#### Dark Matter in the Left Right Twin Higgs Model

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Work done with Shufang Su and Jessica Goodman arXiv:0712.1234v1 [hep-ph]

#### Outline

- Left-Right Twin Higgs Model
- Relic Density Analysis
- Conclusion

 Chacko, Goh, and Harnik: arXiv:hep-ph/0506256v1

• Solution to Little Hierarchy Problem

- To avoid EW precision constraints, add a second Higgs  $\hat{H}$  that couples to gauge bosons only

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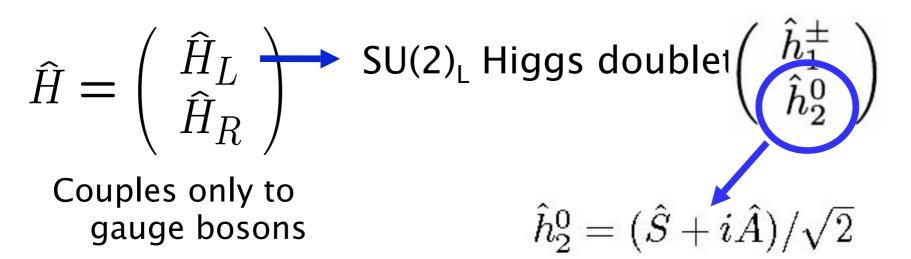
# $\hat{H} = \left(\begin{array}{c} \hat{H}_L\\ \hat{H}_R \end{array}\right)$

Couples only to gauge bosons

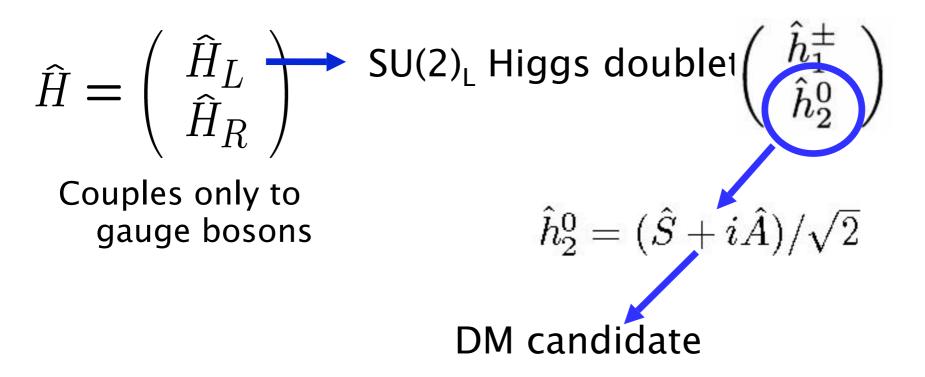
Left-Right Twin Higgs Model  $H = \begin{pmatrix} H_L \\ H_R \end{pmatrix} \xrightarrow{\text{SM Higgs doublet}} h_{SM}$   $\hat{H} = \begin{pmatrix} \hat{H}_L \\ \hat{H}_R \end{pmatrix} \xrightarrow{\text{SU}(2)_{L} \text{ Higgs doublet}} \begin{pmatrix} \hat{h}_1^{\pm} \\ \hat{h}_2^{0} \end{pmatrix}$ 

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- The lighter one of  $\hat{S}/\hat{A}\,$  is stable, weakly interacting

→ Natural WIMP candidates

In addition to the CW potential, we can add terms to the Lagrangian:

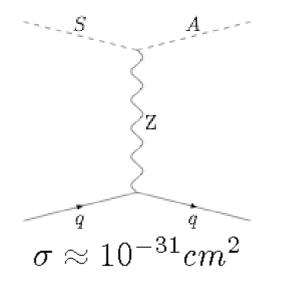
 $\Delta L = -\hat{\mu}^2 \hat{H}_L^{\dagger} \hat{H}_L - \lambda_4 |H_L^{\dagger} \hat{H}_L|^2 - \frac{\lambda_5}{2} ((H_L^{\dagger} \hat{H}_R)^2 + h.c.)$ 

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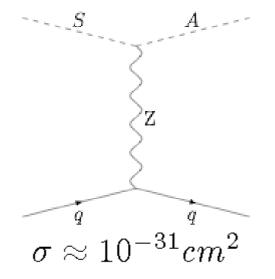
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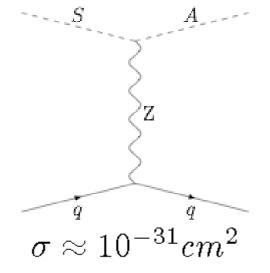


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Impose neutral mass splitting to make it kinematically forbidden.

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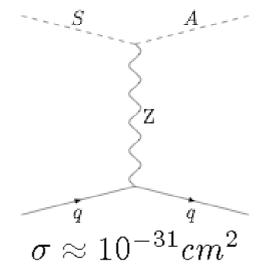


