

LRS@LHC

Hooman Davoudiasl

Brookhaven National Laboratory

Based on: H. D., G. Perez, and A. Soni

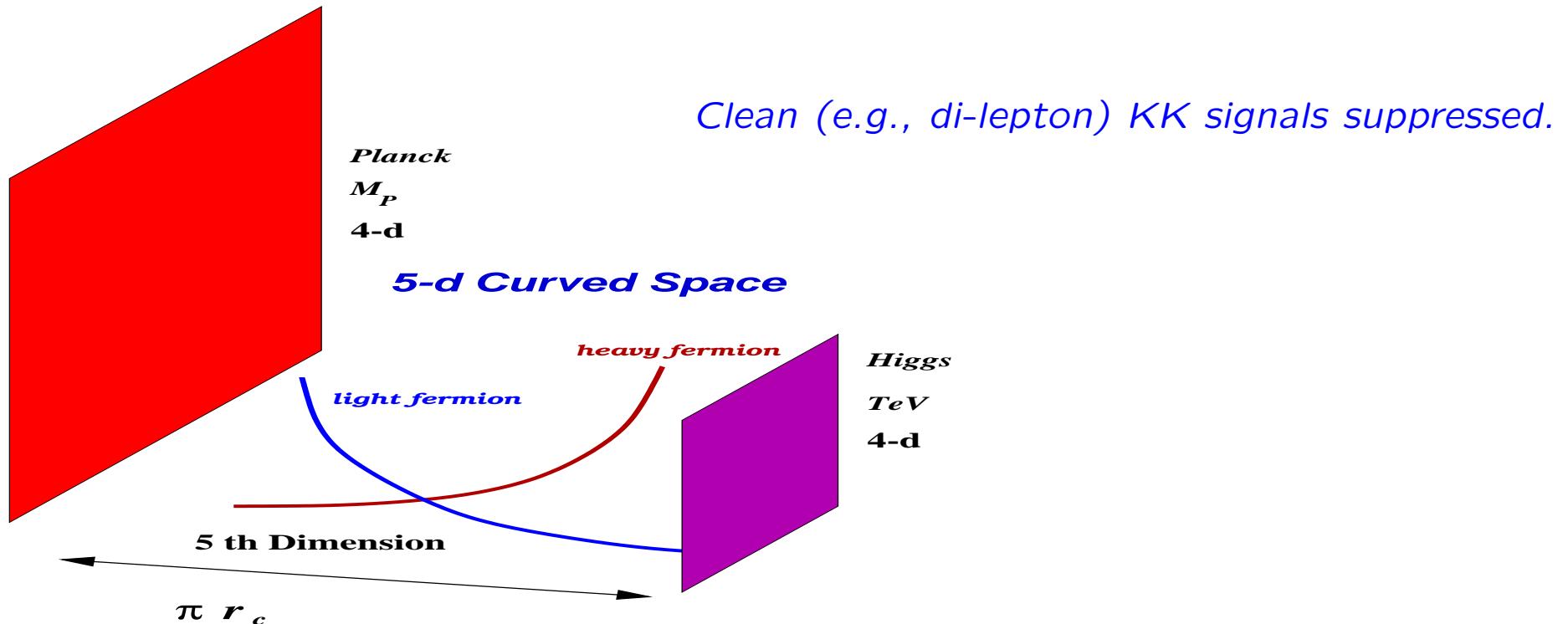
arXiv:0802.0203 [hep-ph]

Pheno 2008

April 28–30, 2008, University of Wisconsin-Madison

Introduction

- Warped 5D Randall-Sundrum (RS) model: Planck-weak hierarchy.
 - $e^{-kr_c\pi} M_5 \sim \text{TeV}$, $kr_c\pi \approx 35$ for $M_5 \sim M_P$. [Randall, Sundrum, 1999](#)
 - Distinct signatures: spin-2, ... TeV-scale resonances.
- SM in 5D bulk, fermion profiles \rightarrow Realistic model of flavor.
- Bulk fermion masses: light (heavy) fermions UV(IR)-localized.



