

μ -problem and neutrino masses in supersymmetry

Raghavendra Srikanth Hundi

University of Hawaii

Work in progress with:

Sandip Pakvasa and Xerxes Tata

μ -problem

The superpotential of Minimal Supersymmetric Standard Model is

$$W_{MSSM} = \mu H_u H_d + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d$$

The superpotential of Minimal Supersymmetric Standard Model is

$$W_{MSSM} = \mu H_u H_d + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d$$

The soft terms in the MSSM are

$$V_{soft} = m_{H_d}^2 |H_d|^2 + m_{H_u}^2 |H_u|^2 + (B_\mu \mu) H_u H_d + \dots$$

The superpotential of Minimal Supersymmetric Standard Model is

$$W_{MSSM} = \mu H_u H_d + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d$$

The soft terms in the MSSM are

$$V_{soft} = m_{H_d}^2 |H_d|^2 + m_{H_u}^2 |H_u|^2 + (B_\mu \mu) H_u H_d + \dots$$

In order to have consistent electroweak symmetry breaking, we should have

$$\mu^2 \sim m_{H_d}^2, m_{H_u}^2, (B_\mu \mu) \sim \text{TeV}^2$$

The superpotential of Minimal Supersymmetric Standard Model is

$$W_{MSSM} = \mu H_u H_d + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d$$

The soft terms in the MSSM are

$$V_{soft} = m_{H_d}^2 |H_d|^2 + m_{H_u}^2 |H_u|^2 + (B_{\mu\mu}) H_u H_d + \dots$$

In order to have consistent electroweak symmetry breaking, we should have

$$\mu^2 \sim m_{H_d}^2, m_{H_u}^2, (B_{\mu\mu}) \sim \text{TeV}^2$$

We should explain the origin of μ -term and soft terms in a single theory.

Soft terms and μ -term in supergravity

Soft terms and μ -term in supergravity

Locally supersymmetric invariant theory is called supergravity.

Soft terms and μ -term in supergravity

Locally supersymmetric invariant theory is called supergravity.

In supergravity we can have fields which are gauge singlets but interact gravitationally.

These are called hidden sector fields.

Soft terms and μ -term in supergravity

Locally supersymmetric invariant theory is called supergravity.

In supergravity we can have fields which are gauge singlets but interact gravitationally.

These are called hidden sector fields.

In supergravity we can write superpotential and Kahler potential as

$$W = \Lambda^2 S + W_{MSSM}, \quad K = S^\dagger S + \sum \Phi_i^\dagger \Phi_i$$

where S is a hidden sector field, Φ_i is any MSSM field and $\Lambda \sim 10^{10}$ GeV.

Soft terms and μ -term in supergravity

Locally supersymmetric invariant theory is called supergravity.

In supergravity we can have fields which are gauge singlets but interact gravitationally.

These are called hidden sector fields.

In supergravity we can write superpotential and Kahler potential as

$$W = \Lambda^2 S + W_{MSSM}, \quad K = S^\dagger S + \sum \Phi_i^\dagger \Phi_i$$

where S is a hidden sector field, Φ_i is any MSSM field and $\Lambda \sim 10^{10}$ GeV.

The scalar potential in supergravity contains all soft terms with corresponding parameters as

$$m_{soft} \sim \frac{\Lambda^2 \langle S \rangle}{M_P^2} \sim \text{TeV}, \quad \text{if } \langle S \rangle \sim M_P.$$

Soft terms and μ -term in supergravity

Locally supersymmetric invariant theory is called supergravity.

In supergravity we can have fields which are gauge singlets but interact gravitationally.

These are called hidden sector fields.

In supergravity we can write superpotential and Kahler potential as

$$W = \Lambda^2 S + W_{MSSM}, \quad K = S^\dagger S + \sum \Phi_i^\dagger \Phi_i$$

where S is a hidden sector field, Φ_i is any MSSM field and $\Lambda \sim 10^{10}$ GeV.

