

Isajet-M and Unified analysis of SUSY GUTs

Analyzing flavor & neutrino sectors with full running

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Outline

- 1 Analyzing Field Theories at Multiple Scales
- 2 Isajet-M Architecture
- 3 Results
- 4 Future work

We want to analyze SUSY GUTs in conjunction with low-scale data.

- We know masses, couplings, etc. at a low scale.
- We know coupling constants unify at some high scale.
- Generally, we have some ansatz for Yukawas, SSB parameters, etc. at the high scale.
- They are connected by the RGEs.

So... we have a set of ODEs with mixed boundary conditions.

Approximations for the RGEs

RGEs are commonly solved in . . .

- 1 The *dominant third-family approximation*.
 - This assumes $Y_{ij} \approx \delta_{i3}\delta_{j3}$.
 - Captures largest effect in quark sector.
- 2 The *diagonal approximation*.
 - Assumes $Y_{ij} \approx \text{diag}(y_{11}, y_{22}, y_{33})$
 - Better: the quark sector is approximately diagonal
 - But the lepton (neutrino) sector is not
- 3 Full three-family matrix couplings.
 - Treats all Yukawas in full generality

Difficulties with Matrix RGEs

- The number of parameters is greatly multiplied:
 - 3 Yukawa matrices eg y_{Uij}
 - 3 trilinear coupling matrices h_{Uij}
 - 5 mass-squared matrices $m_{\bar{U}}^2, m_{\bar{Q}}^2$, etc
 - 1 Yukawa matrix $y_{\nu ij}$
 - 1 trilinear coupling matrix $h_{\nu ij}$
 - 1 right-handed neutrino mass matrix $M_{\nu r ij}$
 - 1 mass-squared matrix $m_{\bar{\nu}}^2$
 - 1 effective coupling matrix κ_{ij}

A total of $16 \times 9 = 144$ matrix parameters.

- Basis issues

Current Codes

Several codes exist to solve the SUSY β -functions:

- SuSpect
- SoftSUSY
- SPheno
- Isajet
- REAP
- Isajet-M

Only SoftSUSY handles full three-family couplings.

Only REAP handles the neutrino sector properly.

We want a code that does *all* of this.

What can we do with Isajet-M?

- 1 Flavor analysis
 - CKM constraints
 - FCNC constraints, $b \rightarrow s \gamma$
- 2 Neutrino sector
 - Neutrino mass splittings, mixings
 - Neutrino phases
- 3 Leptogenesis
- 4 Effects of neutrino RGEs on e.g. $\Omega_{\tilde{Z}_1}$ (cf. Barger, Marfatia, Mustafayev 2008)

