

HIGGS DECAYS AND BRANE GRAVI-PHOTONS

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Outline

- Massive vector fields common feature of many locally invariant brane world models
- Couple vector to the Standard Model via intrinsic and extrinsic curvature as well as vector field strength
- Constrain vector parameters using LEP I, II accelerator data.
- Examine Higgs decays containing brane vectors.

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Presence of 4-d flexible probe brane breaks extra dimensional space-time symmetries

Focus on case of $N \geq 2$ isotropic extra co-dimensions

Associated with broken translations are Nambu-Goldstone boson fields $\phi_i(x)$, $i = 1, \dots, N$ whose dynamics describes motion of probe brane into the extra dimensions

Brane oscillations induces metric on the brane

Dynamics given by Nambu-Goto action

$$S_{NG} = -F_X^4 \int d^4x \sqrt{1 + \frac{1}{F_X^4} \partial_\mu \phi_i \partial^\mu \phi_i}$$

Identify F_X^4 as brane tension

Make extra dimensional translations locally invariant as part of higher dimensional general covariance

Introduce “gravi-photon” gauge field $X_i^\mu(x)$ and replace derivatives by covariant derivative

S_{NG} contains vector mass term: Higgs mechanism

Vector mass M_X is independent mass scale; its value is model dependent

Presence of massive vector common feature of many locally invariant brane world models

Massive vector field Proca action:

$$S_{Proca} = -\frac{1}{4} \int d^4x X_i^{\mu\nu} X_{\mu\nu i} - \frac{1}{2} M_X^2 \int d^4x X_i^\mu X_{\mu i}$$

with field strength $X_i^{\mu\nu} = \partial^\mu X_i^\nu - \partial^\nu X_i^\mu$

Invariant Couplings of XX to Standard Model

Global $O(N)$ symmetry associated with possible orientations of embedding 4-d brane in $D=4+N$ space

X_i is $SU(3) \times SU(2) \times U(1)$ singlet, but transforms as N under the global $O(N)$ (carries label i)
Standard Model fields are $O(N)$ singlets

$O(N)$ invariant couplings to Standard Model require even powers of X_i ; **Massive vectors are stable particles**

• **Induced metric coupling to Standard Model symmetric energy momentum tensor $T_{SM}^{\mu\nu}$**

$$S_{XXT} = \frac{M_X^2}{2F_X^4} \int d^4x X_{\mu i} X_{\nu i} T_{SM}^{\mu\nu}$$

Extrinsic curvature

$$K_i^{\mu\nu} = \frac{1}{2}(\partial^\mu X_i^\nu + \partial^\nu X_i^\mu) + \dots$$

Measures curvature of embedded brane relative to enveloping D-dimensional geometry

- **Coupling to weak hypercharge field strength**

$$S_{XXB} = \frac{M_X^2}{2F_X^4} \int d^4x (\kappa_1 B_{\mu\nu} + \kappa_2 \tilde{B}_{\mu\nu}) X_i^{\mu\rho} K_{i\rho}^\nu$$

$$B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu \text{ with } B_\mu = \cos\theta_W A_\mu - \sin\theta_W Z_\mu$$

- **Coupling to scalar doublet bilinear**

$$S_{XXH} = \frac{M_X^2}{2F_X^4} \int d^4x \left[\lambda_1 K_i^{\mu\nu} K_{i\mu\nu} + \lambda_2 X_{i\mu\nu} X_i^{\mu\nu} + \lambda_3 X_{i\mu\nu} \tilde{X}_i^{\mu\nu} \right] \phi^\dagger \phi$$

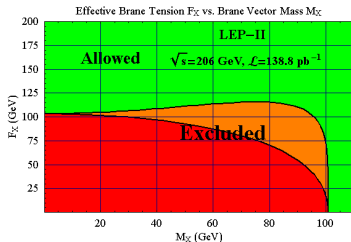
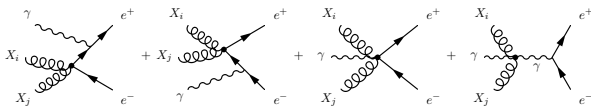
Unitary gauge: $\phi^\dagger \phi = vH + \frac{1}{2}H^2$ with H Higgs scalar

- LEP limits

$e^+e^- \rightarrow \gamma XX$ appears as γ plus missing energy

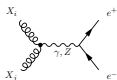
Expt limit: $\sigma(e^+e^- \rightarrow \gamma \cancel{E}) < .45 \text{ pb}$ leads to restriction of allowed M_X, F_X values