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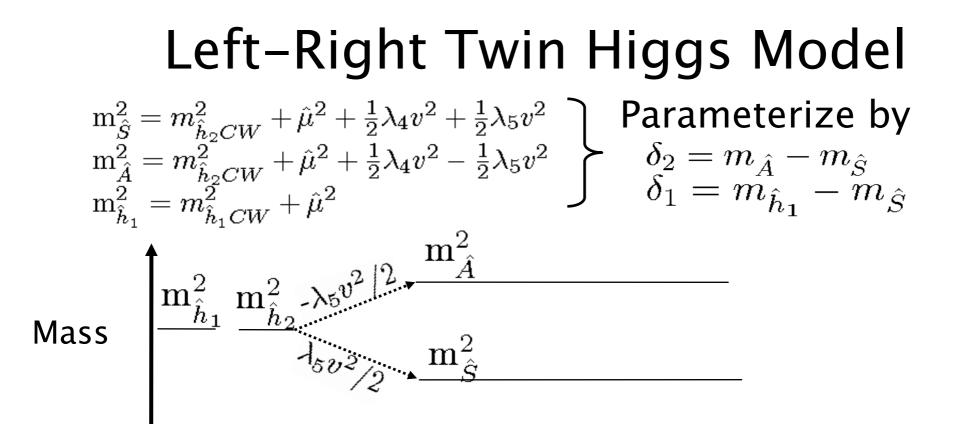
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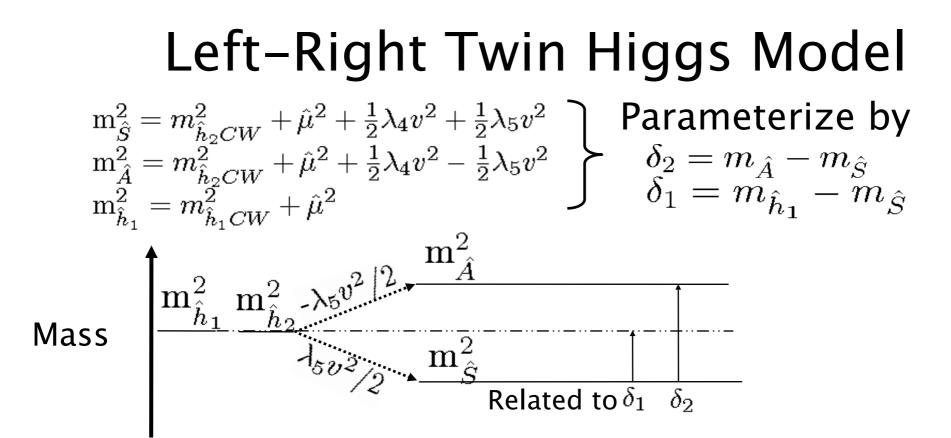
$$\begin{split} \mathbf{m}_{\hat{S}}^2 &= m_{\hat{h}_2 CW}^2 + \hat{\mu}^2 + \frac{1}{2}\lambda_4 v^2 + \frac{1}{2}\lambda_5 v^2 \\ \mathbf{m}_{\hat{A}}^2 &= m_{\hat{h}_2 CW}^2 + \hat{\mu}^2 + \frac{1}{2}\lambda_4 v^2 - \frac{1}{2}\lambda_5 v^2 \\ \mathbf{m}_{\hat{h}_1}^2 &= m_{\hat{h}_1 CW}^2 + \hat{\mu}^2 \end{split}$$

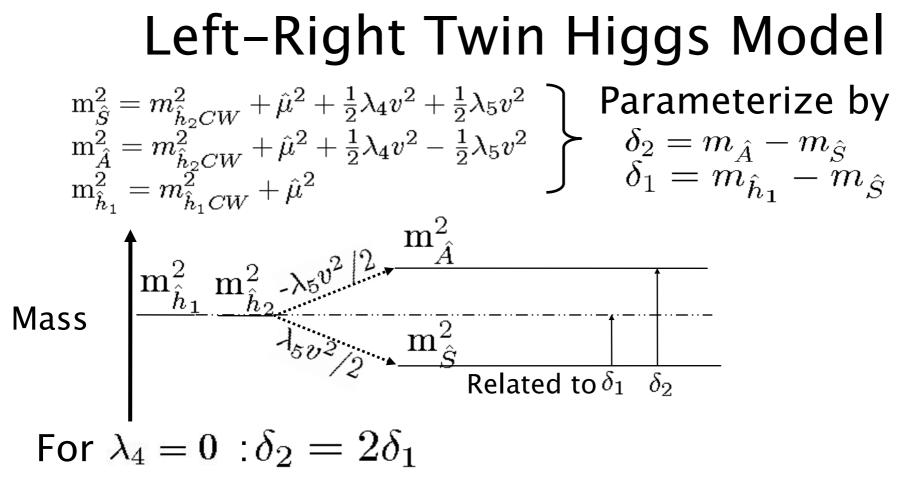
# $\begin{array}{l} \text{Left-Right Twin Higgs Model} \\ m_{\hat{S}}^{2} = m_{\hat{h}_{2}CW}^{2} + \hat{\mu}^{2} + \frac{1}{2}\lambda_{4}v^{2} + \frac{1}{2}\lambda_{5}v^{2} \\ m_{\hat{A}}^{2} = m_{\hat{h}_{2}CW}^{2} + \hat{\mu}^{2} + \frac{1}{2}\lambda_{4}v^{2} - \frac{1}{2}\lambda_{5}v^{2} \\ m_{\hat{h}_{1}}^{2} = m_{\hat{h}_{1}CW}^{2} + \hat{\mu}^{2} \end{array} \right\} \begin{array}{l} \text{Parameterize by} \\ \delta_{2} = m_{\hat{A}} - m_{\hat{S}} \\ \delta_{1} = m_{\hat{h}_{1}} - m_{\hat{S}} \end{array}$

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$$\begin{array}{l} \text{Mass} \end{array}$$







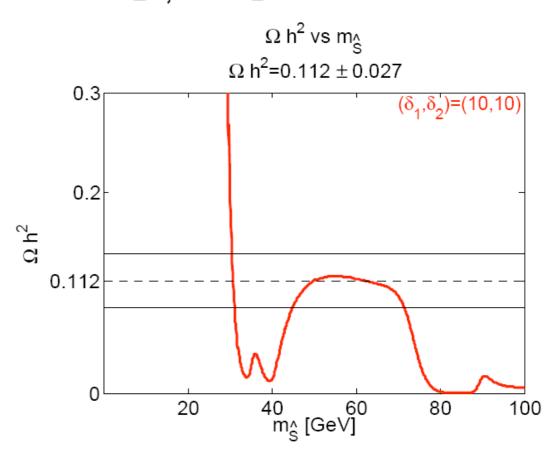
In general, for  $\lambda_4 
eq 0$  :  $\delta_2 
eq 2\delta_1$ 

We looked at both cases, treating splittings as free parameters.