The scalar potential in supergravity contains all soft terms with corresponding parameters as

$$m_{soft} \sim \frac{\Lambda^2 \langle S \rangle}{M_P^2} \sim \text{TeV}, \quad \text{if } \langle S \rangle \sim M_P.$$

To explain the μ -term, let us assume some symmetry group G . Let us assume G forbids the usual $H_u H_d$ term, but allows a term

$$W = \frac{X_1^2}{M_P} H_u H_d,$$

where X_1 is some other hidden sector field. If $\langle X_1 \rangle \sim \Lambda$, the above term gives a μ -parameter of order TeV.

To explain the origin of soft terms and μ -term, we need at least two hidden sector fields with

$$\langle S \rangle \sim M_P, \quad \langle X_1 \rangle \sim \Lambda$$

To explain the origin of soft terms and μ -term, we need at least two hidden sector fields with

$$\langle S \rangle \sim M_P, \quad \langle X_1 \rangle \sim \Lambda$$

To achieve the required scalar vacuum expectation values (vev's), we may have to choose non-minimal Kahler potential.

$$\begin{aligned} W &= \Lambda^2 S + \dots, \\ K &= S^\dagger S + X_1^\dagger X_1 + \left(a \frac{S}{M_P} S^\dagger S + b \frac{S}{M_P} X_1^\dagger X_1 + \dots \right), \end{aligned}$$

where $a, b \sim \mathcal{O}(1)$ constants.

To explain the origin of soft terms and μ -term, we need at least two hidden sector fields with

$$\langle S \rangle \sim M_P, \quad \langle X_1 \rangle \sim \Lambda$$

To achieve the required scalar vacuum expectation values (vev's), we may have to choose non-minimal Kahler potential.

$$\begin{aligned} W &= \Lambda^2 S + \dots, \\ K &= S^\dagger S + X_1^\dagger X_1 + \left(a \frac{S}{M_P} S^\dagger S + b \frac{S}{M_P} X_1^\dagger X_1 + \dots \right), \end{aligned}$$

where $a, b \sim \mathcal{O}(1)$ constants.

We have seen that we need a symmetry group G to forbid the $H_u H_d$ term in the superpotential.

To explain the origin of soft terms and μ -term, we need at least two hidden sector fields with

$$\langle S \rangle \sim M_P, \quad \langle X_1 \rangle \sim \Lambda$$

To achieve the required scalar vacuum expectation values (vev's), we may have to choose non-minimal Kahler potential.

$$\begin{aligned} W &= \Lambda^2 S + \dots, \\ K &= S^\dagger S + X_1^\dagger X_1 + \left(a \frac{S}{M_P} S^\dagger S + b \frac{S}{M_P} X_1^\dagger X_1 + \dots \right), \end{aligned}$$

where $a, b \sim \mathcal{O}(1)$ constants.

We have seen that we need a symmetry group G to forbid the $H_u H_d$ term in the superpotential.

Since gravitational interactions do not respect global symmetries, we can choose a gauge symmetry. The minimal choice for G is a U(1) gauge symmetry.



Neutrino masses

Neutrino masses

The neutrino oscillation data indicates the existence of three flavor neutrinos, with masses $m_\nu \leq .1$ eV.

Neutrino masses

The neutrino oscillation data indicates the existence of three flavor neutrinos, with masses $m_\nu \leq .1$ eV.

Neutrino mass is so small compared to other elementary particle masses.

Neutrino masses

The neutrino oscillation data indicates the existence of three flavor neutrinos, with masses $m_\nu \leq .1$ eV.

Neutrino mass is so small compared to other elementary particle masses.

Let us try to understand the neutrino masses with the known MSSM fields and hidden sector fields.

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Let us suppose that bilinear term (LH_u) exists.

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Let us suppose that bilinear term (LH_u) exists.

Neutrinos will have mixing masses with gauginos and Higgsinos.

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Let us suppose that bilinear term (LH_u) exists.

Neutrinos will have mixing masses with gauginos and Higgsinos.

The mixing mass is $\sim \epsilon$.

$$W_{LV} \sim \lambda L L E^c + \lambda' Q L D^c + \epsilon L H_u$$

Let us suppose that bilinear term (LH_u) exists.

