

T-Anomaly Induced Decays in Little Higgs Model with T-Parity

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Little hierarchy and the little Higgs

SM:

Higgs mass = 115~ 200 GeV requires: $\left| \frac{\delta_q m^2}{m^2} \right| \leq 10 \Rightarrow \Lambda_{\text{SM}} \lesssim 2 - 3 \text{ TeV}$

Precision Electroweak data requires: $\Lambda_{\text{SM}} \gtrsim 10 \text{ TeV}$

A Little Higgs:

a pseudo-Goldstone boson protected by collective symmetry breaking

The SU(5)/SO(5) construction

Arkani-Hamed, Cohen,
Kalz, Nelson, 2002

SU(5) with gauged $[SU(2)_1 \times U(1)_1] \times [SU(2)_2 \times U(1)_2]$

Each $SU(2) \times U(1)$ leaves a SU(3) symmetry intact

Higgs mass needs both $SU(2) \times U(1)$ groups

\Rightarrow leading vector boson loop diagrams lead
no worse than logarithmic divergence

Top sector

Top contribution canceled by a new fermion $t_p \Rightarrow$ reduced effective top Yukawa

New particles

$$\begin{array}{ccc} \text{SU}(5) & \xrightarrow{\Sigma_0} & \text{SO}(5) \\ [\text{SU}(2) \times \text{U}(1)]^2 & & \text{SU}(2)_L \times \text{U}(1) \end{array}$$

- * Massive complex triplet Higgs \mathbf{f}
Massive partners of EW gauge fields
Another top quark

The LH model doesn't solve the little hierarchy problem

* Naturalness: $f \sim 1 \text{ TeV}$ (Casa et.al, 2005)

* precision EW: $f > 4.6 \text{ TeV} (2\sigma)$ (Csaki et.al., 2003)

T parity (Chern and Low, 2004)

- \Rightarrow * triplet f odd, doublet h even under T
- * $H^\dagger f H$ term forbidden, no triplet vev
- * extended fermion sector:
 - T -odd partner to every LH fermion
- * no mixing between W / W_H or Z / Z_H or between f / f_H :
(T -odd particles leave EW physics alone)

OK with precision EW (Casa et.al.,2005)

NOT that OK with fine tuning (J. Hubisz, P. Meade, 2005)

Production rate

Tree level @ LHC, $\sqrt{s}=1.5$ TeV, $m_H=200$ GeV

Final state	σ [fb]
$q^+ q^-$	5.2
$q^+ q^+$	2.6
$T_- \bar{T}_-$	1.5
$q_- W_H^+ + q_- W_H^-$	1.8
$q_- Z_H$	0.90
$Z_H W_H^+ + Z_H W_H^-$	1.6
$W_H^+ W_H^-$	1.0

