Andriy Badin B_s mixing: Standard Model and Beyond

Based on: Andriy Badin, Fabrizzio Gabbiani and Alexey A. Petrov: "Lifetime difference in B_s mixing: Standard Model and beyond" Phys. Lett. B 653 (2007) 230-240

Outline

- B_S mixing in standard model with subleading 1/m² corrections included.
- Contribution from generic ΔB=1 New Physics models.
- Examples of New Physics models.

Experimental Results:

 $\Delta\Gamma/\Gamma=0.121\pm0.083/0.090$ PDG

 $\Delta\Gamma/\Gamma=0.104\pm0.076/0.084$ HFAG

The width difference between heavy and light states is given by:

$$\Delta \Gamma_{B_s} = \Gamma_L - \Gamma_H = -2\Gamma_{12} = -2\Gamma_{12} = -2\Gamma_{21}$$

Where Γ_{ij} are the elements of the decay-width matrix, $i, j = 1, 2, |1\rangle = |B_s\rangle, |2\rangle = |\bar{B}_s\rangle$

$$\Gamma_{12} = \frac{1}{2M_{B_s}} \left\langle \overline{B_s} \left| T \right| B_s \right\rangle \qquad T = \operatorname{Im} i \int d^4 x \operatorname{T} \left(H_{eff}(x) H_{eff}(0) \right)$$



Both vertices are SM:

 $\frac{G_{f}}{\sqrt{2}}V_{bq}^{*}V_{qs}\left[C_{1}\left(\overline{b}^{i}q^{j}\right)_{V-A}\left(\overline{q}^{j}s^{i}\right)_{V-A}+C_{2}\left(\overline{b}^{i}q^{i}\right)_{V-A}\left(\overline{q}^{j}s^{j}\right)_{V-A}\right]\left(\overline{b}^{i}q^{j}\right)_{V-A}\left(\overline{q}^{j}s^{i}\right)_{V-A}$

1/m expansion

$$\Gamma_{21}(B_s) = \frac{1}{2M_{B_s}} \sum_k \langle B_s | \mathcal{T}_k | B_s \rangle = \sum_k \frac{C_k(\mu)}{m_b^k} \langle B_s | \mathcal{O}_k^{\Delta B = 2}(\mu) | B_s \rangle$$

 $\Gamma_{21}(B_s) = -\frac{G_F^2 m_b^2}{12\pi (2M_{B_s})} \left(V_{cb}^* V_{cs} \right)^2 \left\{ \left[F(z) + P(z) \right] \langle Q \rangle + \left[F_S(z) + P_S(z) \right] \langle Q_S \rangle + \delta_{1/m} + \delta_{1/m^2} \right\} \right\}$

Leading order and corrections O(1/m) were calculated by several groups. For example:M. Beneke, G. Buchalla, I. Dunietz, Phys. Rev. D 54 (1996) 4419

However O(1/m) corrections are about 30% of leading contribution. There is a question how well this series converges.

Possible ways to clarify this problem:

• Redefine basis of operators in such a way that NLO corrections would be small (For example Alexander Lenz, Ulrich Nierste, JHEP 0706:072,2007)

• Compute NNLO corrections directly to see their contribution.

$$\Delta \Gamma_{B_s} = \left[0.0005B + 0.1732B_s + 0.0024B_1 - 0.0237B_2 - 0.0024B_3 - 0.0436B_4 + 2 \times 10^{-5}\alpha_1 + 4 \times 10^{-5}\alpha_2 + 4 \times 10^{-5}\alpha_3 + 0.0009\alpha_4 - 0.0007\alpha_5 + 0.0002\beta_1 - 0.0002\beta_2 + 6 \times 10^{-5}\beta_3 - 6 \times 10^{-5}\beta_4 - 1 \times 10^{-5}\beta_5 - 1 \times 10^{-5}\beta_6 + 1 \times 10^{-5}\beta_7 + 1 \times 10^{-5}\beta_8 \right] (\text{ps}^{-1}).$$



New Physics contribution



Generic description of New Physics contribution

Eugene Golowich, Sandip Pakvasa, Alexey A. Petrov Phys. Rev. Lett. 98, 181801 (2007)

$$y = \frac{\Delta\Gamma}{\Gamma} = -\frac{4\sqrt{2}G_{f}}{M_{B}\Gamma_{B}} \sum_{q,q'} V_{bq}^{*} V_{sq'} D_{qq'} \left(K_{1}\delta_{ik}\delta_{jl} + K_{2}\delta_{il}\delta_{jk}\right) \sum_{n=1}^{5} I_{n} \left\langle O_{n}^{ijkl} \right\rangle$$

$$I_{1}(y, z) = -\frac{km_{b}}{48\pi} \left(1 - 2(y + z) + (y - z)^{2}\right)$$

$$\left\langle O_{1}^{ijkl} \right\rangle = \left\langle B_{s} \left| \overline{b}_{k}\Gamma_{\mu}\gamma_{\nu}\Gamma_{2}s_{j}\overline{b}_{l}\Gamma_{1}\gamma^{\nu}\Gamma^{\mu}s_{i} \right| \overline{B_{s}} \right\rangle; \quad I_{2}(y, z) = -\frac{k}{24\pi m_{b}} \left(1 + (y + z) - 2(y - z)^{2}\right)$$

