, etc.
 - Space-based: GLAST
- Typical reach for DM searches:
 - ACT's $\sim 10^{-11} - 10^{-12} \text{ cm}^{-2} \text{ s}^{-1}$
 - GLAST $\sim 10^{-10} \text{ cm}^{-2} \text{ s}^{-1}$



Indirectly Detecting PHDM



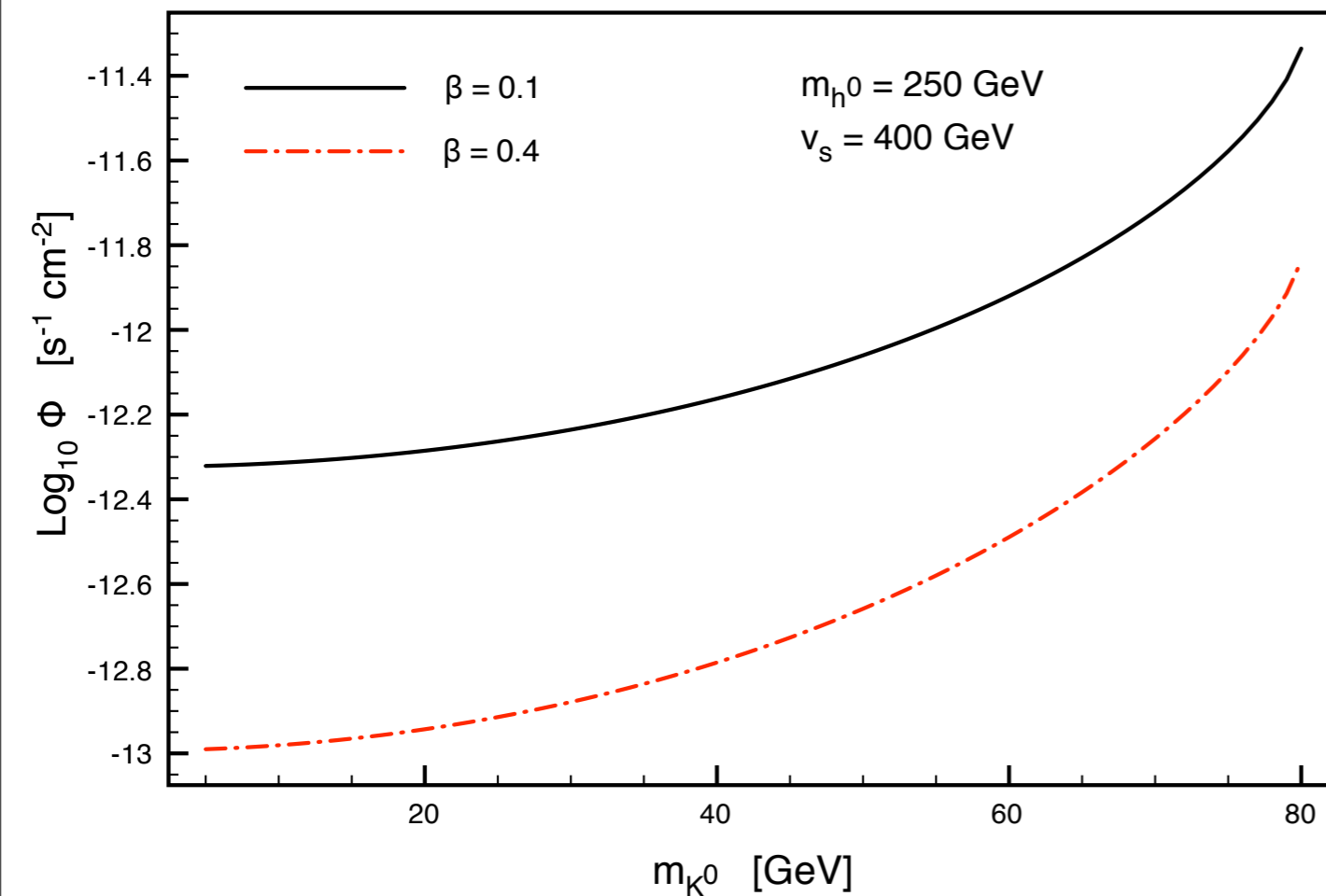
- The flux of photons observed with a line of sight $\Psi(\theta, \varphi)$ and f.o.v. $\Delta\Omega$:

$$\Phi = (1.1 \times 10^{-9} \text{ s}^{-1} \text{ cm}^{-2}) \left(\frac{\sigma_{\gamma\gamma} u}{1 \text{ pb}} \right) \left(\frac{100 \text{ GeV}}{m_{K^0}} \right) \bar{J}(\Psi, \Delta\Omega) \Delta\Omega,$$

- Dependence of flux on DM density distribution is contained in J
 - Value of J very model-dependent
 - Many models predict large spike in neighborhood of galactic center
 - $J \approx 10^3 - 10^7$ for $\Delta\Omega = 10^{-3}$ sr (typical for ACTs)

