

Toward a Minimum Branching for Dark Matter Annihilation into Electromagnetic Final States

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*Toward a Minimum Branching for Dark Matter Annihilation
into Electromagnetic Final States*

JD, Robert J. Scherrer, and Thomas J. Weiler

Electroweak Bremsstrahlung in Dark Matter Annihilation

Nicole F. Bell, JD, Thomas D. Jacques, Thomas J. Weiler

Outline

- Bounds on DM annihilation to Standard Model products

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- Neutrinos only model and higher order corrections from SM processes

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- Neutrinos only model and higher order corrections from SM processes
- We calculate the branching ratio to charged leptons in a neutrinos only model and show that it is typically $\gtrsim 1\%$, when the scale of the dark matter mass(es) exceeds the electroweak scale, M_W

Bounds from Disappearance and Appearance Rates

- Thermal Relic Scale: $\Omega_\chi h^2 = 0.1$ requires an annihilation cross section of approximately $\langle v\sigma \rangle \sim 3 \times 10^{-26} \text{ cm}^3 \text{ sec}^{-1}$

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- Alteration of Halo density profiles

M. Kaplinghat, L. Knox, and M.S. Turner, Phys.Rev.Lett. 85,3335 (2000)

Bounds from Disappearance and Appearance Rates

- Assuming that DM annihilation or decay is into SM particles neutrinos can be used to set an upper bound on the cross section.

J.F. Beacom, N.F. Bell, and G.D. Mack, Phys.Rev.Lett. **99**, 231301 (2007)

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- Measuring ν and γ fluxes leads to bounds on the cross section
- For a $Br(\gamma\gamma)$ of 10^{-4} , the ν bounds generally provide tighter constraints than photons for large dark matter masses ($M_\chi > 100$ MeV)

Embedding Neutrinos Only into Field Theory

- At minimum introduce a particle, B with mass M_B , that mediates the annihilation to neutrinos

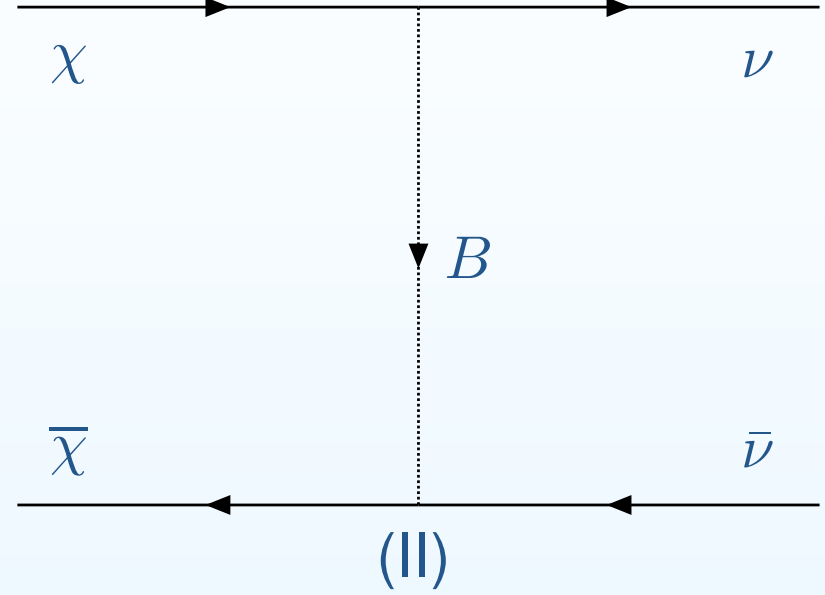
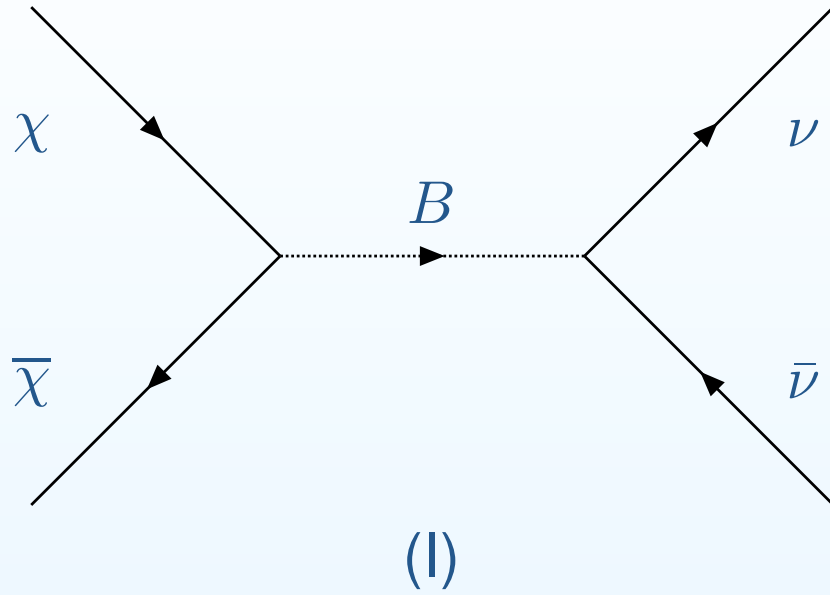
Embedding Neutrinos Only into Field Theory