#### Relic Density Analysis

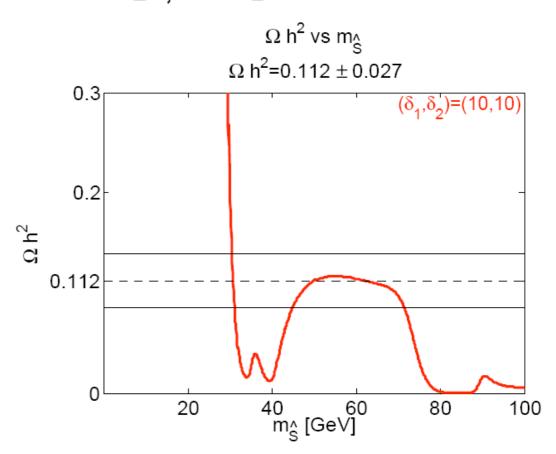
- WMAP:  $0.085 < \Omega h_{CDM}^2 < 0.139$  at  $3\sigma$
- Solve Boltzmann equation  $\frac{dY}{dT} = \sqrt{\frac{\pi g_*(T)}{45}} M_{Pl} \langle \sigma v \rangle (Y^2(T) - Y_{eq}^2(T))$
- micrOmegas: considers co-annihilations when mass splittings are small
- Modest choice of parameters yields
  - Low mass region:  $m_{\hat{S}} < 100~{
    m GeV}$
  - High mass region:  $400~{
    m GeV} < m_{\hat{S}} < {
    m a~few~TeV}$

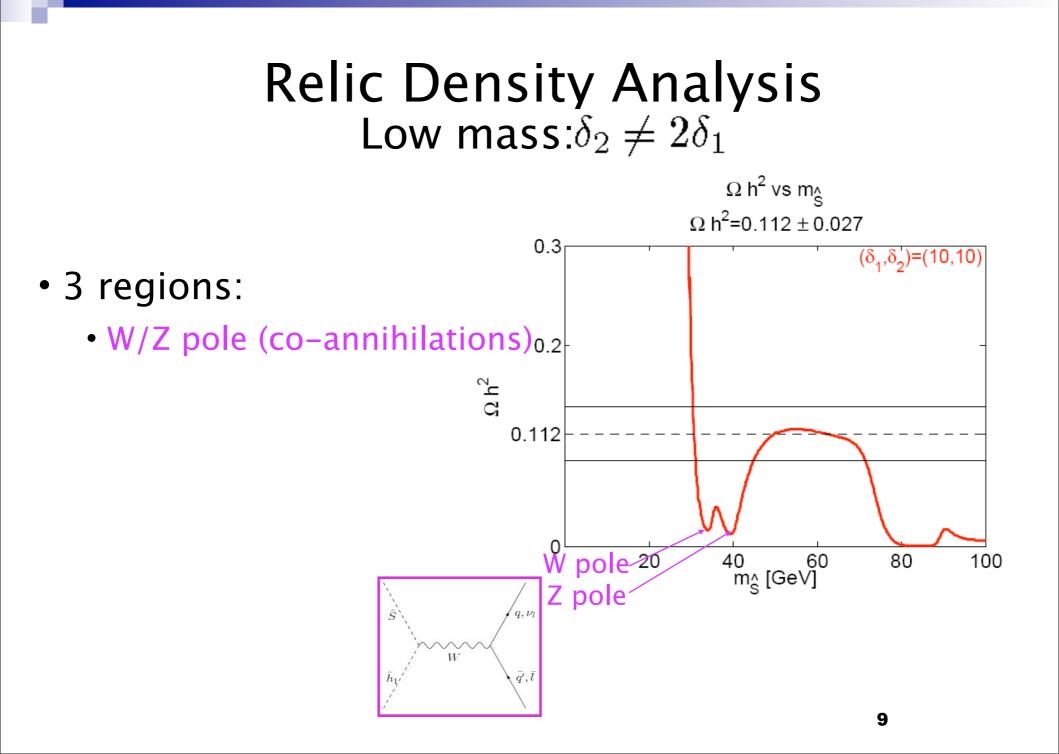
#### Relic Density Analysis Low mass: $\delta_2 \neq 2\delta_1$