Example Application

Bimaximal mixing

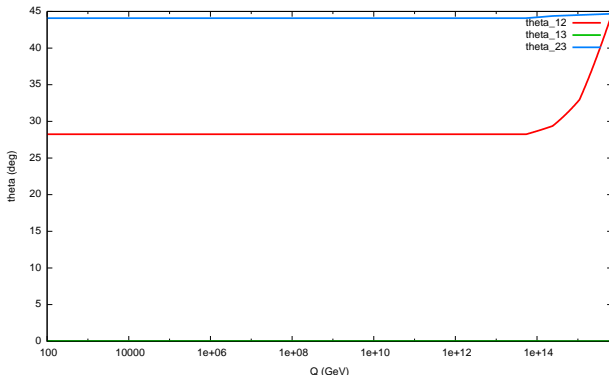


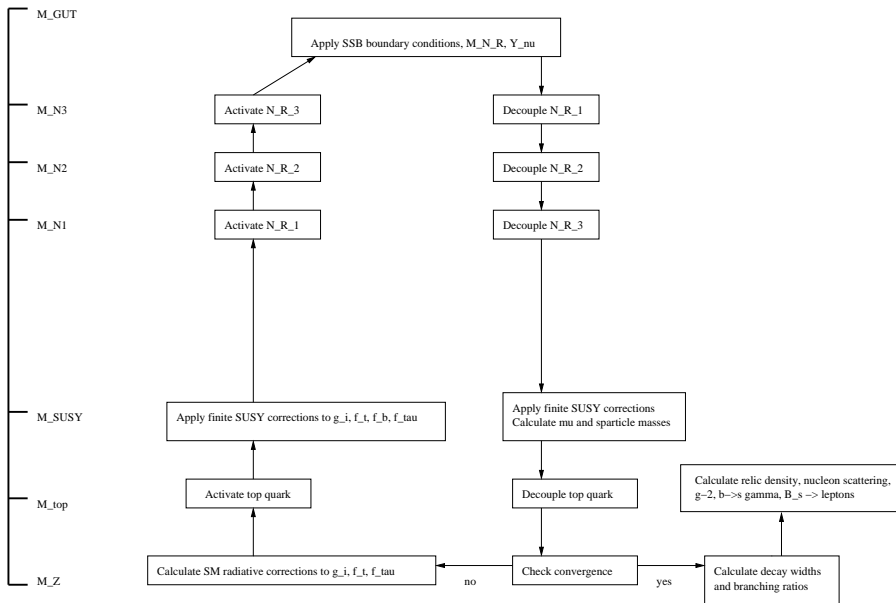
Figure: With Y_ν diagonal at the GUT scale, we see that θ_{12} runs significantly, while the other angles do not. Here we obtain bimaximal mixing at the GUT scale.

Isajet-M Layout

Overall scheme

- 1 Impose boundary conditions at M_Z
- 2 Run up to M_{GUT}
- 3 Impose GUT-scale boundary conditions
- 4 Run down to M_{SUSY}
- 5 Calculate $|\mu|$ and SUSY spectrum
- 6 Run back down to M_Z

Repeat until convergence.



Isajet-M Inputs

We take the following parameters as inputs:

- 1 SUSY model and parameters
 - mSUGRA: $m_0, m_{1/2}, A_0, \tan \beta, \text{sgn}(\mu)$
- 2 SM particle masses and mixings
- 3 Higgs vev
- 4 Right-handed neutrino Majorana mass matrix M_N
- 5 Neutrino Yukawa matrix Y_ν

Isajet-M Outputs

We calculate the following quantities:

- Sparticle spectrum
- Higgs masses
- Decay widths and branching fractions
- Neutrino masses and mass splittings
- Neutrino mixing angles (and phases)
- $\Omega_{\tilde{Z}_1} h^2$, $\text{BF}(b \rightarrow s \gamma)$, $(g - 2)_\mu$, $\sigma(\tilde{Z}_1 p)$, $\text{BF}(B_s \rightarrow \ell \bar{\ell})$.

Low-scale boundary conditions

Boundary conditions at M_Z :

- Use measured masses, v , and V_{CKM} to find Y_e, Y_u, Y_d
- We run in \overline{DR} , so first:
 - Run masses to M_Z in 3-loop QCD plus 2-loop QED \overline{MS} .
 - Apply finite threshold corrections
 - Convert from \overline{MS} into \overline{DR}
- Run from M_Z to m_{top} .
- Correct Y_{top} at m_{top} .

High-scale boundary conditions

Boundary conditions at M_{GUT} :

- M_{GUT} , determined by gauge unification
- Soft SUSY-breaking parameters:
 - Universal gaugino, scalar masses
 - Universal trilinear couplings
- M_{ν_r}, Y_ν

Decoupling

- Particles only contribute to RG evolution when they are kinematically accessible.
- At scales below their mass, they do not contribute.
- We change RGEs at mass thresholds:
 - 1 Sparticles
 - 2 Higgs, top quark
 - 3 Right-handed neutrinos
- Isajet includes finite threshold corrections for sparticles (but not right-handed neutrinos.)
- But these threshold corrections are all applied at M_{SUSY}

EFTs and the κ Matrix

- Seesaw formula not valid below seesaw scale

$$m_\nu = \frac{v^2}{2} Y_\nu^\dagger M_{N_R}^{-1} Y_\nu$$

- Right-handed neutrinos are decoupled and do not appear
- But evolution of ν_l is significant
- Use an EFT
- After integrating out each ν_r , replace it with an effective mass operator κ

EFTs: Diagram

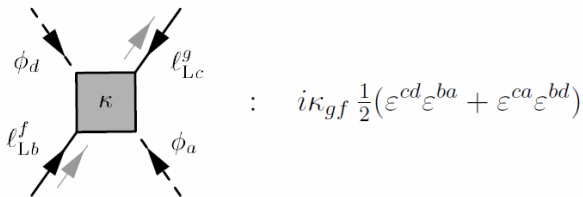


Figure: The effective operator κ and its Feynman rule

Diagram from Antusch et al.