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- Such an embedding proves to be unstable against higher order corrections

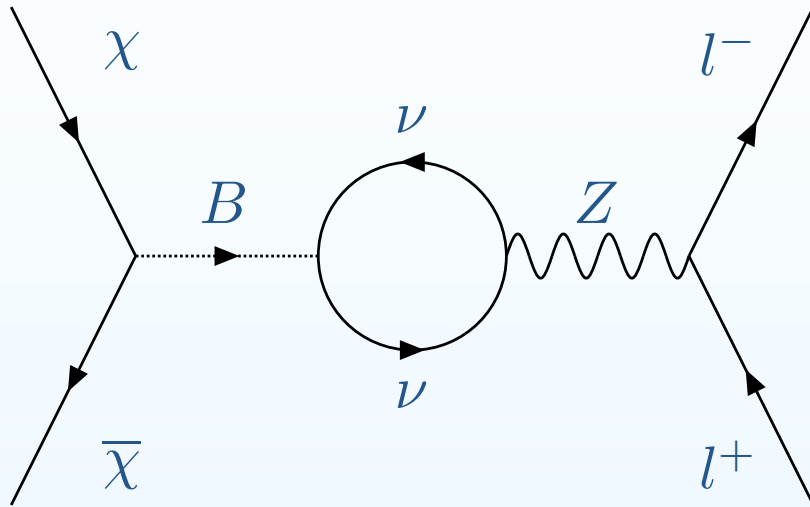
Embedding Neutrinos Only into Field Theory

- At minimum introduce a particle, B with mass M_B , that mediates the annihilation to neutrinos
- Such an embedding proves to be unstable against higher order corrections
- Finely tuned effective theory below some scale Λ , or UV completion introduces a new spectrum at Λ