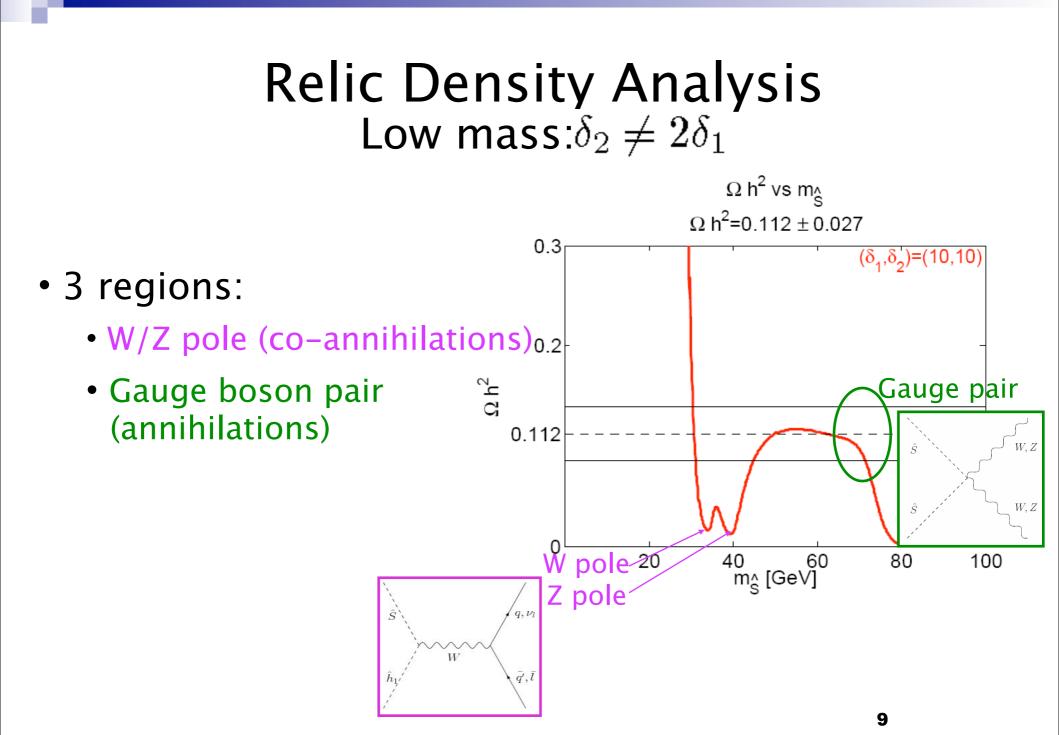


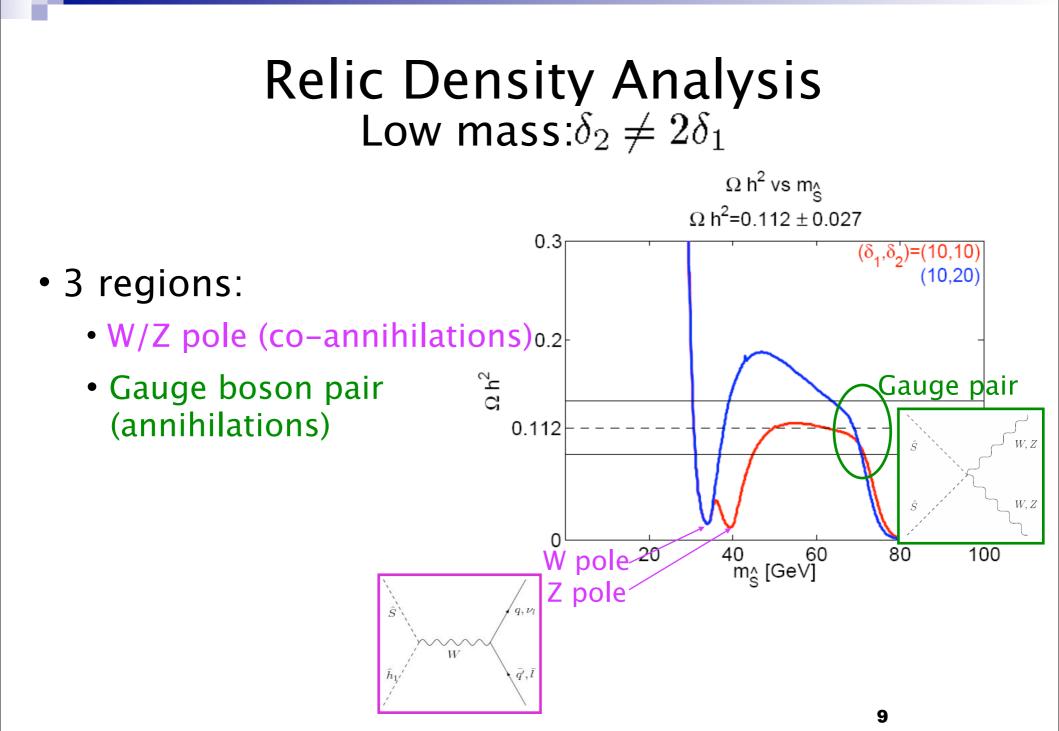
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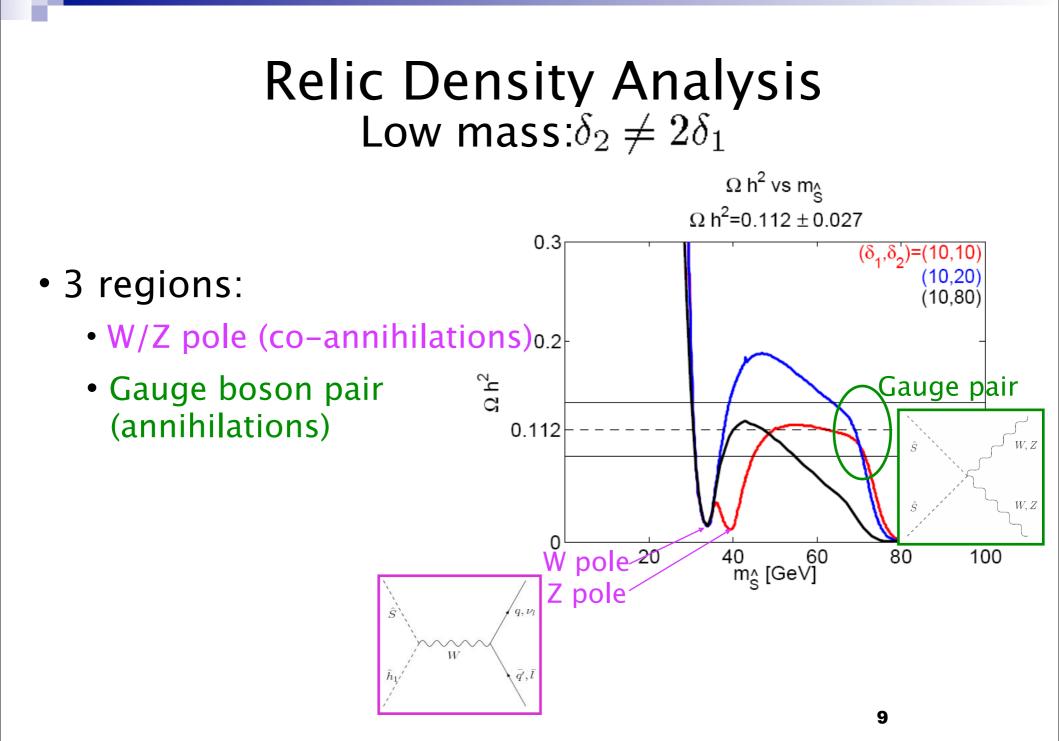
• 3 regions:

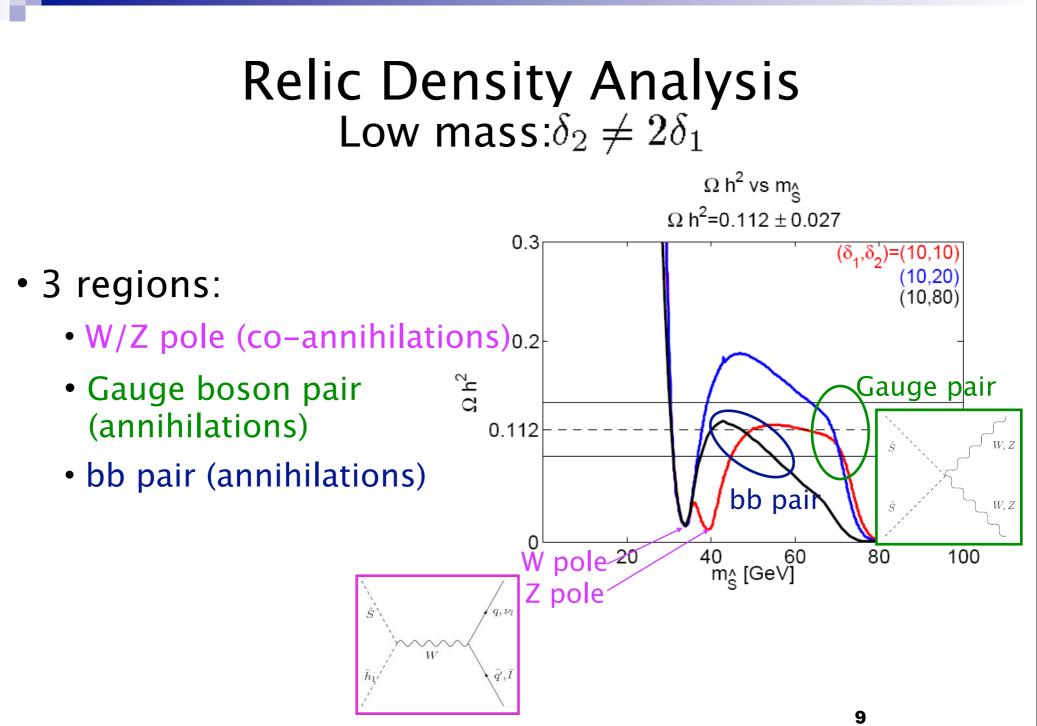








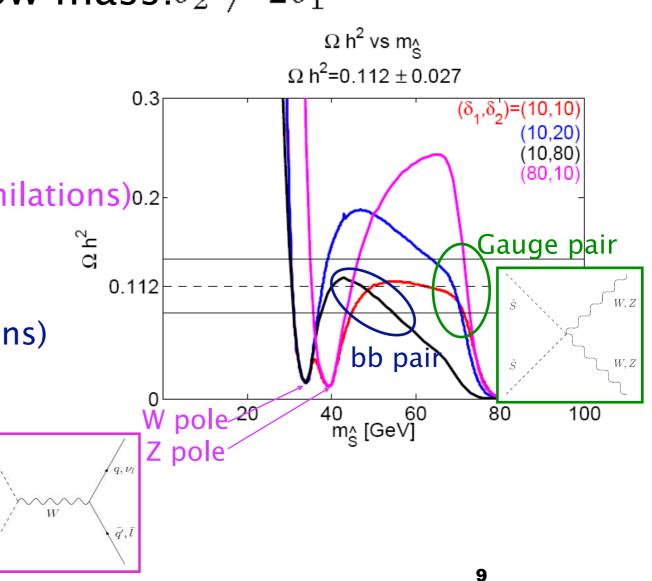




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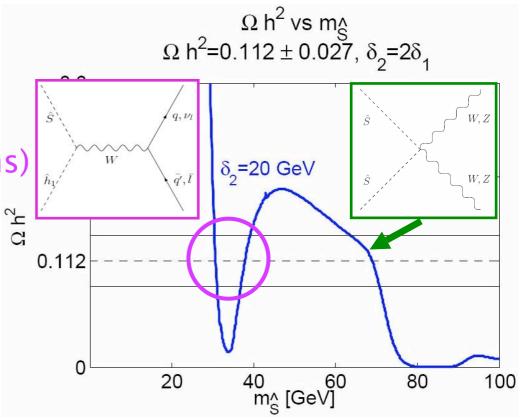


- W/Z pole (co-annihilations)0.2
- Gauge boson pair (annihilations)
- bb pair (annihilations)



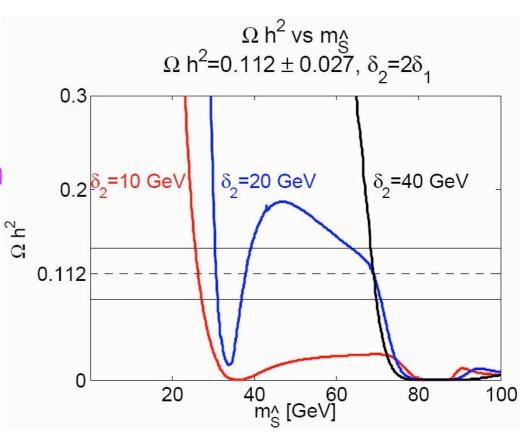
#### Relic Density Analysis Low mass: $\delta_2 = 2\delta_1$

- 2 regions:
  - •W/Z pole (co-annihilations)
  - Gauge boson pair (annihilations)

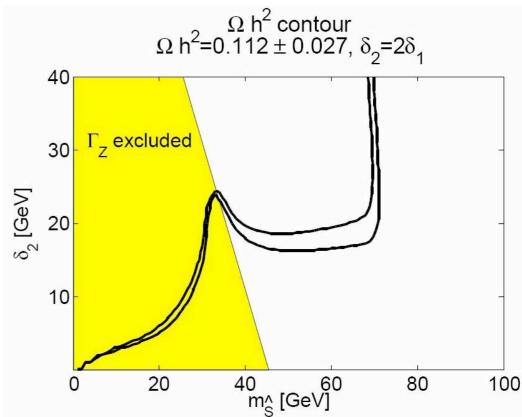


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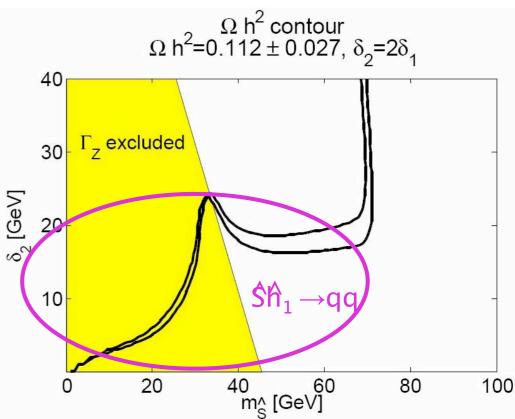
- 2 regions:
  - •W/Z pole (co-annihilation
  - Gauge boson pair (annihilations)
- Change with splittings