Final state:	σ [fb]
$q_- q_-$	9.7
$f X$	0.005
$A_H X$	0.33
$A_H A_H$	0.003

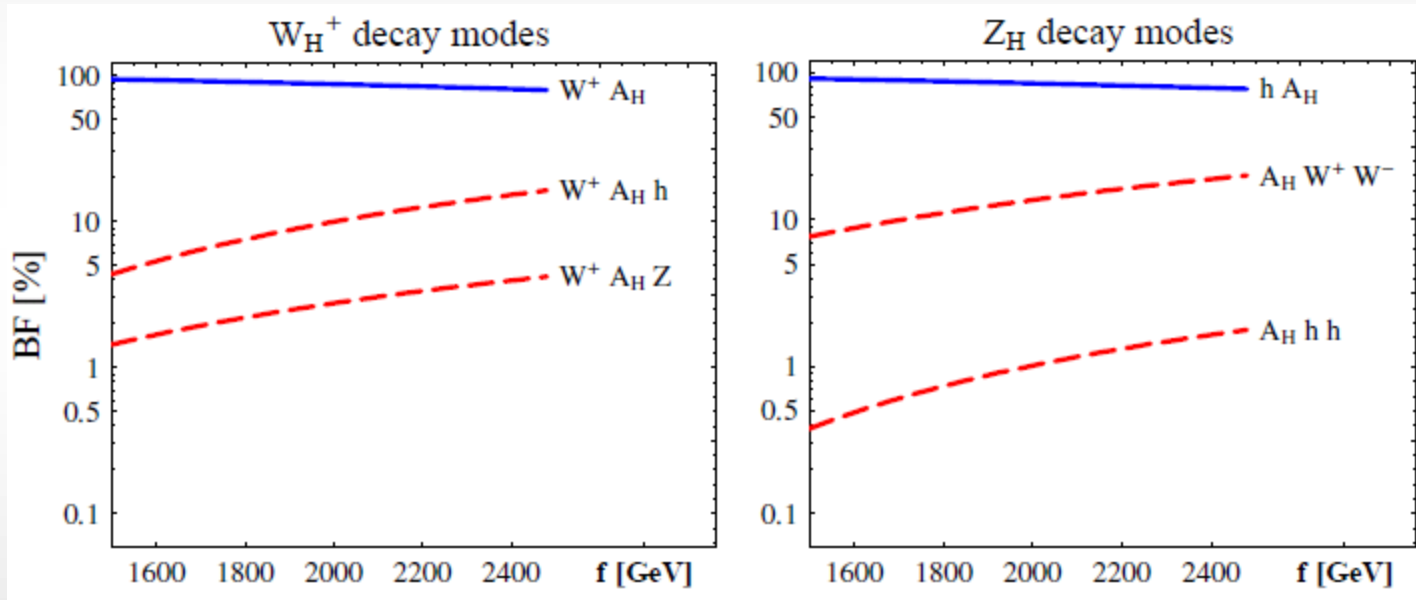
T conservation means:

T -odd particles must be produced in pairs

T_+ is T even and can be produced with SM quarks

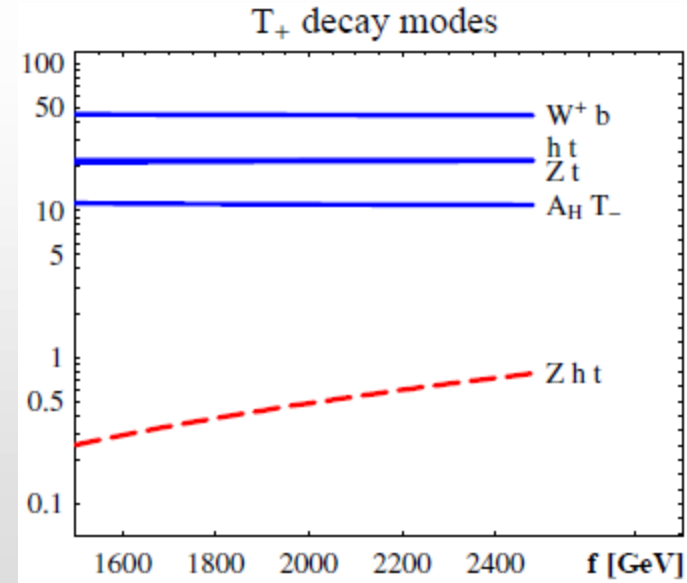
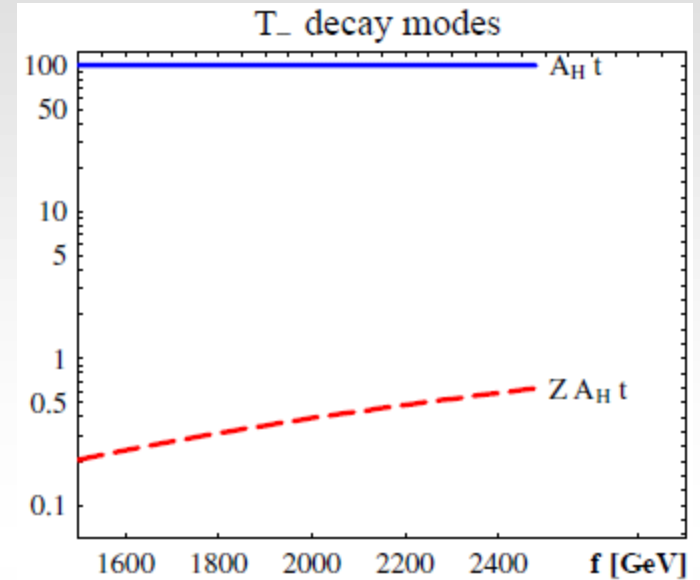
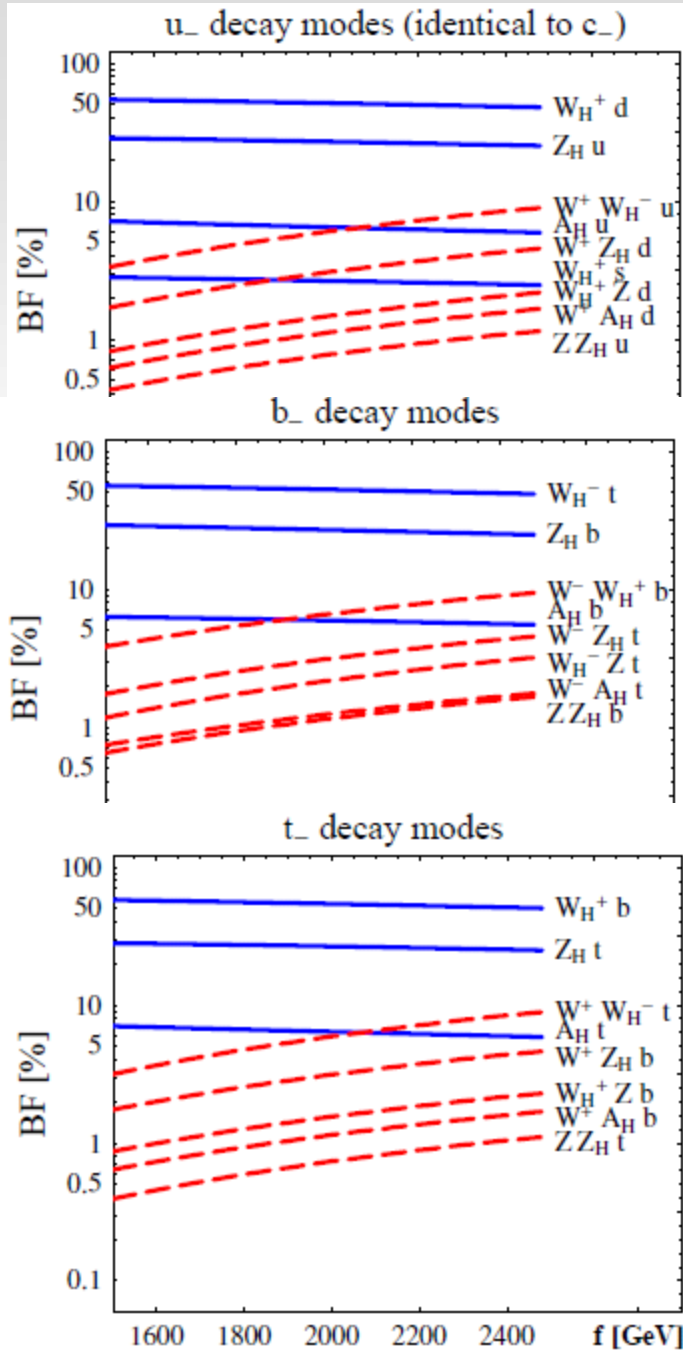
All T -odd particles decay into A_H

\cancel{E}_T factory at LHC?