Decoupling Procedure

Neutrino Sector

- 1 Run down from M_{GUT} to largest M_{ν_R} eigenvalue.
- 2 (n.b. this eigenvalue runs.)
- 3 Rotate into mass eigenbasis (diagonalizing M_{ν_R}).
- 4 Zero this row/col of M_{ν_R} and this *row* of Y_ν .
- 5 Create the appropriate entries in κ .

Neutrino sector comparisons

- We compare Isajet-M to REAP.
- REAP is not a full-fledged code:
 - Doesn't search for unification
 - All parameters must be specified at the high scale
- So, run Isajet-M on a given input
- Extract relevant parameters at M_{GUT}
- Enter them into REAP, compare results

Neutrino comparisons

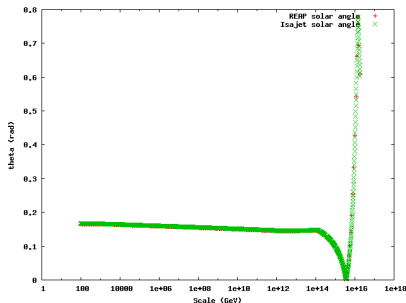


Figure: SPS 1b point, solar angle

$$Y_\nu(M_{\text{GUT}}) = \text{diag}(0.02, 0.1, 1.0)$$

$$M_N(M_{\text{GUT}}) = \begin{pmatrix} -0.5554E+11 & -0.2205E+10 & -0.3448E+11 \\ -0.2205E+10 & -0.1433E+13 & -0.6204E+12 \\ -0.3448E+11 & -0.6204E+12 & -0.1461E+15 \end{pmatrix}$$