•  $\delta_2 = 2\delta_1$  case

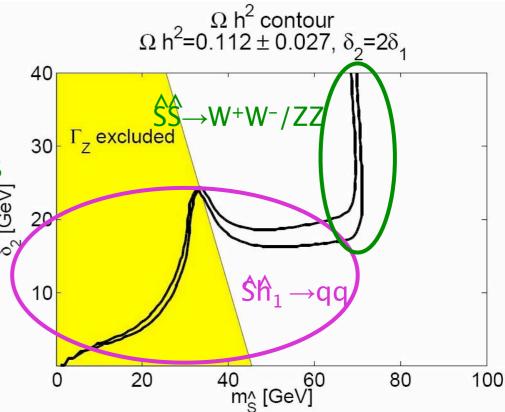


- $\delta_2 = 2\delta_1$  case
  - Co-annihilations: W/Z pole

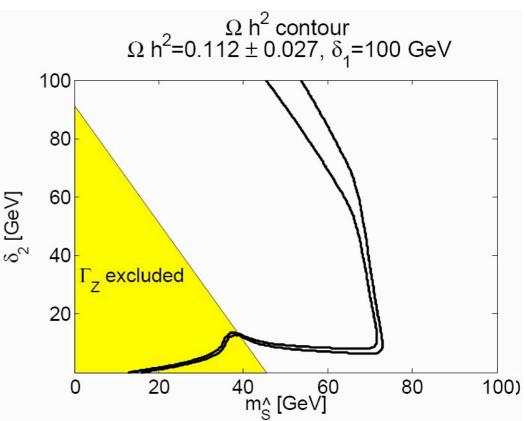


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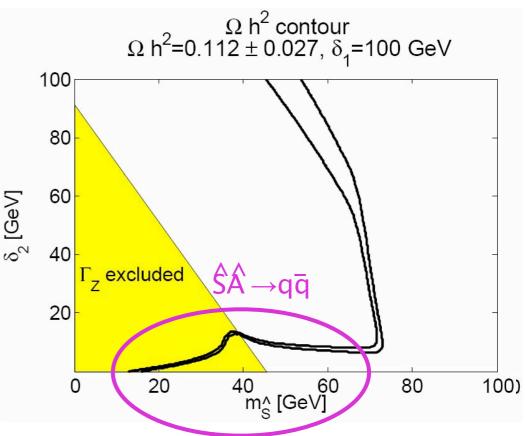
  - Co-annihuauone.
     Annihilations: gauge bosons



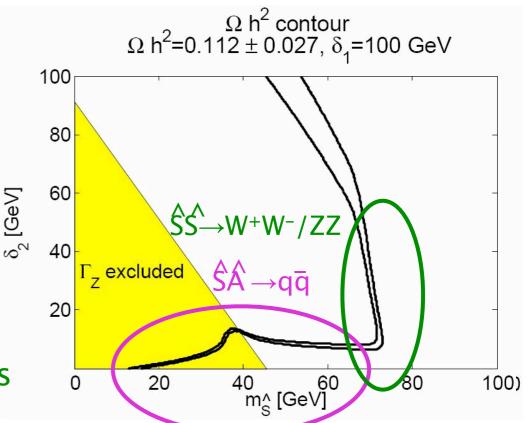
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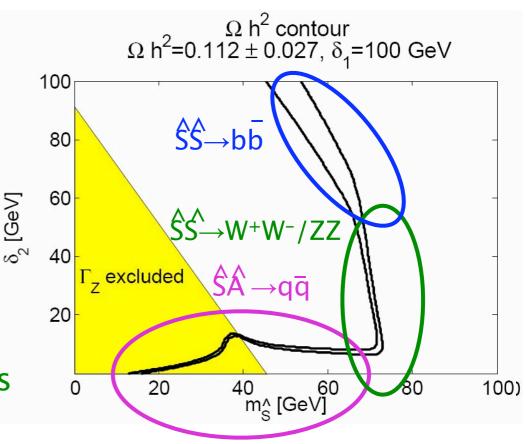
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  - Co-annihilations: W/Z pole

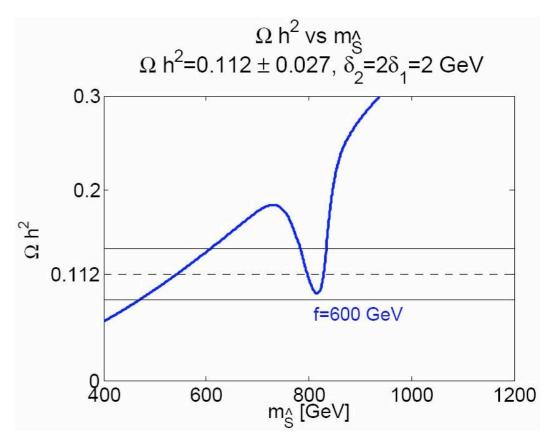


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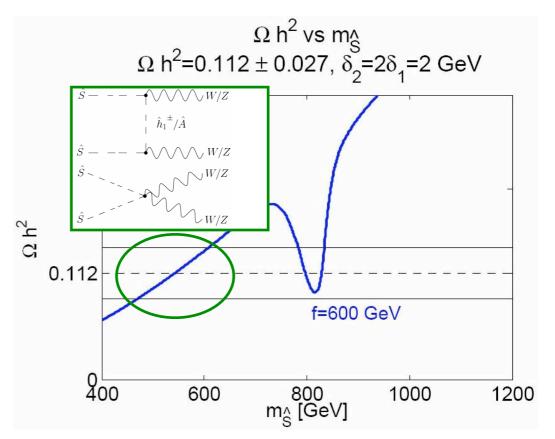
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  - Annihilations: bb pair