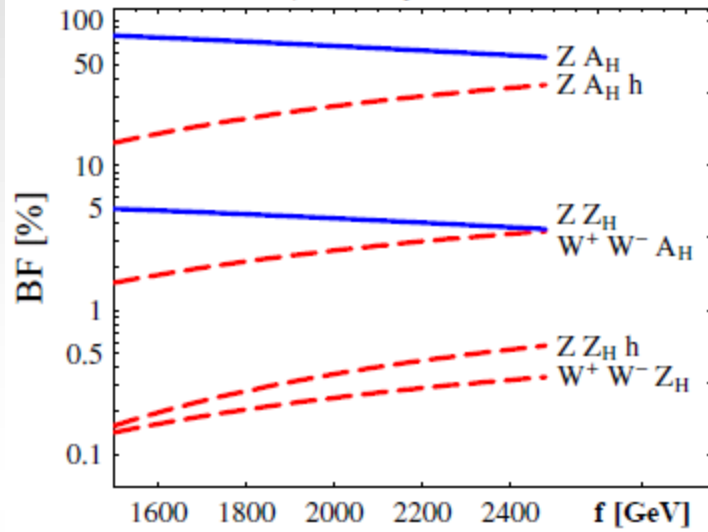
Weak gauge bosons

Heavy quark decays

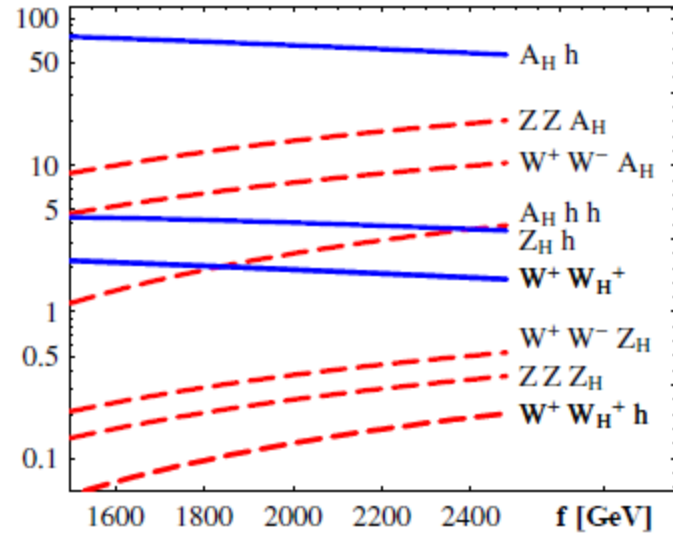


Heavy Higgs fields

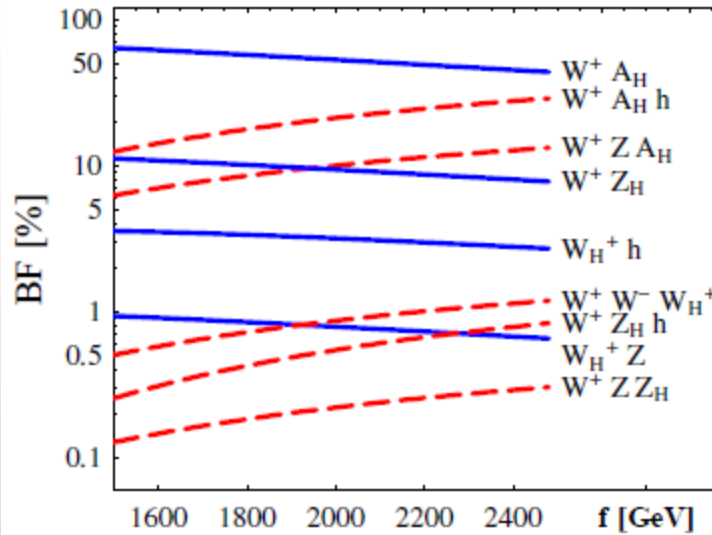
ϕ^0 decay modes



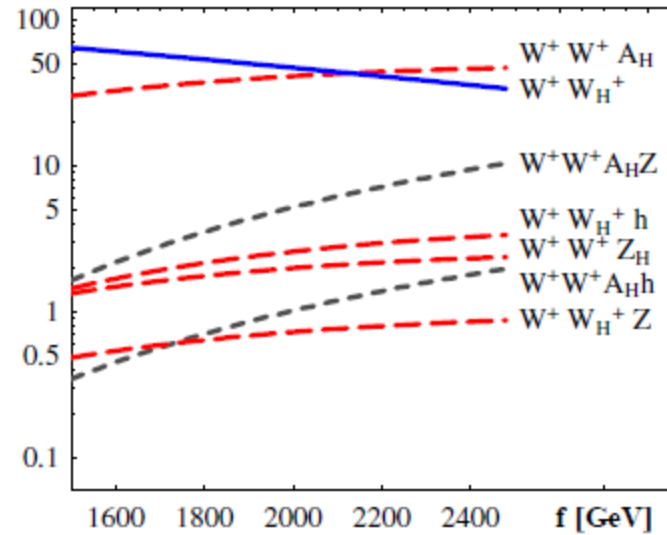
ϕ^p decay modes



ϕ^+ decay modes



ϕ^{++} decay modes



T -violation (A_H isn't stable) (C. Hill, R. Hill, 2007)

Similarity between T parity and KK parity

A 5D Lagrangian \longrightarrow The Chern-Simmons term
breaks the T -parity

4D
 \swarrow

Wess-Zumino-Witten term \longrightarrow T parity broken
in 4D

$$\Gamma_{WZW} = \int d^4x \frac{\tilde{g}N}{24\sqrt{3}\pi^2} \epsilon^{\mu\nu\rho\sigma} \tilde{B}_\mu [$$

$$-\frac{1}{3}g_1^2 [B_\nu \partial_\rho B_\sigma] + 2g_2^2 \text{Tr}[W_\nu \partial_\rho W_\sigma] - \frac{3ig_2^3}{2} \text{Tr}[W_\nu W_\rho W_\sigma]$$

$$+ \frac{ig_1}{2F^2} F_{\nu\rho}^B [H^\dagger (D_\sigma H) - (D_\sigma H^\dagger) H] - \frac{ig_2}{F^2} [H^\dagger F_{\nu\rho}^W (D_\sigma H) - (D_\nu H^\dagger) F_{\rho\sigma}^W H]$$