Neutrinos will have mixing masses with gauginos and Higgsinos.

The mixing mass is $\sim \epsilon$.

The mass scale of gauginos and Higgsinos is $M \sim \mathcal{O}(100)$ GeV.

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Let us suppose that bilinear term (LH_u) exists.

Neutrinos will have mixing masses with gauginos and Higgsinos.

The mixing mass is $\sim \epsilon$.

The mass scale of gauginos and Higgsinos is $M \sim \mathcal{O}(100)$ GeV.

The mass of light neutrinos is $m_\nu \sim \epsilon^2/M$.

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Let us suppose that bilinear term (LH_u) exists.

Neutrinos will have mixing masses with gauginos and Higgsinos.

The mixing mass is $\sim \epsilon$.

The mass scale of gauginos and Higgsinos is $M \sim \mathcal{O}(100)$ GeV.

The mass of light neutrinos is $m_\nu \sim \epsilon^2/M$.

For $\epsilon \sim 10^{-4}$ GeV, neutrino mass of .1 eV can be explained.

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Let us suppose that bilinear term (LH_u) exists.

Neutrinos will have mixing masses with gauginos and Higgsinos.

The mixing mass is $\sim \epsilon$.

The mass scale of gauginos and Higgsinos is $M \sim \mathcal{O}(100)$ GeV.

The mass of light neutrinos is $m_\nu \sim \epsilon^2/M$.

For $\epsilon \sim 10^{-4}$ GeV, neutrino mass of .1 eV can be explained.

Let us suppose $U(1)'$ forbids LH_u term, instead, let us have

$$W = \frac{X_2^3}{M_P^2} LH_u.$$

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Let us suppose that bilinear term (LH_u) exists.

Neutrinos will have mixing masses with gauginos and Higgsinos.

The mixing mass is $\sim \epsilon$.

The mass scale of gauginos and Higgsinos is $M \sim \mathcal{O}(100)$ GeV.

The mass of light neutrinos is $m_\nu \sim \epsilon^2/M$.

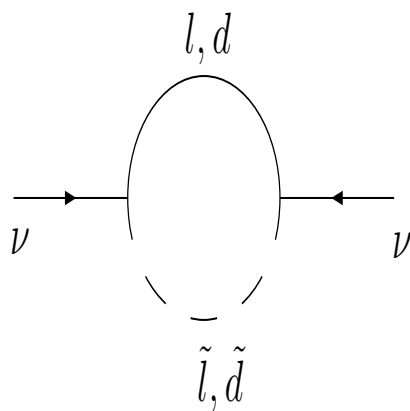
For $\epsilon \sim 10^{-4}$ GeV, neutrino mass of .1 eV can be explained.

Let us suppose $U(1)'$ forbids LH_u term, instead, let us have

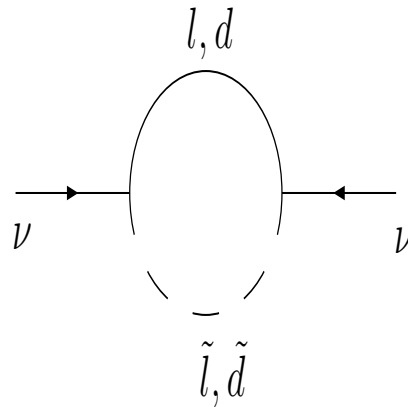
$$W = \frac{X_2^3}{M_P^2} LH_u.$$

For $\langle X_2 \rangle \sim \Lambda$, we get $\frac{\langle X_2 \rangle^3}{M_P^2} \sim 10^{-4}$ GeV.

The remaining two terms λLLE^c and $\lambda' QlD^c$, generate neutrino masses at one-loop level.