Oblique Corrections

$$S_{tree} \approx 2\pi (v/\kappa)^2 \quad ; \quad T_{tree} \approx \frac{\pi}{2\cos\theta_W^2} (v/\kappa)^2 \underbrace{(kr_c\pi)}_{\approx 35}$$

- $\kappa \equiv ke^{-kr_c\pi} \sim m_{KK}^{\text{gauge}}/2.5$
- KK-tower mixing via EWSB (modified zero-modes) $\propto kr_c\pi$
- $m_{KK}^{\text{gauge}} \sim 3 \text{ TeV} \rightarrow T_{tree} \sim 3 \text{ (RS)}$; data: $|S| \sim |T| \sim 0.1 - 0.3$.

Little Randall-Sundrum (LRS) Model:

- Truncated; $kr_c\pi = 6$ ($M_5 \sim 10^3 \text{ TeV}$).
- Flavor same as RS: λ_5 (5D Yukawa) and fermion IR-profiles.
- Suppression of T_{tree} in LRS.

$$m_{KK}^{\text{gauge}} \sim 3 \text{ TeV} \rightarrow T_{tree} \sim 3/y$$

LRS truncation factor: $y \equiv (kr_c\pi|_{RS})/(kr_c\pi|_{LRS}) \approx 6$.

- S_{tree} : Largely *universal* shift in light fermion-gauge coupling.
 - (i) zero-mode-KK mixing $\sim \sqrt{kr_c\pi}$.
 - (ii) Fermion coupling to KK modes $\sim 1/\sqrt{kr_c\pi}$.

$\Rightarrow S_{tree}$ unchanged after LRS truncation.

- δT : UV-sensitive loops, cutoff operators: $m_{KK} \gtrsim 10$ TeV.
 - LRS: δT loop and cutoff contributions same as RS.

Need gauged 5D custodial symmetry:

$$SU(2)_L \times SU(2)_R \times U(1)_X.$$

- T : Eliminate tree, cutoff, UV-sensitive loops.
- $S, T \Rightarrow m_{KK} \gtrsim 3$ TeV. [Agashe, Delgado, May, Sundrum, 2003](#)
- Same for RS and LRS.

Non-Oblique Corrections & Flavor Physics

- LRS: λ_5 unscaled \rightarrow non-universality unchanged vis-à-vis RS.

- RS and $Zb\bar{b}$:

- $m_{KK} \gtrsim 3$ TeV with custodial symmetry and \mathbb{Z}_2 .
- Otherwise: $m_{KK} \gtrsim 5$ TeV. [Agashe, Contino, Da Rold, Pomarol, 2006](#)

- LRS and $Zb\bar{b}$:

- Gauge KK-mixing: truncated.
- Untruncated: $\underbrace{b_L - b_R^{KK}}_{\text{Yukawa}} \propto m_{KK}(b_R)^{-2} \rightarrow m_{KK}(b_R) \gtrsim 4$ TeV
- Implies $m_{KK}(\text{gauge}) \gtrsim 3$ TeV in LRS! (no \mathbb{Z}_2 , but custodial for T)

- $\Delta F = 2$: tree-level KK gluon exchange contribution to ϵ_K .

[Agashe, Perez, Soni, 2004](#)

- Dominant contribution from $(V - A) \times (V + A)$: [UTfit Collaboration, 2007](#)
- $\rightarrow m_{KK} \gtrsim 20$ TeV (**RS**); $\mathcal{O}(30\%)$ uncertainty [Csaki, Falkowski, Weiler, 2008](#)

LRS bound smaller by $1/\sqrt{y} \approx 1/2.5$.