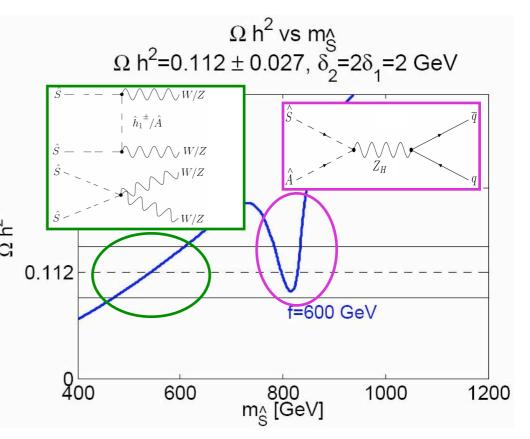


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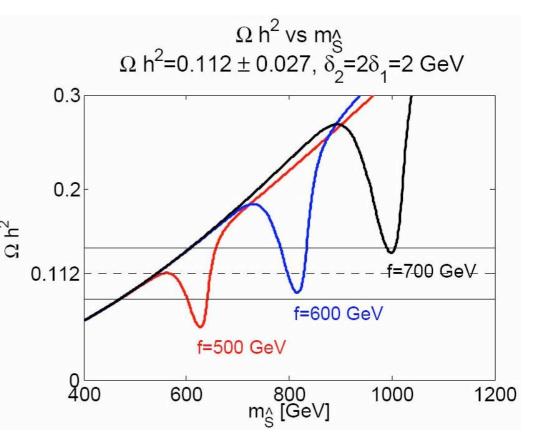
- Two regions:
  - Bulk (annihilations)



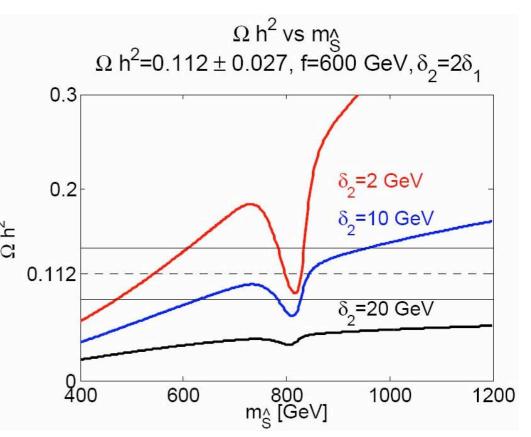
- Two regions:
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  - Z<sub>H</sub> pole (co-annihilations) mg~m<sub>ZH</sub>/2



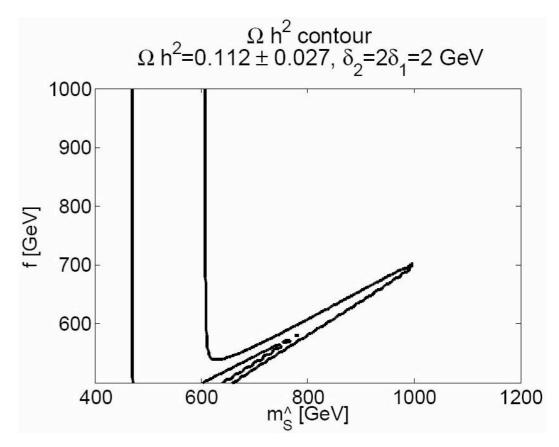
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- Regions change by:
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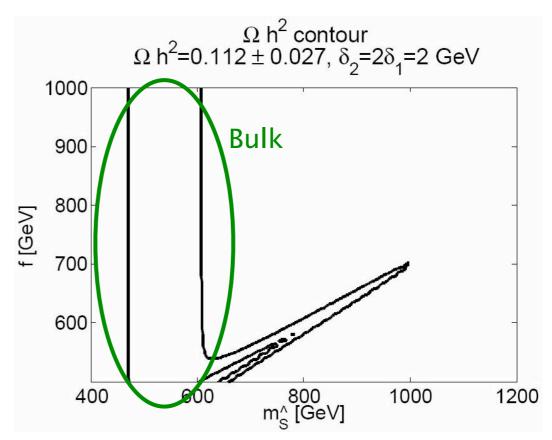
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  - changing  $\delta_{1,2}$



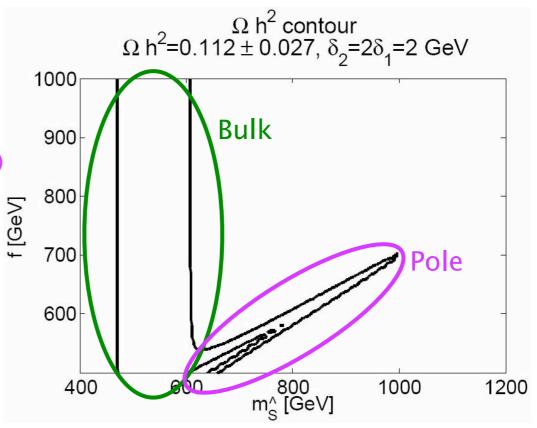
• m<sub>s</sub> –f contour



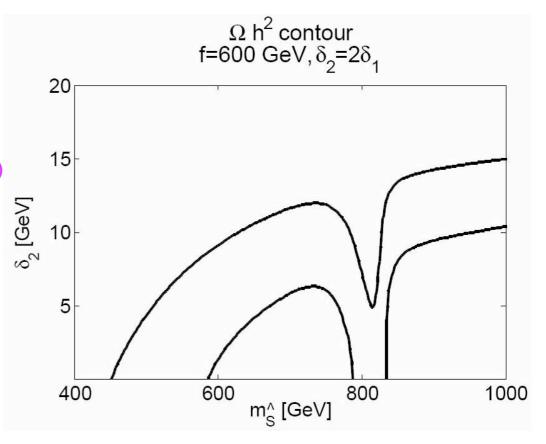
- m<sub>s</sub> –f contour
  - Bulk: m<sub>§</sub> constant