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Wess-Zumino-Witten term \longrightarrow T parity broken in 4D

$$\Gamma_{WZW} = \int d^4x \frac{\tilde{g}N}{24\sqrt{3}\pi^2} \epsilon^{\mu\nu\rho\sigma} \tilde{B}_\mu \left[\frac{1}{3} \text{Tr}[D_\nu \partial_\rho D_\sigma] + 2g_2^2 \text{Tr}[W_\nu \partial_\rho W_\sigma] - \frac{3ig_2^3}{2} \text{Tr}[W_\nu W_\rho W_\sigma] \right]$$

$$+ \frac{ig_1}{2F^2} F_{\nu\rho}^B [H^\dagger (D_\sigma H) - (D_\sigma H^\dagger) H] - \frac{ig_2}{F^2} [H^\dagger F_{\nu\rho}^W (D_\sigma H) - (D_\nu H^\dagger) F_{\rho\sigma}^W H]$$

The effective WZW term:

$$\mathcal{L}_{\text{WZW}} \supset -\frac{K\tilde{g}g_2^2 N_{\text{WZ}} v_{\text{SM}}^2}{48\sqrt{3}\pi^2 f^2} \epsilon^{\mu\nu\rho\sigma} B_{H\mu} \left[\sec^2 \theta_W Z_\sigma \partial_\nu Z_\rho + (D_\nu^A W_\rho^+) W_\sigma^- + (D_\nu^A W_\rho^-) W_\sigma^+ \right]$$