LRS Phenomenology and Golden Modes

- $g_{KK}|_{UV} \sim g_4/\sqrt{kr_c\pi}, \quad g_{KK}|_{IR} \sim g_4\sqrt{kr_c\pi}.$
- (i) Broad KK states become narrower by y .
- (ii) Width into light states (e^+e^- , $u\bar{u}$, ...) enhanced by $y \rightarrow \text{BR} \sim y^2$.
- (iii) $\sigma(f_i\bar{f}_j \rightarrow KK \rightarrow f_k\bar{f}_l) \propto \Gamma(KK \rightarrow f_i\bar{f}_j) \text{BR}(KK \rightarrow f_k\bar{f}_l)$
- (i) \oplus (ii) \oplus (iii) $\Rightarrow \mathcal{S} \sim y^3$ and $\mathcal{B} \sim 1/y$ (over the width); $\mathcal{S}/\mathcal{B} \sim y^4$.

LRS, $y \approx 6 \Rightarrow \mathcal{S} \rightarrow \mathcal{O}(100)\mathcal{S}$; $\mathcal{S}/\mathcal{B} \rightarrow \mathcal{O}(1000)\mathcal{S}/\mathcal{B}$!

$M_{Z'} \sim 4 - 5 \text{ TeV}$ and $L = 100 \text{ fb}^{-1}$: $Z' \rightarrow \ell^+\ell^-$, $\ell = e, \mu$.

Compare with RS: $M_{Z'} \sim 2 \text{ TeV}$ and $L = 1000 \text{ fb}^{-1}$. [Agashe et al., 2007](#)

Revived prospects for golden modes!

Holography

- $1/g_4^2 = \tau_{\text{UV}} + \tau_{\text{IR}} + \underbrace{\log(k/\kappa)}_{kr_c\pi}/(kg_5^2)$ $\tau_{\text{UV},\text{IR}}$ small QM corrections
- Dual large N CFT: $kg_5^2 \sim 16\pi^2/N$.
LRS dual to *larger N CFT* : $N^{\text{LRS}} \sim yN^{\text{RS}}$.
- $N \rightarrow \infty$: inter-composite interactions weaker (e.g. T_{tree}).
- S_{tree} : mainly universal vertex correction, unchanged.
[Mixing] \times [light fermion coupling]: $\sim 1/\sqrt{N} \times \sqrt{N}$.
- LRS: λ_5 unchanged \rightarrow separate “Flavor CFT”: $N_F \sim N^{\text{RS}} < N^{\text{LRS}}$.
Non-universality from compositeness held fixed.

Truncation & IR Probes of UV Scales:

- Other truncations are possible:
 - TeV $\rightarrow 10^{10}$ GeV (ν_R mass, . . .): $y \approx 2$; $\mathcal{S} \rightarrow 10\mathcal{S}$.
 - TeV $\rightarrow 10^{15}$ GeV (GUT): $y \approx 1.3$; $\mathcal{S} \rightarrow 2\mathcal{S}$.
- TeV data: Relative BR into IR (H, t, \dots) and UV (e, μ, \dots) channels.
 \Rightarrow Size of 5D slice/conformal window.

UV-Completion:

- Recent LRS UV-model: K. McDonald arXiv:0804.0654 [hep-th].
 - 6D geometry, 2 warped directions.
 - LRS a 5D slice (4-brane), with $M_5 \sim 10^3$ TeV \rightarrow TeV.
 - Warping along the 6th dimension: $M_6 \sim M_P \rightarrow 10^3$ TeV.
 - Generalized to n -warped spacetime.

Summary & Conclusions

- LRS: predictive warped model of flavor, cutoff at $\mathcal{O}(10^3)$ TeV.
- Separate gauge and flavor dynamics.
- Many contributions to precision data suppressed.
- Return of the golden modes (e.g. di-leptons)!
 - LHC discovery prospects much better than the RS counterpart.
- Dual CFT: larger N dynamics for the weakly gauged sector.
- Collider physics very sensitive to truncation.
 - TeV data → UV-scale/conformal window.
- UV-completion: stabilize $\Lambda_{UV}^{LRS} \sim 10^3$ TeV (hierarchy).