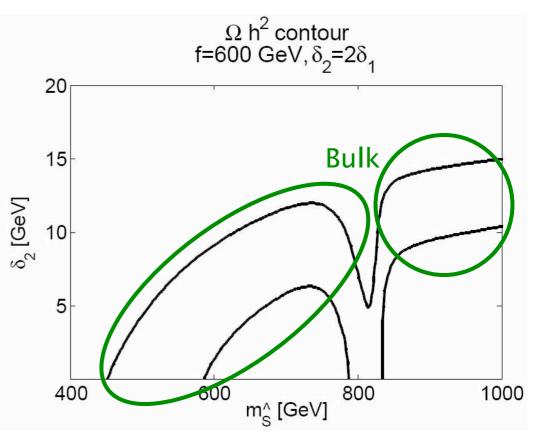
- $m_{\hat{s}}$  –f contour
  - Bulk: m<sub>§</sub> constant
  - Pole: m<sub>s</sub> varies (m<sub>s</sub>~ m<sub>zH</sub>/2)

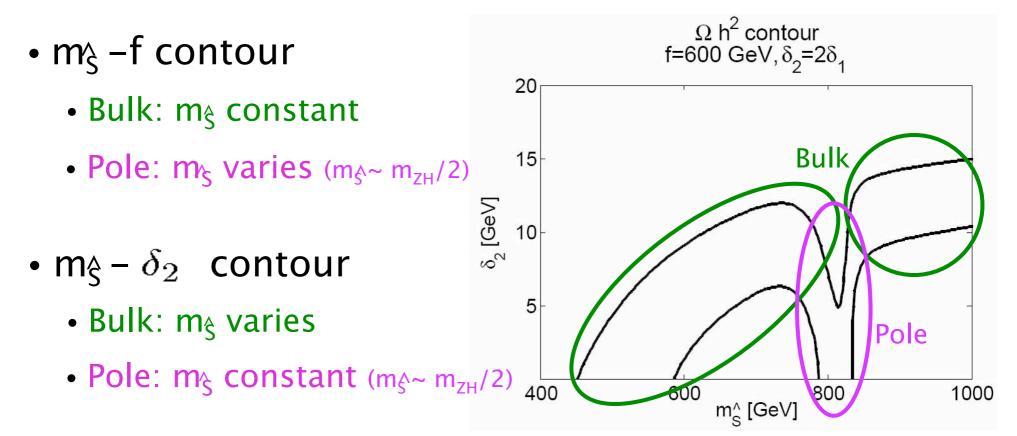


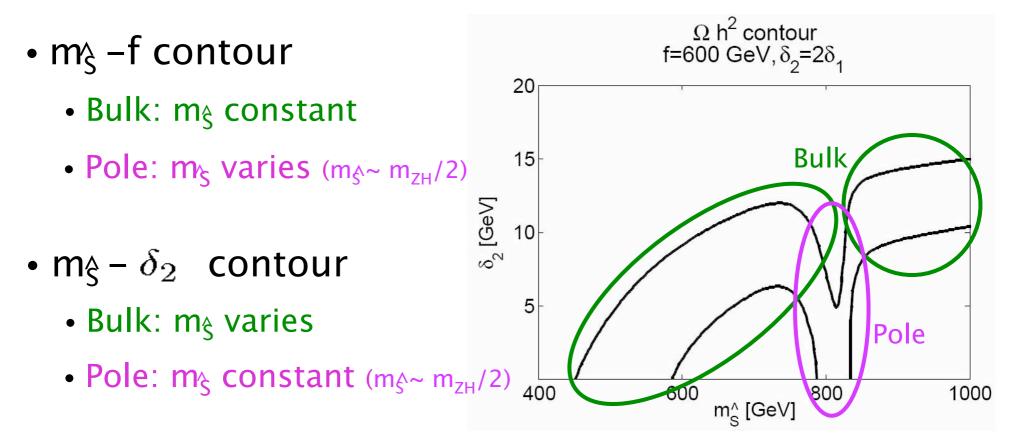
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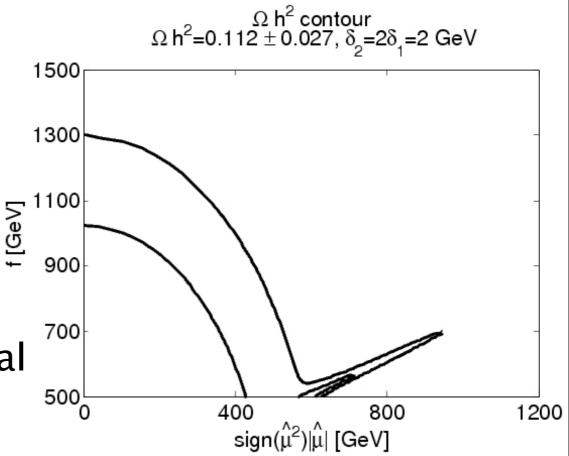
Large areas of parameter space where WMAP results are accessible

• Recall for  $\lambda_4 = 0$ :

$$\begin{split} \mathbf{m}_{\hat{S}}^2 &= m_{\hat{h}_2 C W}^2 + \hat{\mu}^2 + \frac{1}{2} \lambda_5 v^2 \\ \mathbf{m}_{\hat{A}}^2 &= m_{\hat{h}_2 C W}^2 + \hat{\mu}^2 - \frac{1}{2} \lambda_5 v^2 \\ \mathbf{m}_{\hat{h}_1}^2 &= m_{\hat{h}_1 C W}^2 + \hat{\mu}^2 \end{split}$$

- There exists regions where  $\hat{\mu}^2 = 0$
- Bulk mass is then given entirely by CW potential

$$\begin{split} \mathbf{m}_{\hat{S}}^{-} &= m_{\hat{h}_{2}CW}^{-} + \frac{1}{2}\lambda_{5}v^{-} \\ \mathbf{m}_{\hat{A}}^{2} &= m_{\hat{h}_{2}CW}^{2} - \frac{1}{2}\lambda_{5}v^{2} \\ \mathbf{m}_{\hat{h}_{1}}^{2} &= m_{\hat{h}_{1}CW}^{2} \end{split}$$



# Conclusion

- Left Right Twin Higgs Model provides a natural dark matter candidate
- Can obtain WMAP results with a wide range of splittings for low mass region
- High mass region requires a little tuning (splittings of a few GeV), and works with minimal setup ( $\hat{\mu}^2 = 0$ )
- Thank you!