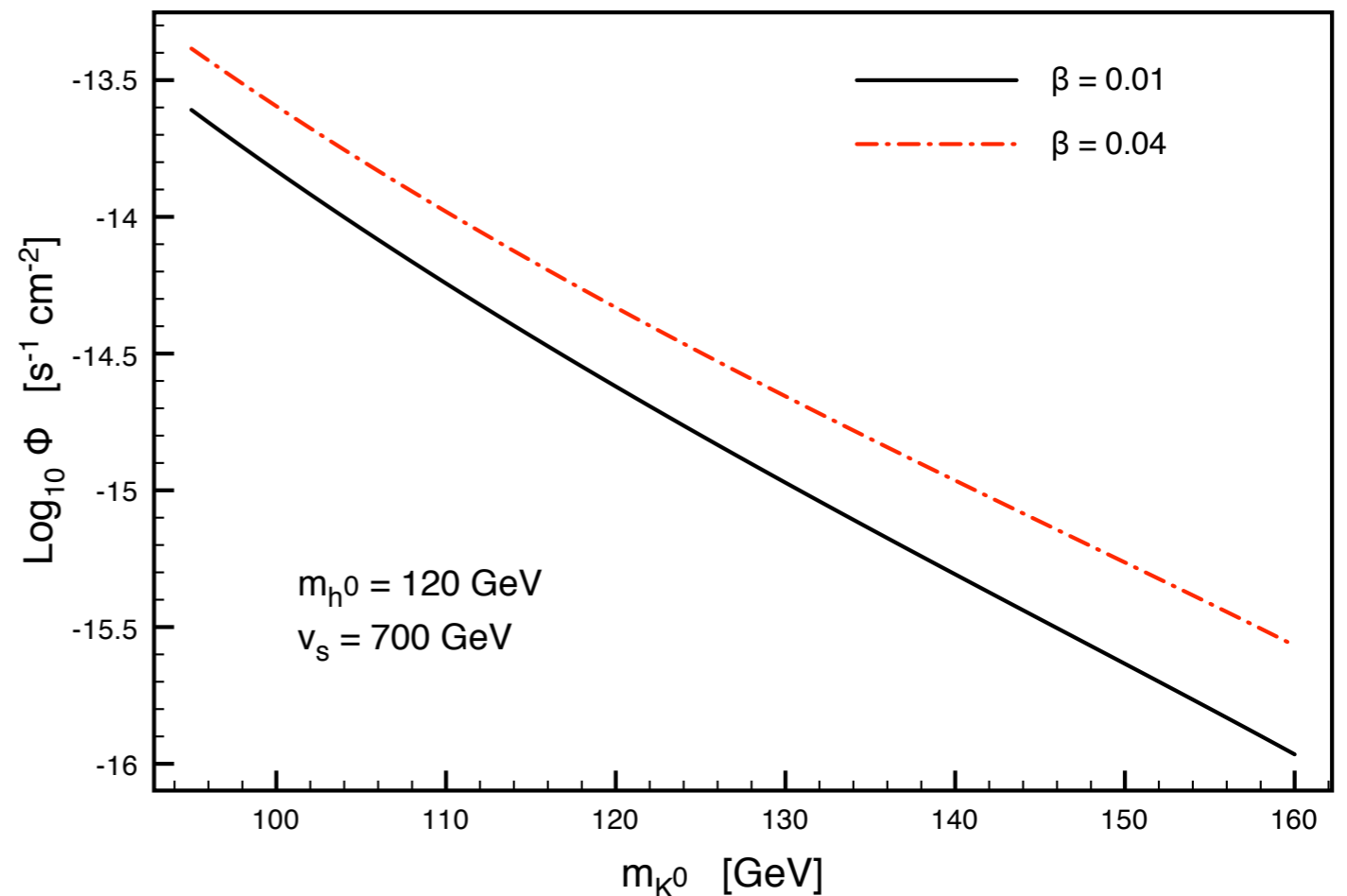
“Light K^0 ” Scenario:

- Assume no significant clumping (i.e., $J(\Psi, \Delta\Omega) \Delta\Omega = 1$)
- Beyond the reach of GLAST (reach = $10^{-10} \text{ cm}^{-2} \text{ s}^{-1}$)
- With considerable clumping (e.g., $J(\Psi, \Delta\Omega) \Delta\Omega = 10^2$), observation at ACTs possible



“Heavy K^0 ” Scenario:

- Tough to detect for larger masses
- Still possible at ACTs for larger values of clumping
- Full scan needed, though



Conclusions

- Private Higgs model provides a simple way to account for large hierarchy observed in the fermion mass spectrum
- Idea: Introduce one Higgs doublet per fermion + discrete symmetries
- Induce EWSB via vev of a gauge singlet scalar S
- PH-enomenology consists of a handful of scalars below $\sim \text{TeV}$ and a possible DM candidate
- Perturbative Unitarity: PH sector softens bounds on SM-like Higgs mass
- Private Higgs Dark Matter (PHDM):
 - Easily account for relic abundance of DM with “natural” values of parameters
 - If PHDM clumps in the region of the galactic core, observation with gamma ray telescopes is possible