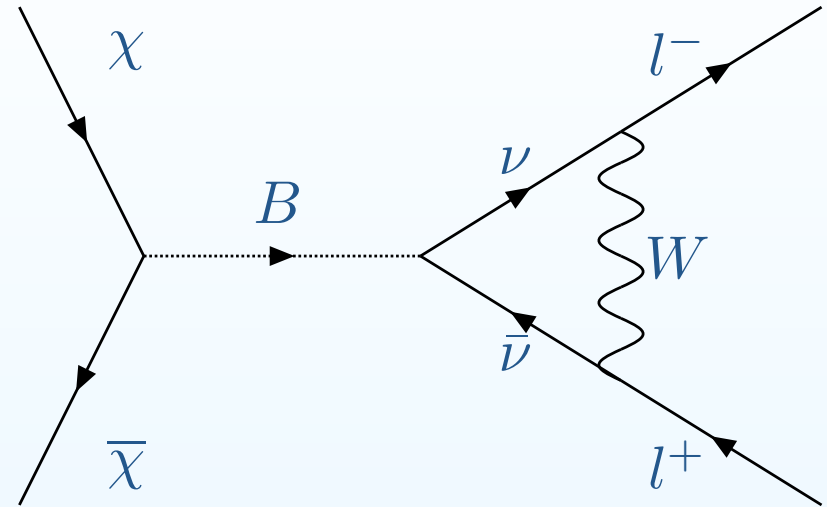
Neutrinos Only at Tree Level



Loop Corrections

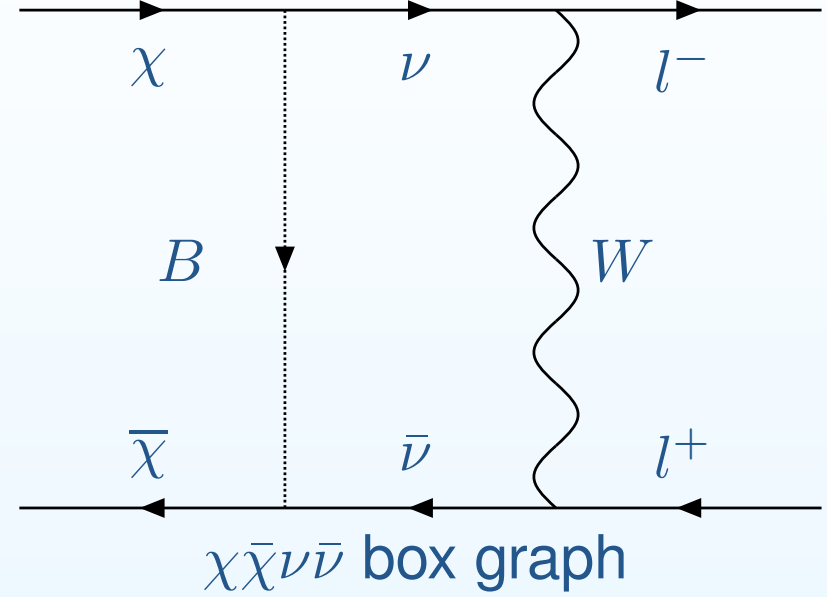
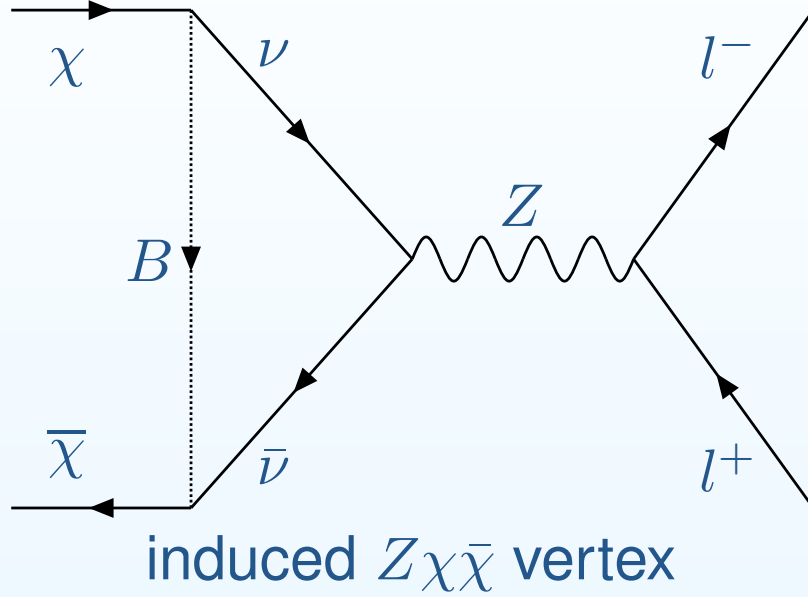