Neutrino comparisons

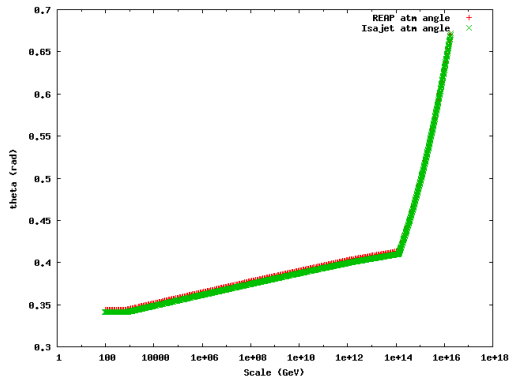


Figure: SPS 1b point, atmospheric angle

Neutrino comparisons

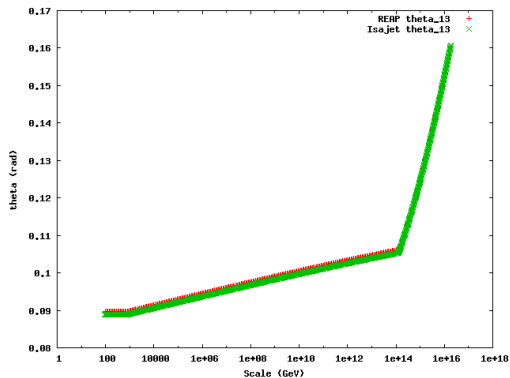


Figure: SPS 1b point, θ_{13}

Neutrino comparisons

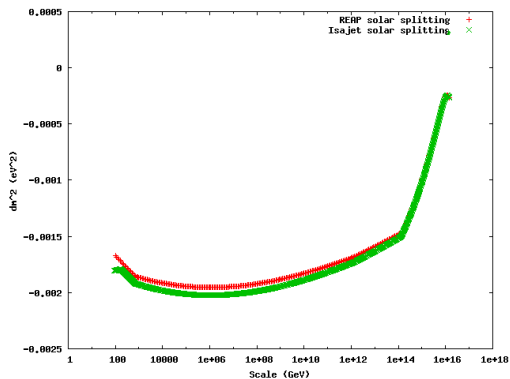


Figure: SPS 1b point, $\delta m^2_{\text{solar}}$

Neutrino comparisons

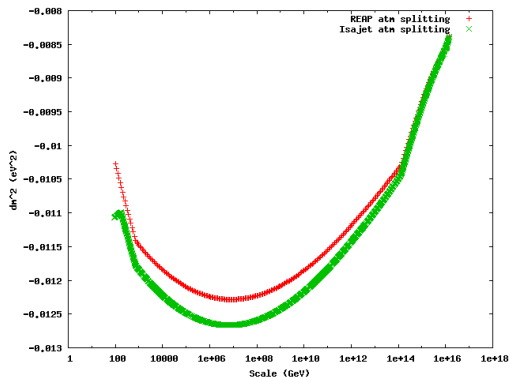


Figure: SPS 1b point, δm^2_{atm}

Neutrino comparisons

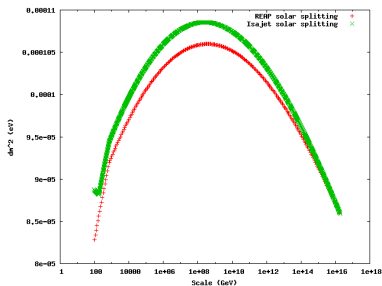


Figure: $\delta m^2_{\text{solar}}$

SPS1b point with $Y_\nu(M_{\text{GUT}}) = \text{diag}(0.001, 0.01, 0.1)$

$$M_N(M_{\text{GUT}}) = \begin{pmatrix} 0.4070E + 12 & 0.2320E + 13 & -0.1990E + 14 \\ 0.2320E + 13 & -0.1350E + 14 & 0.1160E + 15 \\ -0.1990E + 14 & 0.1160E + 15 & -0.1000E + 16 \end{pmatrix}$$

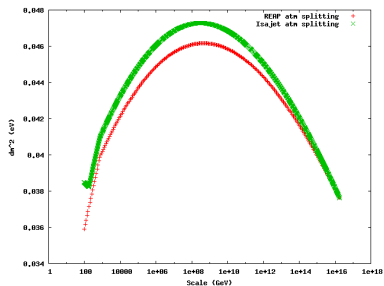


Figure: δm^2_{atm}

Neutrino comparisons

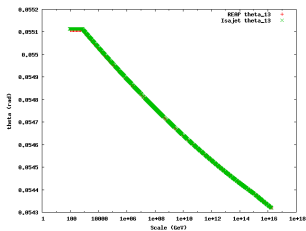


Figure: θ_{13}

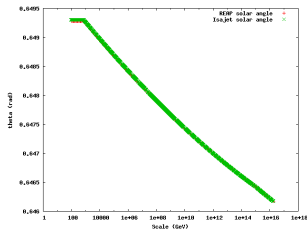


Figure: $\delta m^2_{\text{solar}}$

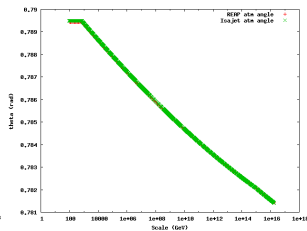


Figure: δm^2_{atm}

Neutrino mass/mixing comparison

SPS1a:

Code	θ_{12}	θ_{13}	θ_{23}	m_{ν_1} (eV)	m_{ν_2} (eV)	m_{ν_3} (eV)
SPS1a						
Isajet-M	0.142	0.105	0.41	0.219	0.216	0.192
REAP	0.1406	0.1058	0.4136	0.2114	0.2081	0.185
SPS1b						
Isajet-M	0.166	0.089	0.341	0.219	0.215	0.187
REAP	0.1654	0.08962	0.3444	0.2109	0.2069	0.1806
SPS2						
Isajet-M	0.142	0.105	0.411	0.213	0.21	0.186
REAP	0.1403	0.106	0.4144	0.2045	0.2014	0.179
SPS3						
Isajet-M	0.14	0.106	0.416	0.218	0.214	0.191
REAP	0.1382	0.1071	0.4193	0.2097	0.2066	0.1838
SPS4						
Isajet-M	0.195	0.056	0.209	0.214	0.209	0.169
REAP	0.1945	0.05647	0.2108	0.2062	0.2009	0.1629
SPS5						
Isajet-M	0.138	0.107	0.419	0.21	0.207	0.184
REAP	0.1366	0.1079	0.4227	0.2034	0.2004	0.1784