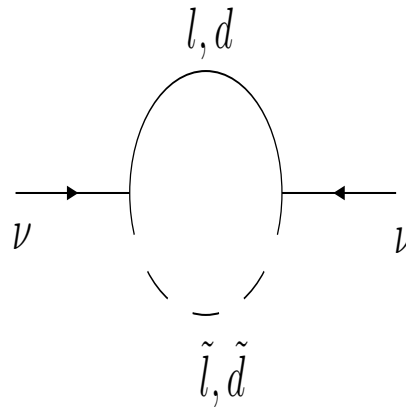


The remaining two terms λLLE^c and $\lambda' QlD^c$, generate neutrino masses at one-loop level.



$$m_\nu \sim \frac{\lambda^2}{32\pi^2} m_\tau + 3 \frac{\lambda'^2}{32\pi^2} m_b$$

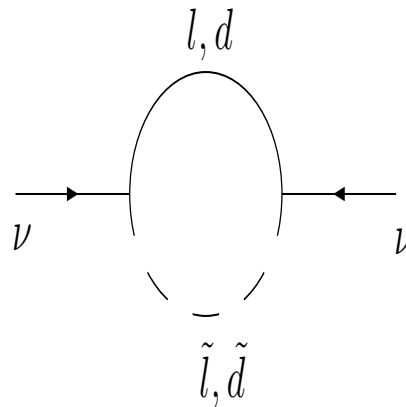
The remaining two terms λLLE^c and $\lambda' QlD^c$, generate neutrino masses at one-loop level.



$$m_\nu \sim \frac{\lambda^2}{32\pi^2} m_\tau + 3 \frac{\lambda'^2}{32\pi^2} m_b$$

We need $\lambda, \lambda' \sim 10^{-4}$ in order to explain the smallness of neutrino masses.

The remaining two terms λLLE^c and $\lambda' QlD^c$, generate neutrino masses at one-loop level.

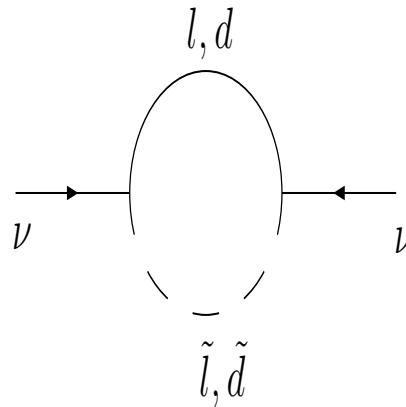


$$m_\nu \sim \frac{\lambda^2}{32\pi^2} m_\tau + 3 \frac{\lambda'^2}{32\pi^2} m_b$$

We need $\lambda, \lambda' \sim 10^{-4}$ in order to explain the smallness of neutrino masses. In our supergravity set-up, we can forbid LLE^c , QLD^c , and let us have

$$W \sim \frac{Y_1}{M_P} LLE^c + \frac{Y_2}{M_P} QLD^c.$$

The remaining two terms λLLE^c and $\lambda' QlD^c$, generate neutrino masses at one-loop level.



$$m_\nu \sim \frac{\lambda^2}{32\pi^2} m_\tau + 3 \frac{\lambda'^2}{32\pi^2} m_b$$

We need $\lambda, \lambda' \sim 10^{-4}$ in order to explain the smallness of neutrino masses. In our supergravity set-up, we can forbid LLE^c , QLD^c , and let us have

$$W \sim \frac{Y_1}{M_P} LLE^c + \frac{Y_2}{M_P} QLD^c.$$

We have $\frac{\langle Y_{1,2} \rangle}{M_P} \sim \frac{\Lambda}{M_P} \sim 10^{-7}$.

To summarize, we have a gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$.

To summarize, we have a gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$.
We can write

$$W = \Lambda^2 S + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d \\ + \frac{X_1^2}{M_P} H_u H_d + \frac{X_2^3}{M_P^2} h_i L_i H_u$$

To summarize, we have a gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$.
We can write

$$W = \Lambda^2 S + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d \\ + \frac{X_1^2}{M_P} H_u H_d + \frac{X_2^3}{M_P^2} h_i L_i H_u$$

The hidden sector fields should have vevs:

$$\langle S \rangle \sim M_P, \quad \langle X_1 \rangle \sim \Lambda, \quad \langle X_2 \rangle \sim \Lambda$$