oblique correction, B-Z mixing

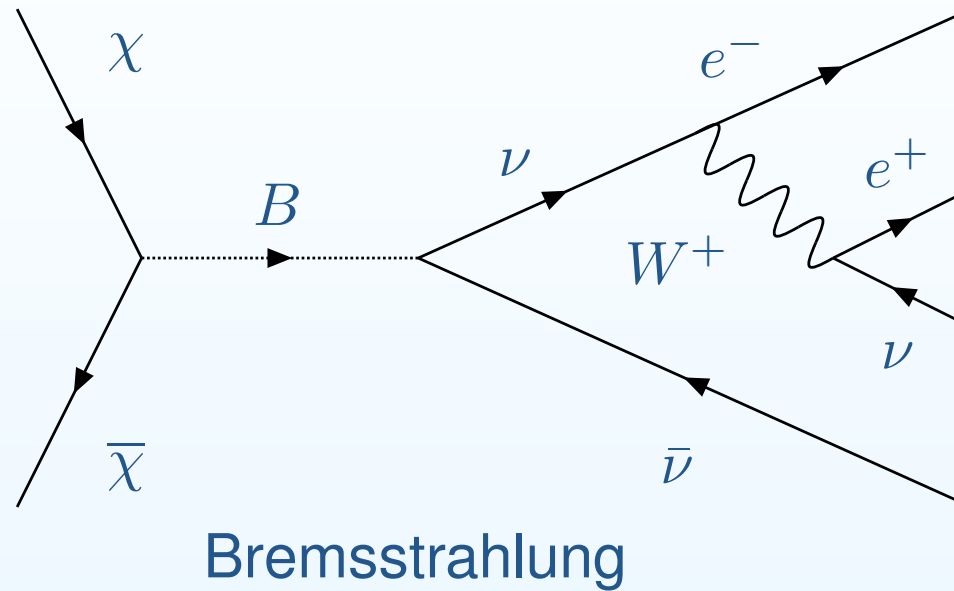


induced $B l^+ l^-$ vertex

Loop Corrections



Bremsstrahlung



first estimated by
M. Kachelriess and P.D. Serpico, Phys.Rev.D 76, 063516 (2007).

The Box Diagram

- The (scalar) box diagram is logarithmically divergent
- There is no counterterm to suppress it
- We introduce a Pauli-Villars regulator which gives the cutoff scale Λ
- The final ratio

$$\mathcal{R} = \frac{\langle v \sigma(\chi\bar{\chi} \rightarrow l^+l^-) \rangle}{\langle v \sigma(\chi\bar{\chi} \rightarrow \nu\bar{\nu}) \rangle}$$

will depend on three independent ratios of the four variables M_χ , M_B , Λ , and M_W .

Box Formula

We work in unitary gauge where the W propagator is given by

$$\frac{-i(g_{\mu\nu} - k_\mu k_\nu / M_W^2)}{k^2 - M_W^2 + iM_W \Gamma_W}$$

The overall ratio becomes

$$\mathcal{R} = \left(\frac{M_B^2 + M_\chi^2}{128 \cdot 8\pi \cdot M_W^2} \right)^2 \left[\int d\xi \ln \left| \frac{\Delta_\Lambda^2}{\Delta_B^2} \right| \right]^2$$

where

$$\begin{aligned} \Delta_B^2 &\equiv \xi_1 M_B^2 + \xi_2 M_W^2 - (\xi_1 - \xi_1^2 - 2\xi_1 \xi_2 + 4\xi_3 \xi_4) M_\chi^2 \\ \Delta_\Lambda^2 &\equiv \xi_1 \Lambda^2 + \xi_2 M_W^2 - (\xi_1 - \xi_1^2 - 2\xi_1 \xi_2 + 4\xi_3 \xi_4) M_\chi^2 \end{aligned}$$