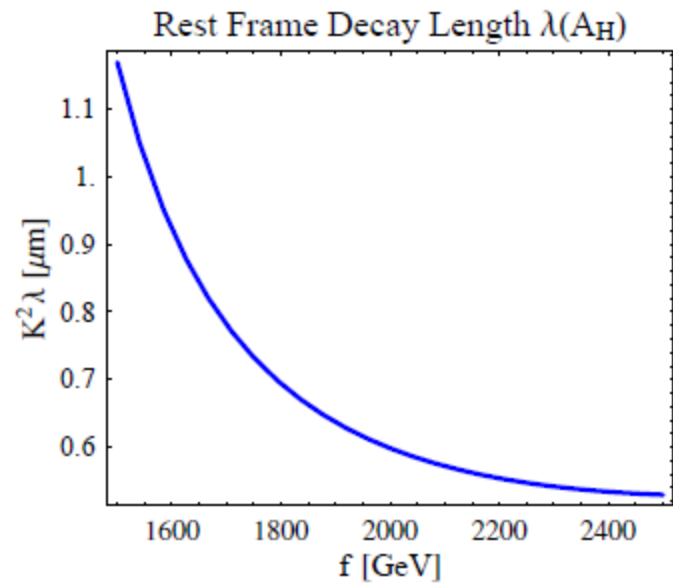
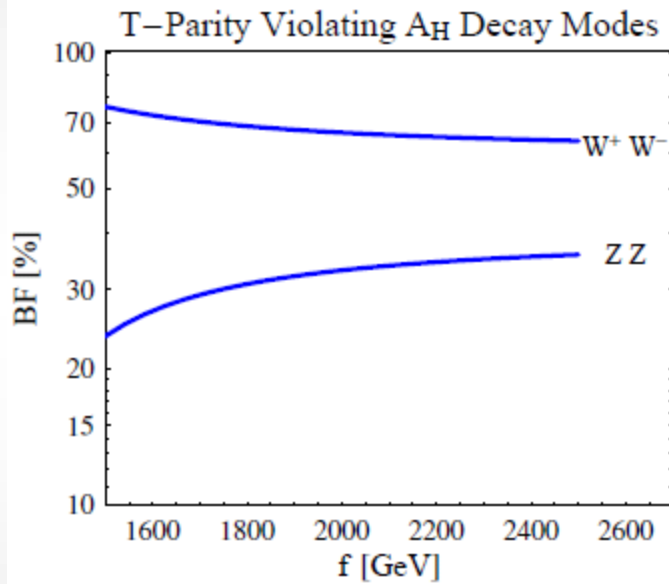
A_H is the dominant component of B_H

$$B_H = A_H \cos \theta_H + Z_H \sin \theta_H \qquad \sin \theta_H = \frac{5gg'}{4(5g^2 - g'^2)} \frac{v_{\text{SM}}^2}{f^2}$$

The A_H decay widths:

$$\Gamma(B_H \rightarrow ZZ) = \frac{1}{2\pi} \left(\frac{K\tilde{g}^3 N_{\text{WZ}}}{144\pi^2} \right)^2 \frac{m_Z^2}{m_{B_H}} \left(1 - \frac{4m_Z^2}{m_{B_H}^2} \right)^{\frac{5}{2}}$$

$$\Gamma(B_H \rightarrow W^+W^-) = \frac{1}{\pi} \left(\frac{K\tilde{g}^3 N_{\text{WZ}}}{144\pi^2} \right)^2 \frac{m_W^2}{m_{B_H}} \left(1 - \frac{4m_W^2}{m_{B_H}^2} \right)^{\frac{5}{2}}$$



Implications

- A_H is unstable. (decay instantaneous if kinematically allowed)

No massive missing particle at collider from LHT model.

A_H decays into WW or ZZ

multiple (4 and more) W channels

