

# Neutrino CP Violation and Baryon Asymmetry in the Universe

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# Motivation

- **Neutrino Oscillations**

- $|\Delta m_{atm}^2| = (2.1 - 2.7) \times 10^{-3} eV^2 (2\sigma),$   
 $|\Delta m_{sol}^2| = (7.3 - 8.1) \times 10^{-5} eV^2 (2\sigma)$

Standard model prediction,

$$M_\nu^2 \sim \left( \frac{1}{M_{Pl}} (v_u)^2 \right)^2 \sim (10^{-10} - 10^{-8}) eV^2 \text{ with}$$
$$M_{Pl} \sim 2 \cdot 10^{18} \text{ GeV}$$

- $\sin^2 \theta_{23} = 0.38 - 0.63 (2\sigma), \sin^2 \theta_{12} = 0.28 - 0.37 (2\sigma)$   
 $\sin^2 \theta_{13} \leq 0.050 (3\sigma), \delta \text{ no constraints.}$

**Schwartz hep-ph/0710.5027**

# Motivation

- **Baryogenesis**

- Sakharov's conditions

- Baryon number violation

- C and CP violation

- Nonequilibrium dynamics

- $\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma} = (6.15 \pm 0.25) \times 10^{-10},$

SM induced CP violation is too small,  $10^{-20}$ .

**Neutrino data and baryon asymmetry require new physics beyond SM. We wish to connect these two phenomena.**

# Framework

## - Seesaw Mechanism

$$(1) \quad L^{D+M} = -\frac{1}{2}\bar{N}MN^c - \bar{N}Y_\nu\nu H_u$$

Left neutrino masses are inversely proportional to heavy RH neutrino masses

$$(2) \quad M_\nu = v_u^2 Y_\nu^T M^{-1} Y_\nu$$

If the number of RH neutrinos  $l \leq 2$ ,  $\det(M_\nu) = 0$ , there must be at least one massless left handed neutrino. If  $m_1 = 0$ , normal hierarchy, if  $m_3 = 0$ , inverted hierarchy.

# Framework

## - Resonant Leptogenesis

For fulfilling of thermal baryogenesis, hierarchical RH neutrino  $\rightarrow M > 2 \times 10^9 \text{ GeV} \rightarrow T_{reheat} \geq 10^8 \text{ GeV} \rightarrow$  overproduction of gravitinos  $\rightarrow$  conflicts with observed light elements abundance.

Davidson and Ibarra hep-ph/0202239

Kohri *et. al.* hep-ph/0507245

Nearly degeneracy of RH neutrino masses enhances CP violation resonantly, resolves this problem.

# The Model

## Normal hierarchy

We assumed one texture 0 to realize the connection of neutrino and baryon asymmetry. The neutrino Yukawa coupling

$$Y_\nu = \begin{pmatrix} a_1 & a_2 & a_3 e^{i\phi} \\ 0 & b_2 & b_3 \end{pmatrix}, \quad M_N = M \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

(3)

Through seesaw mechanism,

$$M_\nu = \frac{v_u^2}{M} \begin{pmatrix} 0 & a_1 b_2 & a_1 b_3 \\ a_1 b_2 & 2a_2 b_2 & a_2 b_3 + a_3 b_2 e^{i\phi} \\ a_1 b_3 & a_2 b_3 + a_3 b_2 e^{i\phi} & 2a_3 b_3 e^{i\phi} \end{pmatrix}.$$

(4)

$$(5) \quad \tan \theta_{13} = \sin \theta_{12} \sqrt{\frac{m_2}{m_3}}.$$

$\sin^2 \theta_{13} = 0.044$ , within  $3\sigma$  accuracy.  $\theta_{12}$  and  $\theta_{23}$  within  $2\sigma$  accuracy.

# The Model

## Inverted hierarchy

The neutrino Yukawa coupling

$$(6) \quad Y_\nu = \begin{pmatrix} a_1 & a_2 e^{i\phi} & a_3 \\ b_1 & b_2 & 0 \end{pmatrix}$$

Through seesaw mechanism, the LH neutrino Majorana mass matrix is

$$M_\nu = \frac{v_u^2}{M} \begin{pmatrix} 2a_1 b_2 & a_1 b_2 + a_2 b_1 e^{i\phi} & a_3 b_1 \\ a_2 b_1 e^{i\phi} + a_1 b_2 & 2a_2 b_2 e^{i\phi} & a_3 b_2 \\ a_3 b_1 & a_3 b_2 & 0 \end{pmatrix}.$$