Range of Validity

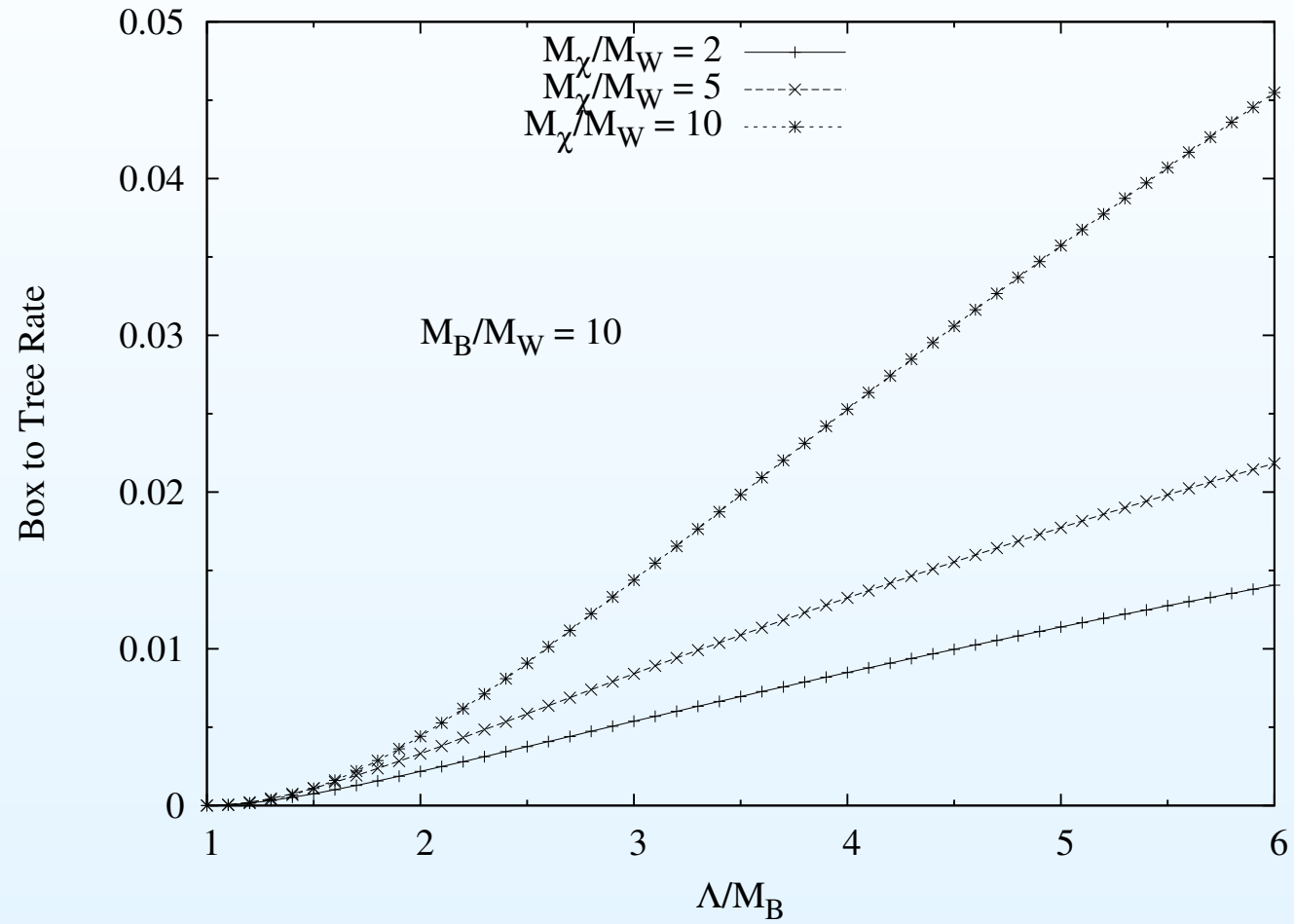
M_B and M_χ are constrained by the requirement that the value $\Omega_\chi h^2 = 0.1$ requires an annihilation cross section of $\langle v\sigma \rangle \gtrsim 3 \times 10^{-26} \text{ cm}^3 \text{ sec}^{-1}$. In the model presented here, in the non-relativistic $s \sim (2 M_\chi)^2$ and large $M_B^2 \gg s$ limits, we have

$$\langle v \sigma(\chi\bar{\chi} \rightarrow \nu\bar{\nu}) \rangle = \left(\frac{g_B^4}{8\pi} \right) \left(\frac{2 M_\chi^2}{M_B^4} \right)$$