- Induced metric coupling (to $T_{SM}^{\mu\nu}$) only

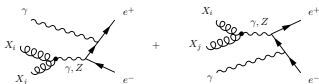


- Coupling to $B^{\mu\nu}$ only

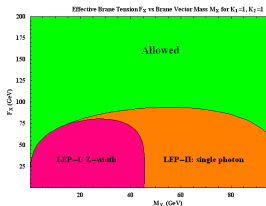
Allowed invisible Z decay width: $\Gamma_{Z \rightarrow XX} \leq 2 \text{ MeV}$



and limits from $e^+e^- \rightarrow \gamma XX$



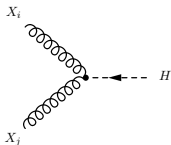
Transverse vector modes required



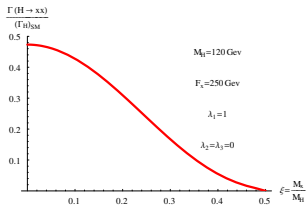
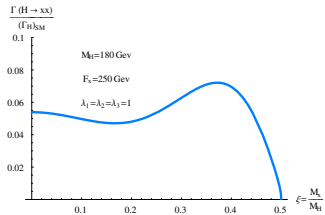
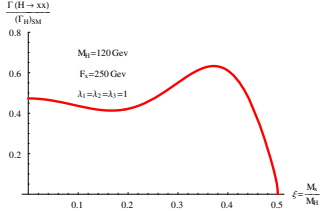
No interference in $e^+e^- \rightarrow \gamma XX$ cross section between graphs containing from S_{XXT} and S_{XXB}

Higgs decays containing X vectors

- $H \rightarrow XX$: Appears as invisible Higgs decay

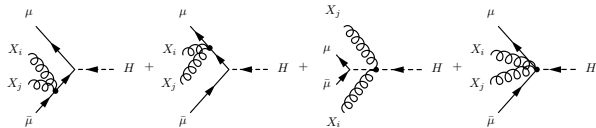


$$\frac{\Gamma(H \rightarrow XX)}{(\Gamma_H)_{SM}} = \frac{1}{(\Gamma_H)_{SM}} \frac{N \sin^2 \theta_W M_H^7 M_W^2}{32\pi^2 \alpha F_X^8} \sqrt{1 - 4\xi^2} \left[\left(\frac{\lambda_1}{4} + \lambda_2 \right)^2 \left(\frac{1}{4} - \xi^2 + 3\xi^4 \right) + \left(\frac{\lambda_1^2}{16} - \lambda_2 \right) \frac{1}{2} (1 - 2\xi^2)^2 (1 - 4\xi^2) + \frac{1}{4} \left(\frac{\lambda_1}{4} - \lambda_2 \right)^2 (1 - 4\xi^2)^2 + 8\lambda_3^2 \xi^4 (1 - 4\xi^2) \right] ; \quad \xi \equiv \frac{M_X}{M_H}$$



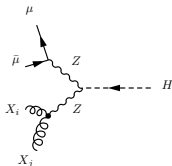
- $H \rightarrow \mu \bar{\mu} X X$

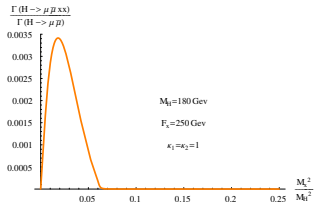
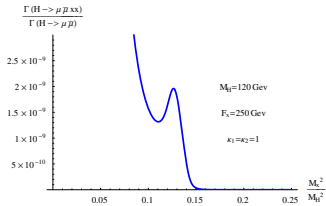
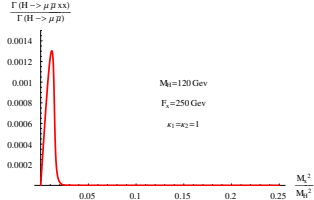
Couplings involving $(T_{\mu\nu})_{SM}$



Amplitude varies as $\frac{m_\mu}{F_X^4}$.

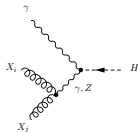
In general, suppressed relative to graph containing extrinsic curvature coupling S_{XXB}



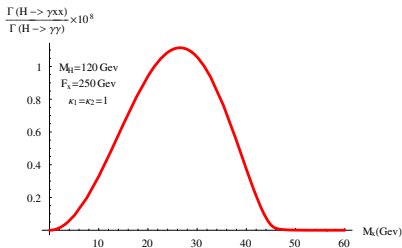


- $H \rightarrow \gamma XX$

Appears as photon plus missing energy



For the $H\gamma\gamma$ and $H\gamma Z$ vertices use the effective couplings $F_{\mu\nu}F^{\mu\nu}H$ and $F_{\mu\nu}Z^{\mu\nu}H$; For the γXX and ZXX vertices use S_{XXB}



Summary

- 4-d probe brane breaks extra dimensional translation invariances. Gauging broken translations leads to massive, Standard Model singlet, stable Proca vector fields
- Coupled vectors to the Standard Model using both intrinsic and extrinsic curvatures
- Limits on effective Lagrangian parameters (M_X, F_X) from LEP I, LEP II data
- Rate for $H \rightarrow XX$ (invisible Higgs decay) could prove comparable to any Standard Model Higgs decay rate for presently allowed M_X, F_X values
- $\Gamma(H \rightarrow \mu\bar{\mu}XX) \sim 10^{-3}\Gamma(H \rightarrow \mu\bar{\mu})$ for small range of M_X values which varies as function of M_H