$$\left\langle O_{2}^{ijkl} \right\rangle = \left\langle B_{s} \left| \overline{b}_{k}\Gamma_{\mu}\widehat{p}\Gamma_{2}s_{j}\overline{b}_{l}\Gamma_{1}\widehat{p}\Gamma^{\mu}s_{i} \right| \overline{B_{s}} \right\rangle; \quad I_{3}(y, z) = \frac{k}{8\pi}\sqrt{z} \left(1 + y - z\right)$$

$$\left\langle O_{3}^{ijkl} \right\rangle = \left\langle B_{s} \left| \overline{b}_{k}\Gamma_{\mu}\widehat{p}\Gamma_{2}s_{j}\overline{b}_{l}\Gamma_{1}\widehat{p}\Gamma^{\mu}s_{i} \right| \overline{B_{s}} \right\rangle; \quad I_{4}(y, z) = -\frac{k}{8\pi}\sqrt{y} \left(1 - y + z\right)$$

$$\left\langle O_{4}^{ijkl} \right\rangle = \left\langle B_{s} \left| \overline{b}_{k}\Gamma_{\mu}\widehat{p}\Gamma_{2}s_{j}\overline{b}_{l}\Gamma_{1}\Gamma^{\mu}s_{i} \right| \overline{B_{s}} \right\rangle; \quad I_{5}(y, z) = \frac{km_{b}}{4\pi}\sqrt{zy}$$

$$\left\langle O_{5}^{ijkl} \right\rangle = \left\langle B_{s} \left| \overline{b}_{k}\Gamma_{\mu}\Gamma_{2}s_{j}\overline{b}_{l}\Gamma_{1}\Gamma^{\mu}s_{i} \right| \overline{B_{s}} \right\rangle; \quad k = \frac{m_{b}}{2} \left(1 - 2(y + z) + (y - z)^{2}\right)^{V_{2}};$$

$$y = \left(\frac{m_{q}}{m_{b}}\right)^{2}; z = \left(\frac{m_{q}}{m_{b}}\right)^{2}$$

 $\mathbf{I}_{1,2}$ and $D_{qq'}$ are parameters of a New Physics model

Multi Higgs Model

This model contains charged Higgs bosons as parts of the extended Higgs sector and provides new flavor-changing interactions mediated by charged Higgs bosons

$$\mathcal{H}_{ChH}^{\Delta B=1} = -\frac{\sqrt{2}G_F}{M_H^2} \ \overline{b}_i \overline{\Gamma}_1 q_i' \ \overline{q}_j \overline{\Gamma}_2 s_j$$

where $\overline{\Gamma}_i, \ i = 1, 2$ are

 $\overline{\Gamma}_{1} = m_{b}V_{cb}^{*} \cot\beta P_{L} - m_{c}V_{cb}^{*} \tan\beta P_{R}$ $\overline{\Gamma}_{2} = m_{s}V_{cs} \cot\beta P_{R} - m_{c}V_{cs} \tan\beta P_{L}$

For values of $M_H = 85 GeV$ and $\cot \beta = 0.05$ it gives $y_{ChH} \approx 0.006$ Dependence of y_{ChH} on mass of Higgs boson: solid line - $\tan \beta = 20$, dashed line - $\tan \beta = 10$, dotted line - $\tan \beta = 5$, dash-dotted line - $\tan \beta = 3$



Left-Right Symmetric Model

Left–Right Symmetric Model (LRSM) assumes the extended $SU(2)_L xSU(2)_R$ symmetry of the theory, which restores parity at high energies

$$\overline{\Gamma}_{1,2} = \gamma^{\mu} P_R$$

$$D_{qq'} = V_{cb}^{*(R)} V_{cs}^{(R)} \frac{G_F^{(R)}}{\sqrt{2}}$$

here $\frac{G_F^{(R)}}{\sqrt{2}} = g_R^2 / 8M_{W(R)}^2$ and for future calculations we take $g_L = kg_R$

One of possible realizations of such scenario which gives the biggest numerical value of y_{LR} is a "Non-manifest LR" $(V_{ij}^{(R)} \approx 1)$ with $M_{W^{(R)}} = 1 \ TeV$ value of $y_{LR} \approx -0.04$ was obtained. In case of "manifest LR" $((V_{ij}^{(R)} = V_{ij}))$ contribution from this model is less.

Dependence of y_{LR} on mass of $W^{(R)}$ boson. Solid line "Non manifest" case, dashed - "manifest"



Dependence of $\Delta\Gamma/\Gamma$ on $M_W^{(R)}$ in non-manifest LRSM. Solid line: $\kappa = 1$ dashed line: $\kappa = 1.5$, dotted line: $\kappa = 2$.



Conclusions

- The most precise theoretical prediction for Bs lifetime difference obtained. Corrections of the order of 1/m_b² included and appeared to be small. Results are in good agreement with available experimental data
- We considered the most general four-fermion effective Hamiltonian, which can be generated by any reasonable extension of the Standard Model and derived its contribution to Γ_{Bs}
- Several models considered:
 - Multi Higgs Model can increase lifetime difference. Contribution from this model of NP is rather small. y = 0.006, compare to SM – y=0.10
 - Contribution from Left-Right Symmetric Model is rather large and negative :y= -0.04, compare to SM y = 0.10. Studies of mixing help to put constraints onto parameters of this model.