# A 4th generation in the MSSM

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# Outline

#### Motivations

- Constraints on  $m_{b'}$  and  $m_{t'}$  from EW precision data and the axial anomaly  $(\Delta M_Z)$
- Limits on  $\tan \beta$  from the introduction of a 4th generation due to perturbativity requirements
- Introduction of model to extend limits on  $\tan\beta$
- Conclusions

# **Motivation**

- Why not have a 4th generation?
- Current experimental bounds of 4th generation (PDG):  $m_{t'} > 220 \text{ GeV}$ ,  $m_{b'} > 190 \text{ GeV}$ ,  $m_{\tau'} > 100 \text{ GeV}$ ,  $m_{\nu'} > 50 \text{ GeV}$
- To show there is room for a 4th generation in the MSSM
- The requirement of perturbativity limits  $1 < \tan \beta < 1.17$ , we would like to extend this range
- To investigate what masses are possible for b' and t'
- LHC will explore the region around 1 TeV giving a good chance to either find or rule out a 4th generation

#### **Masses of 4th Generation Quarks**

- The electroweak precision variables S, T, and U limit the difference between  $m_{b'}$  and  $m_{t'}$ .
- New chiral fermions increase S, T, and U. For example<sup>a</sup>:

$$S_f = \frac{N_c}{6\pi} \left\{ 2(4Y+3)x_1 + 2(3-4Y)x_2 - 2Y\ln\frac{x_1}{x_2} + \left[ \left(\frac{3}{2} + 2Y\right)x_1 + Y \right] G(x_1) + \left[ \left(\frac{3}{2} - 2Y\right)x_2 - Y \right] G(x_2) \right\}$$
  
Where  $x_1 = \frac{m_1^2}{m_Z^2}$  and  $G(x) = -4\sqrt{4x-1} \arctan\frac{1}{\sqrt{4x-1}}$ 

$$S = -0.13 \pm 0.10$$
  
 $T = -0.13 \pm 0.11$   
 $U = 0.20 \pm 0.12$ 

<sup>a</sup>He, Polonsky, Su, hep-ph/0102144v2

#### $m_{b'}$ vs $m_{t'}$ as constrained by S, T, U

 $m_h = 115 \text{ GeV}, m_{\nu'} = 100 \text{ GeV}, m_{\tau'} = 150 \text{ GeV}, \tan \beta = 1.5$  $m_{t'}$ 375 350 325 300 275





# **Axial Anomaly**

Loop corrections to the Z<sup>0</sup> propagator are directly dependent upon heavy chiral fermions<sup>a</sup>



- Their result generalized to 4 generations:  $\Delta M_Z = \frac{1.6 \times 10^{-5}}{2M_Z} (m_t^2 + m_{t'}^2 - m_{b'}^2)$
- Not considered in previous work on the 4th generation, this further constrains the difference between the masses of the b' and t'

<sup>&</sup>lt;sup>a</sup>Anselm, Dombey, Leader, Phys Lett B312 (1993)

#### $m_{b'}$ vs $m_{t'}$ with $\Delta M_Z$ Constraint





# **Limits on** $\tan \beta$

- In the MSSM  $\tan \beta < 60$  due to b Yukawa exploding
- MSSM with 4 generations (MSSM4) was explored by Gunion, McKay, and Pois <sup>a</sup>. They found  $\tan \beta < 3$
- The problem is that for  $\tan \beta > 3$ , the Yukawa coupling of the b' quark becomes non-perturbative near 1 TeV
- Current values of  $m_t$  and the lower limits on  $m_{b'}$  and  $m_{t'}$  cause the constraint on  $\tan \beta$  to be even more restrictive
- With current data we find that no  $tan \beta$  allows perturbativity up to the GUT scale

<sup>&</sup>lt;sup>a</sup>Phys. Lett. B 334 pp. 339-347 (1994)

# $\tan \beta$ vs Cutoff Scale ( $\lambda > 2$ )



# **Our Model**

#### Motivations:

- The limit on  $\tan\beta$  is too narrow
- We introduce two new vector like Higgs doublets at the TeV scale to redefine the Yukawa coupling of the b', t', and  $\tau'$  particles
- This allows  $\tan \beta$  to be as large as 2.5 while the Yukawa couplings remain perturbative **up to the GUT scale**
- Increasing  $\tan\beta$  will help increase the upper limit on the Higgs mass
- Below the scale of new physics we have the low-energy effective theory of the MSSM

#### Model



 $\lambda_{b'} = \lambda_{b'}^{(1)} + \lambda_D \frac{M}{M_H} \Theta(M - \mu)$ 



 $\lambda_{t'} = \lambda_{t'}^{(1)} + \lambda_U \frac{M'}{M_H} \Theta(M' - \mu)$ 

#### Model

The extra particles in the loops preserve the perturbativity of the Yukawa couplings  $\lambda_{b'}^{(1)}$  and  $\lambda_D$  up to the GUT scale



$$16\pi^2 \frac{d}{dt} \lambda_{b'}^{(1)} = \lambda_{b'} \left[ 6 \left( \lambda_{b'}^{(1)} \right)^2 + 3\lambda_D^2 + \lambda_U^2 + \left( \lambda_{t'}^{(1)} \right)^2 + \left( \lambda_{\tau'}^{(1)} \right)^2 - c_i^d g_i^2 \right]$$

# **Conclusions and Outlook**

- The 4th generation is still very much alive, but the gap to find it is closing
- The LHC will probe the likely mass range of the 4th generation extensively
- We have one model that allows Yukawa couplings to be perturbative up to GUT scale and increases the limits on  $\tan \beta$
- From the axial anomaly  $|m_{t'} m_{b'}| < 10 \text{ GeV}$
- Increased Higgs production due to new particles in the loop contributing to gluon fusion
- Larger  $\tan \beta$  increases upper bound on  $m_h$ . At tree level,  $m_h^2 = m_Z^2 \cos^2 2\beta + ...$
- New quarks increase higgs mass through new loop corrections:  $\Delta(m_h^2)_{\rm new} \sim (60 \ {\rm GeV})^2$