Comparison among codes

- At worst, we have a 5% difference from REAP.
- Similar differences exist for sparticle spectra among production codes:

Sparticle	Isajet	Isajet-M	Softsusy	Spheno	Suspect	(max-min)	diff[%]
χ_1^0	63.58	63.65	63.15	63.43	63.45	0.43	0.68
χ_2^0	116.63	116.4	116.92	117.00	117.09	0.46	0.39
χ_3^0	284.42	284.8	290.97	290.10	288.51	6.55	2.27
χ_4^0	305.78	306.5	307.15	306.56	307.57	1.79	0.58
$\chi_{1\pm}^\pm$	116.08	115.9	116.61	116.16	116.11	0.53	0.45
$\chi_{2\pm}^\pm$	304.50	305.3	305.92	308.77	308.55	4.27	1.39
\tilde{u}_L	397.88	397.94	394.65	398.50	397.19	3.84	0.97
\tilde{u}_R	388.28	388.38	387.80	388.56	387.43	1.13	0.29
$\tilde{\tau}_1$	95.57	94.14	95.35	95.30	94.55	1.02	1.07
$\tilde{\tau}_2$	154.54	151.31	155.18	154.98	154.59	3.23	2.1

$m_0 = 80$ GeV, $m_{1/2} = 170$ GeV, $A_0 = -250$ GeV, $\tan \beta = 10$, $\mu > 0$

Data from Kraml et al. automated comparison tool at <http://cern.ch/kraml/comparison/>.

Complex phases

- Complex phases in Y_ν can have enormous effect on the neutrino sector:

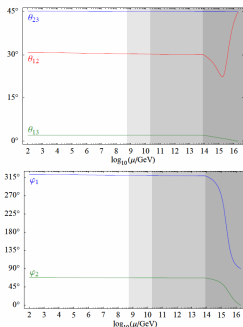


Figure 8: Highly non-linear running of θ_{12} and the Majorana phases in an example with large imaginary entries in the neutrino Yukawa matrix (see Eq. (29)). We used the MSSM with $\tan\beta = 10$, $M_{\text{GUT}} = 1 \text{ TeV}$ and the following initial conditions at $M_{\text{GUT}} = 2 \cdot 10^{16} \text{ GeV}$: $\theta_{12} = \theta_{21} = \pi/4$, $\theta_{13} = 0$, $\varphi_1 = \pi/2$, $\varphi_2 = 0$, normal hierarchy, $m_1 = 0.08 \text{ eV}$, $\Delta m_{21}^2 = 1.1 \cdot 10^{-4} \text{ eV}^2$, $\Delta m_{31}^2 = 4 \cdot 10^{-3} \text{ eV}^2$.

- They are unconstrained
- We are now complexifying Isajet-M

Searches in parameter space

- Goal: find models that match all observable data
- Simple ansatz at M_{GUT} for Y 's, M_N
- Choose grid on parameter space, test points

Summary

- To analyze GUTs, we have to solve the β -functions with mixed boundary conditions.
- To understand flavor physics and neutrino mixings, we need to evolve the couplings in **full matrix form**.

Future outlook

- Analyze parameter space for various SUSY GUTs, looking for viable points.
- Complexify and add analysis of baryogenesis/leptogenesis constraints.



R Dermisek and S Raby.

Bi-large neutrino mixing and CP violation in an SO(10) SUSY GUT for fermion masses.

Phys Lett B, 2005.

[hep-ph/0507045](#)



S Antusch, J Kersten, M Lindner, M Ratz, MA Schmidt.

Running neutrino mass parameters in see-saw scenarios.

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D Pierce, J Bagger, K Matchev, R Zhang.

Precision corrections in the minimal supersymmetric standard model.

Nucl Phys B, 1997.

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