To summarize, we have a gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$.
We can write

$$W = \Lambda^2 S + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d \\ + \frac{X_1^2}{M_P} H_u H_d + \frac{X_2^3}{M_P^2} h_i L_i H_u$$

The hidden sector fields should have vevs:

$$\langle S \rangle \sim M_P, \quad \langle X_1 \rangle \sim \Lambda, \quad \langle X_2 \rangle \sim \Lambda$$

$$K = S^\dagger S + X_1^\dagger X_1 + X_2^\dagger X_2 + \Phi_i^\dagger \Phi_i \\ + \left(a \frac{S}{M_P} S^\dagger S + b \frac{S}{M_P} X_1^\dagger X_1 + c \frac{S}{M_P} X_2^\dagger X_2 + d \frac{S}{M_P} \Phi_i^\dagger \Phi_i + \text{h.c.} \right)$$

To summarize, we have a gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$.
We can write

$$W = \Lambda^2 S + Y_{ij}^u Q_i U_j^c H_u + Y_{ij}^d Q_i D_j^c H_d + Y_{ij}^e L_i E_j^c H_d \\ + \frac{X_1^2}{M_P} H_u H_d + \frac{X_2^3}{M_P^2} h_i L_i H_u$$

The hidden sector fields should have vevs:

$$\langle S \rangle \sim M_P, \quad \langle X_1 \rangle \sim \Lambda, \quad \langle X_2 \rangle \sim \Lambda$$

$$K = S^\dagger S + X_1^\dagger X_1 + X_2^\dagger X_2 + \Phi_i^\dagger \Phi_i \\ + \left(a \frac{S}{M_P} S^\dagger S + b \frac{S}{M_P} X_1^\dagger X_1 + c \frac{S}{M_P} X_2^\dagger X_2 + d \frac{S}{M_P} \Phi_i^\dagger \Phi_i + \text{h.c.} \right)$$

$$F_{ab} = \frac{\delta_{ab}}{g_a^2} \left(1 + f \frac{S}{M_P} \right)$$

Stability of proton

Stability of proton

In our present model, we have effective operators:

$$W_{LV} \sim \lambda LLE^c + \lambda' QLD^c + \epsilon LH_u$$

Stability of proton

In our present model, we have effective operators:

$$W_{LV} \sim \lambda L L E^c + \lambda' Q L D^c + \epsilon L H_u$$

In addition to the above operators, if we have baryon number violating operators

$$W_{BV} \sim \lambda'' U^c D^c D^c,$$

proton can decay to some lighter particles.

Stability of proton

In our present model, we have effective operators:

$$W_{LV} \sim \lambda L L E^c + \lambda' Q L D^c + \epsilon L H_u$$

In addition to the above operators, if we have baryon number violating operators

$$W_{BV} \sim \lambda'' U^c D^c D^c,$$

proton can decay to some lighter particles.

Life time of proton is $\tau_P \geq 10^{32}$ Years.

Stability of proton

In our present model, we have effective operators:

$$W_{LV} \sim \lambda L L E^c + \lambda' Q L D^c + \epsilon L H_u$$

In addition to the above operators, if we have baryon number violating operators

$$W_{BV} \sim \lambda'' U^c D^c D^c,$$

proton can decay to some lighter particles.

Life time of proton is $\tau_P \geq 10^{32}$ Years.

In a supersymmetric model, we can also have dimension-5 operators:

$$W_5 \sim \frac{a_1}{M_P} Q Q Q L + \frac{a_2}{M_P} U^c U^c D^c E^c$$

Stability of proton

In our present model, we have effective operators:

$$W_{LV} \sim \lambda L L E^c + \lambda' Q L D^c + \epsilon L H_u$$

In addition to the above operators, if we have baryon number violating operators

$$W_{BV} \sim \lambda'' U^c D^c D^c,$$

proton can decay to some lighter particles.

Life time of proton is $\tau_P \geq 10^{32}$ Years.

In a supersymmetric model, we can also have dimension-5 operators:

$$W_5 \sim \frac{a_1}{M_P} Q Q Q L + \frac{a_2}{M_P} U^c U^c D^c E^c$$

In our specific model with the help of additional $U(1)'$, we can forbid W_{BV} and W_5 .