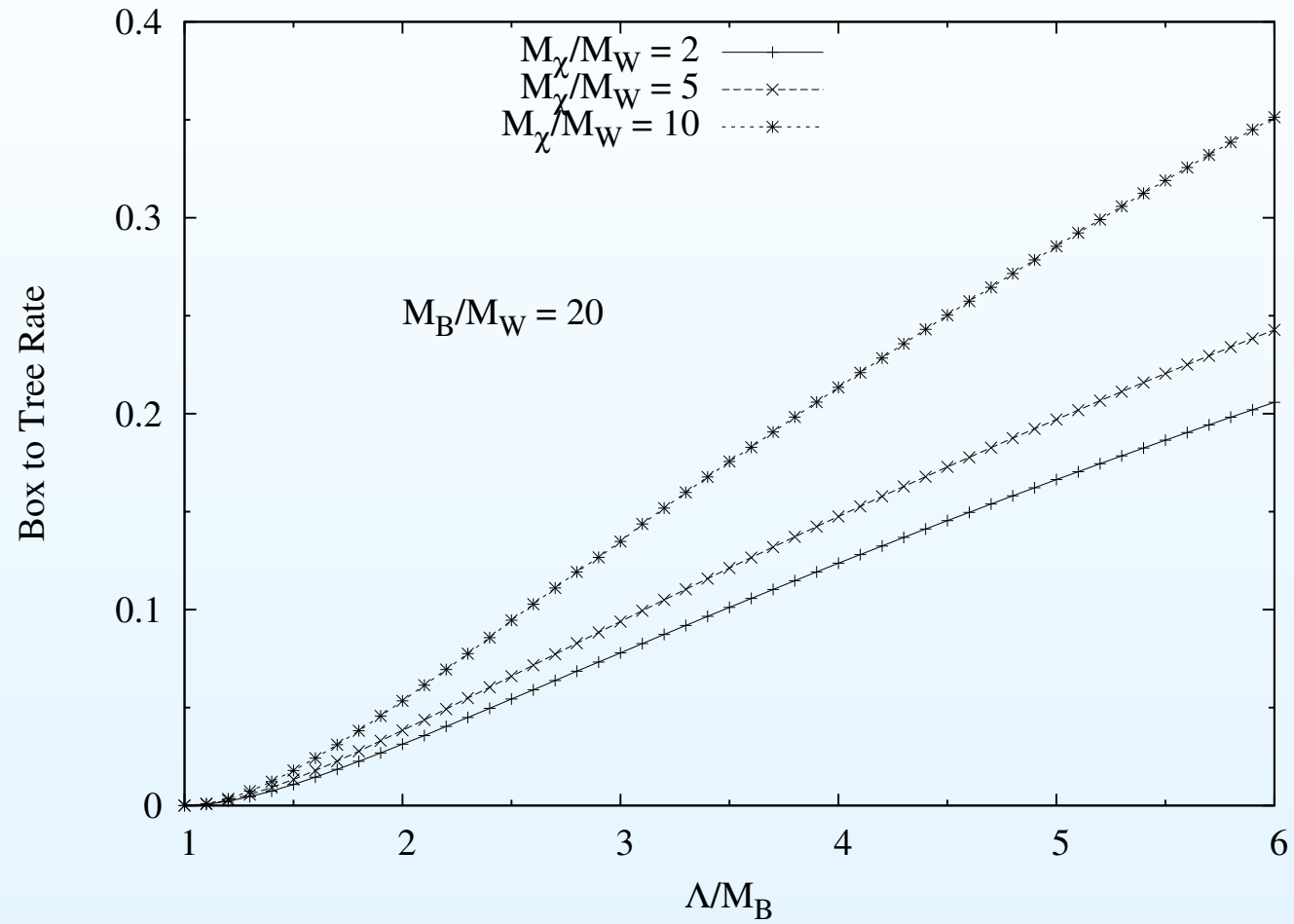
If we consider only $M_\chi > 100 \text{ GeV}$, and restrict the coupling to the perturbative regime $g_B^2 \leq 4\pi$, then we obtain the bound