With inverted hierarchy, we made a prediction

$$(7) \quad \sin^2 \theta_{12} \simeq \frac{1}{2} - \frac{\sin \theta_{13} \tan \theta_{23} \cos \delta}{|\tan^2 \theta_{23} + \sin^2 \theta_{13} e^{2i\delta}|}$$

$$(8) \quad -56^\circ \leq \delta \leq 56^\circ$$

# Baryogenesis

Sakharov's conditions are satisfied in the expanding universe and the out of equilibrium CP violating decays of RH neutrinos. When two RH neutrino have nearly degenerate masses, the CP violation effect is enhanced resonantly by a factor  $\frac{(M_2^2 - M_1^2)M_1\Gamma_2}{(M_2^2 - M_1^2)^2 + M_1^2\Gamma_2^2}$ . The mass split is induced at one loop

$$(9) \quad 16\pi^2 \frac{d}{dt} M_N = 2M_N Y_\nu Y_\nu^\dagger + \text{t.c.}$$

and

$$(10) \quad M_N \rightarrow M \begin{pmatrix} \delta_N & 1 \\ 1 & \delta'_N \end{pmatrix} = M \begin{pmatrix} \delta_N & 1 \\ 1 & \delta_N^* \end{pmatrix},$$



# Baryogenesis

4 parameters of the Yukawa coupling constants of E.(3) can be determined with neutrino oscillation data. We defined