Anomalies

In our present model we need to satisfy the following anomalies:

$$[SU(3)_c]^2 - U(1)', \quad [SU(2)_L]^2 - U(1)', \quad [U(1)_Y]^2 - U(1)'$$
$$[\text{gravity}]^2 - U(1)', \quad U(1)_Y - [U(1)']^2, \quad [U(1)']^3$$

In our present model we need to satisfy the following anomalies:

$$[SU(3)_c]^2 - U(1)', \quad [SU(2)_L]^2 - U(1)', \quad [U(1)_Y]^2 - U(1)'$$
$$[\text{gravity}]^2 - U(1)', \quad U(1)_Y - [U(1)']^2, \quad [U(1)']^3$$

We have found that we can satisfy these anomalies at the cost of introducing some additional fields into the model.

In our present model we need to satisfy the following anomalies:

$$[SU(3)_c]^2 - U(1)', \quad [SU(2)_L]^2 - U(1)', \quad [U(1)_Y]^2 - U(1)'$$
$$[\text{gravity}]^2 - U(1)', \quad U(1)_Y - [U(1)']^2, \quad [U(1)']^3$$

We have found that we can satisfy these anomalies at the cost of introducing some additional fields into the model.

Specifically, we need vector-like triplets and some additional hidden sector fields.

Conclusions and future work

Conclusions and future work

Conclusions:

Conclusions and future work

Conclusions:

- We have attempted to solve μ -problem and neutrino mass problem in a supergravity set-up.

Conclusions and future work

Conclusions:

- We have attempted to solve μ -problem and neutrino mass problem in a supergravity set-up.
- To solve these problems we have introduced an additional gauge symmetry $U(1)'$.

Conclusions and future work

Conclusions:

- We have attempted to solve μ -problem and neutrino mass problem in a supergravity set-up.
- To solve these problems we have introduced an additional gauge symmetry $U(1)'$.
- The gauge symmetry $U(1)'$ can forbid the baryon number violating and dimension-5 operators, that would cause problems to the stability of proton.

Conclusions and future work

Conclusions:

- We have attempted to solve μ -problem and neutrino mass problem in a supergravity set-up.
- To solve these problems we have introduced an additional gauge symmetry $U(1)'$.
- The gauge symmetry $U(1)'$ can forbid the baryon number violating and dimension-5 operators, that would cause problems to the stability of proton.
- The anomalies with respect to the $U(1)'$ can also be satisfied.

Conclusions and future work

Conclusions:

- We have attempted to solve μ -problem and neutrino mass problem in a supergravity set-up.
- To solve these problems we have introduced an additional gauge symmetry $U(1)'$.
- The gauge symmetry $U(1)'$ can forbid the baryon number violating and dimension-5 operators, that would cause problems to the stability of proton.
- The anomalies with respect to the $U(1)'$ can also be satisfied.

Future work:

Conclusions and future work

Conclusions:

- We have attempted to solve μ -problem and neutrino mass problem in a supergravity set-up.
- To solve these problems we have introduced an additional gauge symmetry $U(1)'$.
- The gauge symmetry $U(1)'$ can forbid the baryon number violating and dimension-5 operators, that would cause problems to the stability of proton.
- The anomalies with respect to the $U(1)'$ can also be satisfied.

Future work:

- We are studying the phenomenological prospects of the gauge boson and gaugino corresponding to the $U(1)'$.

Conclusions and future work

Conclusions:

- We have attempted to solve μ -problem and neutrino mass problem in a supergravity set-up.
- To solve these problems we have introduced an additional gauge symmetry $U(1)'$.
- The gauge symmetry $U(1)'$ can forbid the baryon number violating and dimension-5 operators, that would cause problems to the stability of proton.
- The anomalies with respect to the $U(1)'$ can also be satisfied.

Future work:

- We are studying the phenomenological prospects of the gauge boson and gaugino corresponding to the $U(1)'$.
- We are studying the physical implications of the additional fields that have been introduced to cancel the anomalies.