$$1 < \frac{M_B}{M_\chi} < 26,$$

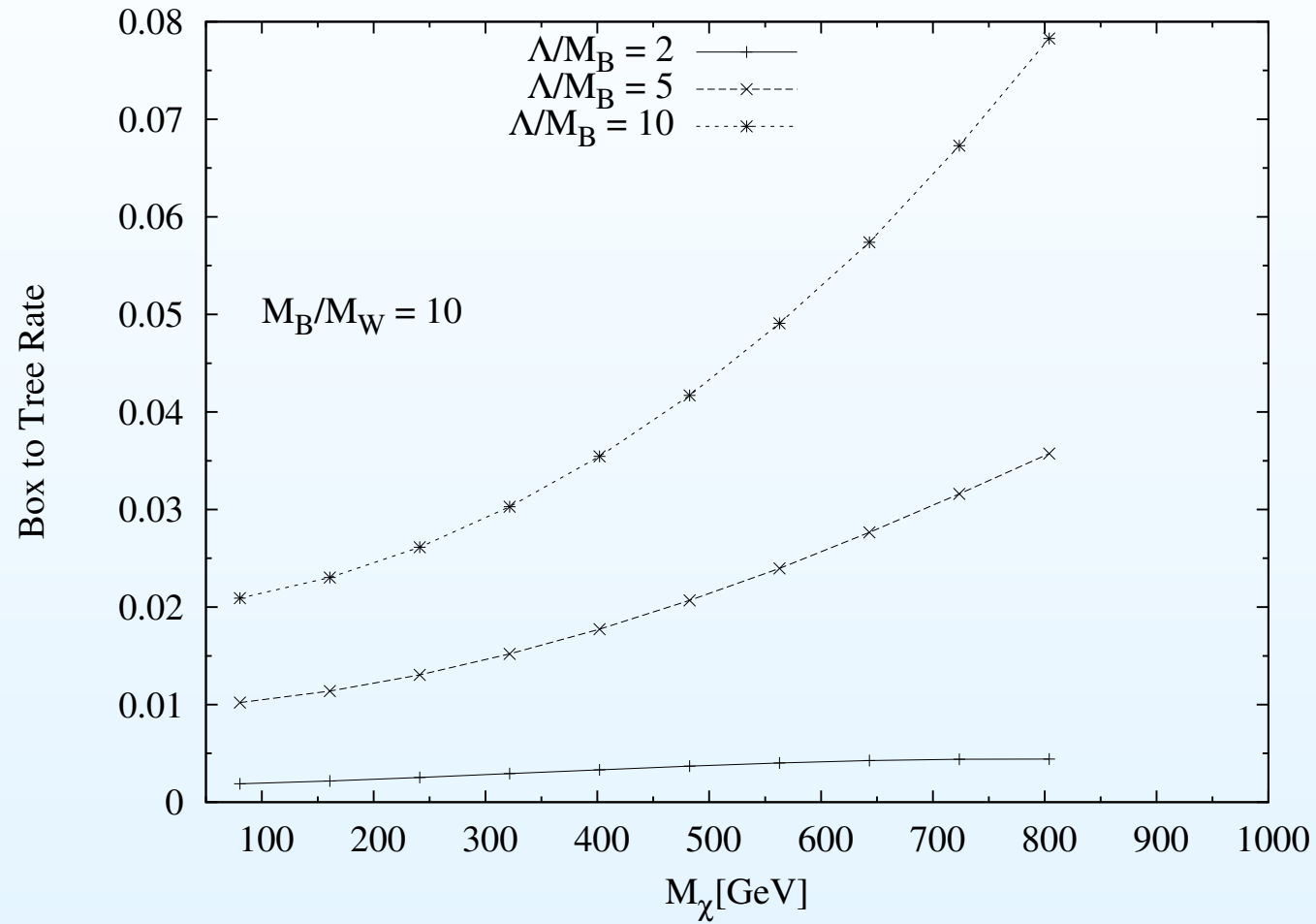
Running Λ



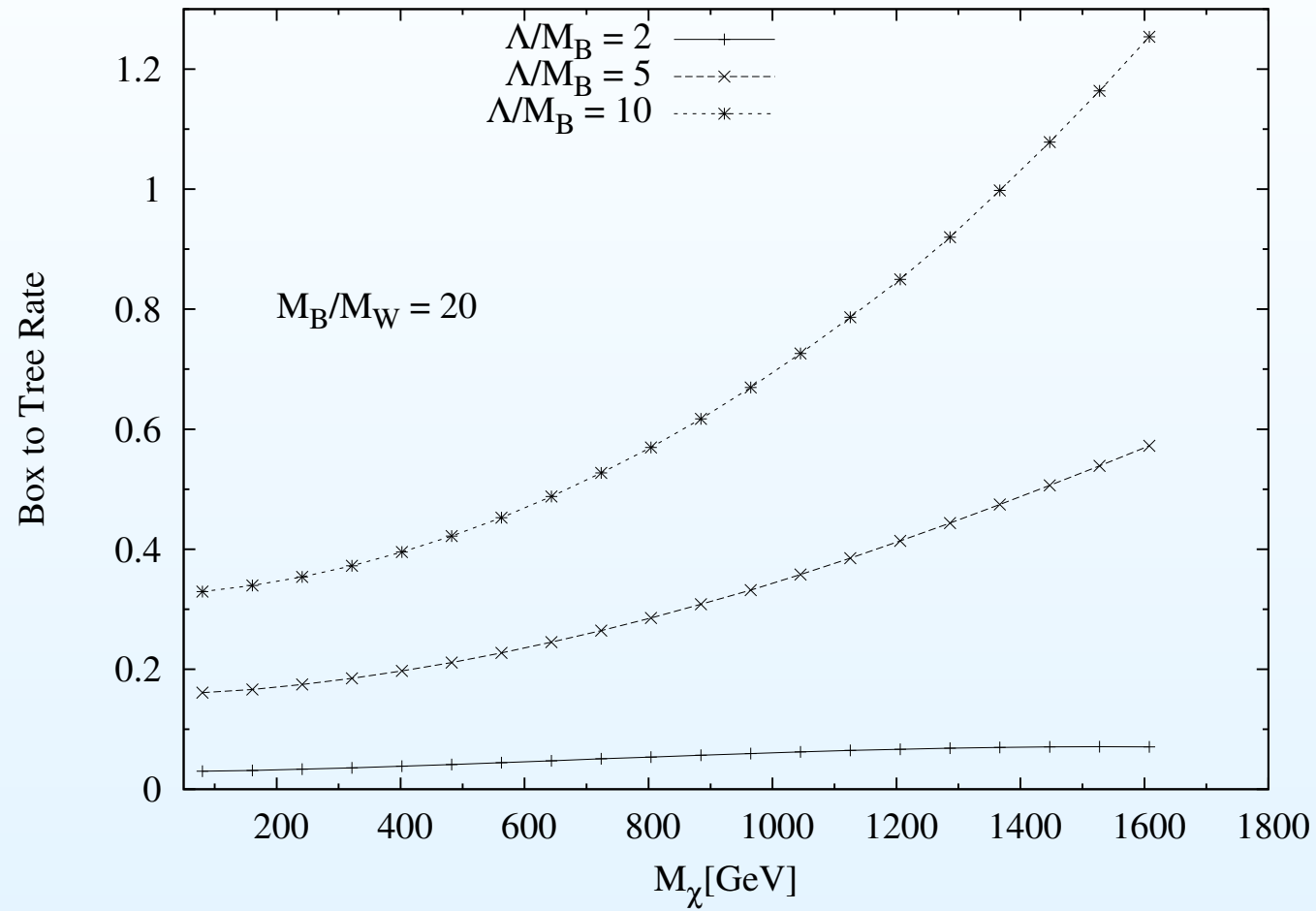
Running Λ



Running M_χ



Running M_χ



Conclusions

- Whenever neutrino-antineutrino pairs are produced in the final state, they can always be converted into l^+l^- pairs through W (or Z) exchange
- The branching ratio to charged leptons in a neutrinos only model is typically $\gtrsim 1\%$, when the scale of the dark matter mass(es) exceeds the electroweak scale, M_W
- This is true for both the box diagram as well as W or Z -strahlung.
- Neutrino and photon limits are likely to be competitive in the 100GeV to a few TeV range