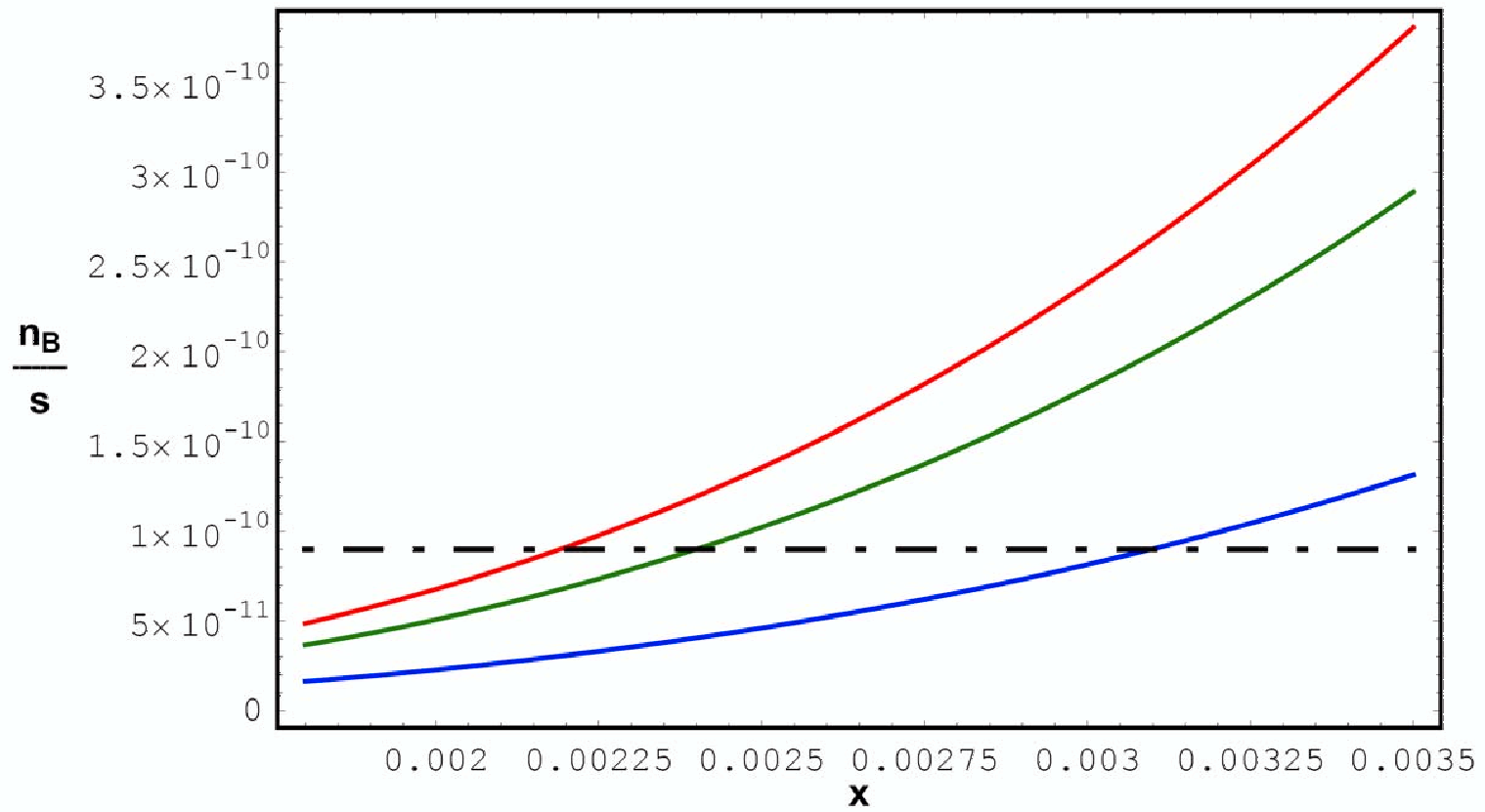
$$x = \frac{a_3}{b}.$$

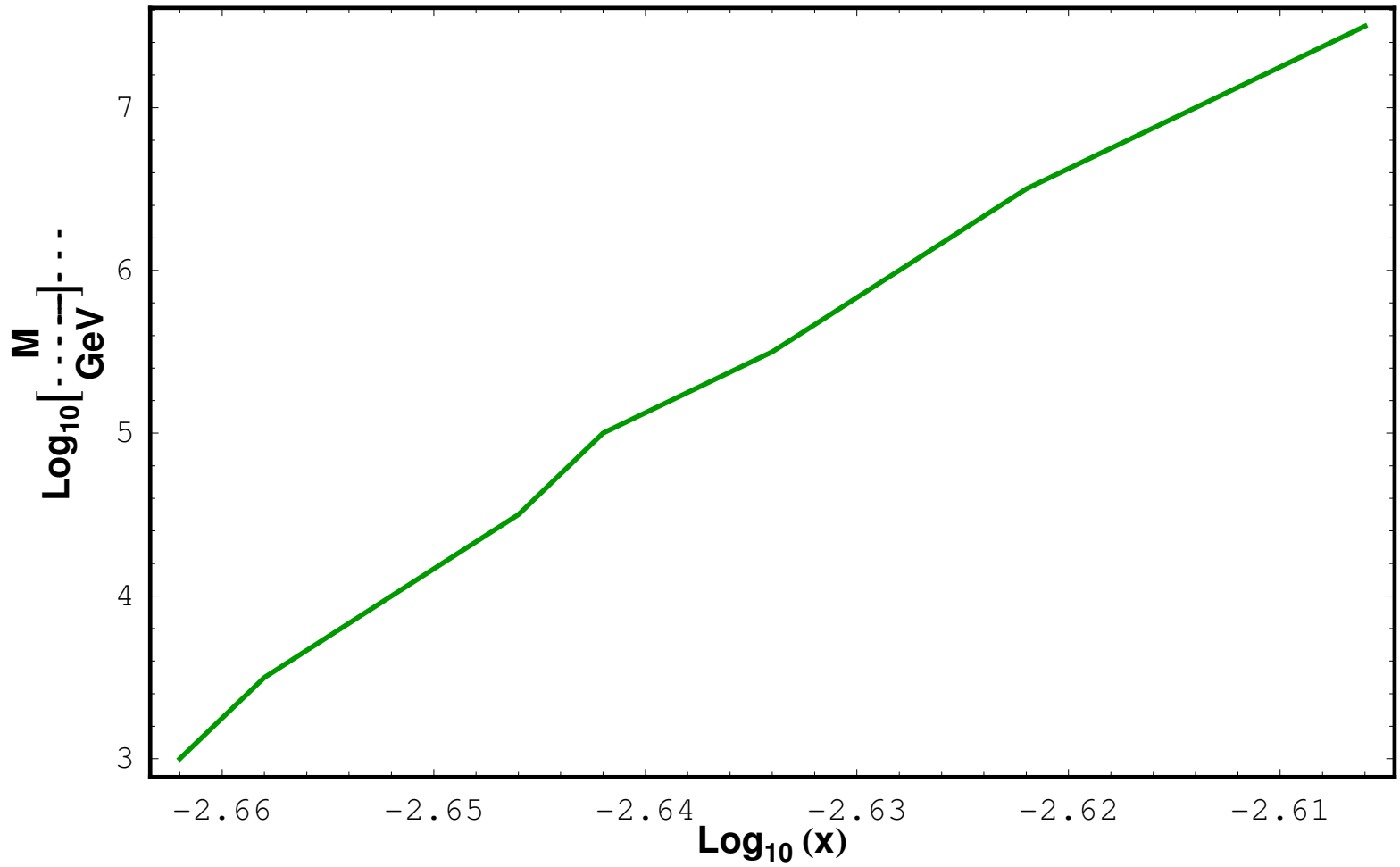
$\delta$  has no constraints. We only studied baryogenesis with normal hierarchy. Study with inverted hierarchy was done before.

*Babu and et.al hep-ph/0705.4419.*

We plot our results below.

$\frac{n_B}{s}$  vs.  $x$  with  $M = 10^6$  GeV. Cases:  $\delta = \pi/4, \pi/3, \pi/2$





$M$  vs.  $x$  with  $\frac{n_B}{s} = 9 \cdot 10^{-11}$  and  $\delta = \frac{\pi}{3}$ .

# Conclusions

- We investigated two new models of texture zeros, with the predictions:

(i) For normal hierarchy:

$$\tan \theta_{13} = \sin \theta_{12} \sqrt{\frac{m_2}{m_3}}$$

(ii) For inverted hierarchy:

$$\sin^2 \theta_{12} \simeq \frac{1}{2} - \frac{\sin \theta_{13} \tan \theta_{23} \cos \delta}{|\tan^2 \theta_{23} + \sin^2 \theta_{13} e^{2i\delta}|}$$

- The observed baryon asymmetry is generated through resonant leptogenesis. The CP violation in  $\nu$ -oscillations and leptonic asymmetries are *directly related*.