

# (MINIMAL) WALKING TECHNICOLOR

## CONSTRAINTS AND UNITARITY

Roshan Foadi  
roshan@fysik.sdu.dk

Southern Denmark University

PHENO 2008, April 29<sup>th</sup> 2008

In collaboration with:

**Mads T. Frandsen, Matti O. Järvinen, and Francesco Sannino**  
(Southern Denmark University, Odense)

## TECHNICOLOR:

- 1 Take a strongly interacting gauge theory with a chiral symmetry.
- 2 Let the gauge coupling be strong enough to trigger spontaneous chiral symmetry breaking.
- 3 Let the electroweak gauge group be a subgroup of the chiral symmetry group. (Weinberg 79, Susskind 79)

## TECHNICOLOR:

- 1 Take a strongly interacting gauge theory with a chiral symmetry.
- 2 Let the gauge coupling be strong enough to trigger spontaneous chiral symmetry breaking.
- 3 Let the electroweak gauge group be a subgroup of the chiral symmetry group. (Weinberg 79, Suskind 79)

Example: **QCD** !

How do the SM fermions interact with the **technifermion condensate** to acquire mass ?

## EXTENDED TECHNICOLOR:

A new gauge theory, at a higher scale, with SM fermions and technifermions in the same multiplet.

(Lane and Eichten 80)

How do the SM fermions interact with the **technifermion condensate** to acquire mass ?

## EXTENDED TECHNICOLOR:

A new gauge theory, at a higher scale, with SM fermions and technifermions in the same multiplet.

(Lane and Eichten 80)

## PROBLEM:

Tension between SM fermion masses and flavor changing neutral currents.

**Walking dynamics** enhances the technifermion bilinears, not the SM fermion bilinears.

WALKING TECHNICOLOR:

Enhanced SM fermion masses and suppressed FCNC. (Holdom 81)

**Walking dynamics** enhances the technifermion bilinears, not the SM fermion bilinears.

WALKING TECHNICOLOR:

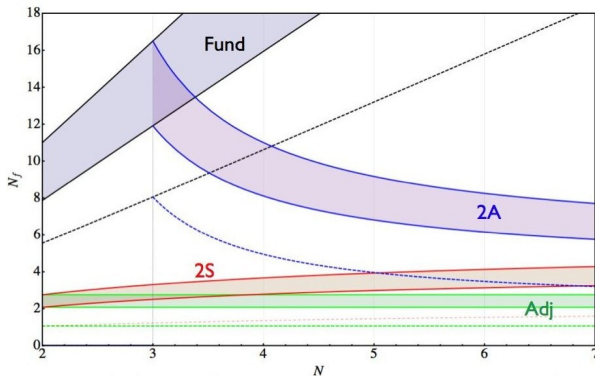
Enhanced SM fermion masses and suppressed FCNC. (Holdom 81)

QUESTION:

Which gauge theories do display walking behavior ?



## Conformal Windows in Ladder Approximation



(Sannino and Dietrich 06)

# $S$ PARAMETER IN WALKING TECHNICOLOR

The smaller the number of new fermions, the smaller the perturbative  $S$  parameter.

PERTURBATIVE  $S$

$$S_{\text{pert}} = \frac{N_D}{6\pi}$$

# $S$ PARAMETER IN WALKING TECHNICOLOR

The smaller the number of new fermions, the smaller the perturbative  $S$  parameter.

PERTURBATIVE  $S$

$$S_{\text{pert}} = \frac{N_D}{6\pi}$$

Take  $N_c = 2$ . Then, close to the conformal window:

- $S_{\text{pert}} = \frac{8}{6\pi} \simeq 0.42$  in the **fundamental representation**.
- $S_{\text{pert}} = \frac{3}{6\pi} \simeq 0.16$  in the **adjoint** representation.

# S PARAMETER IN WALKING TECHNICOLOR

The smaller the number of new fermions, the smaller the perturbative  $S$  parameter.

PERTURBATIVE  $S$

$$S_{\text{pert}} = \frac{N_D}{6\pi}$$

Take  $N_c = 2$ . Then, close to the conformal window:

- $S_{\text{pert}} = \frac{8}{6\pi} \simeq 0.42$  in the **fundamental representation**.
- $S_{\text{pert}} = \frac{3}{6\pi} \simeq 0.16$  in the **adjoint** representation.

(MINIMAL) WALKING TECHNICOLOR:

Walking Technicolor with small  $N_f$  in a higher dimensional representation. (Sannino and Tuominen 05)

## MINIMAL WALKING TECHNICOLOR:

$N_c = 2, N_f = 2$  in the **adjoint** or **two-index symmetric**.

## NEXT TO MINIMAL WALKING TECHNICOLOR:

$N_c = 3, N_f = 2$  in the **two-index symmetric**.

## MINIMAL WALKING TECHNICOLOR:

$N_c = 2$ ,  $N_f = 2$  in the **adjoint** or **two-index symmetric**.

## NEXT TO MINIMAL WALKING TECHNICOLOR:

$N_c = 3$ ,  $N_f = 2$  in the **two-index symmetric**.

MWT on the lattice: tests in progress. (Catteral and Sannino 07; Del Debbio, Frandsen, Panagopoulos, and Sannino 08)

Alternatives: **Partially gauged Technicolor**, and **Split Technicolor**.  
(Dietrich and Sannino 05)

## NEW PHYSICS:

$$\mathcal{L}_{\text{newphysics}} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + i\bar{Q}_L\gamma_\mu D^\mu Q_L + i\bar{Q}_R\gamma_\mu D^\mu Q_R$$

## NEW PHYSICS:

$$\mathcal{L}_{\text{newphysics}} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + i\bar{Q}_L\gamma_\mu D^\mu Q_L + i\bar{Q}_R\gamma_\mu D^\mu Q_R$$

- Plus possible lepton doublets or matter in the adjoint.
- Plus ETC interactions.



## (M)WT IN THE EFFECTIVE LAGRANGIAN

Walking dynamics and small  $S$  in the effective Lagrangian: Weinberg Sum Rules. (R.F., Frandsen, Ryttov, and Sannino 07)

## (M)WT IN THE EFFECTIVE LAGRANGIAN

Walking dynamics and small  $S$  in the effective Lagrangian: Weinberg Sum Rules. (R.F., Frandsen, Ryttov, and Sannino 07)

- Model the spin-1 spectrum with a vector and an axial resonances and a close-to-continuum spectrum in the near conformal region.
- Use some trustable approximation to compute the contribution of the near conformal region to the vacuum polarization amplitudes. (Appelquist and Sannino 98)

# WEINBERG SUM RULES

FIRST WSR: NO IMPORTANT CONTRIBUTIONS FROM THE NEAR CONFORMAL REGION.

$$F_V^2 - F_A^2 = F_\pi^2$$

# WEINBERG SUM RULES

FIRST WSR: NO IMPORTANT CONTRIBUTIONS FROM THE NEAR CONFORMAL REGION.

$$F_V^2 - F_A^2 = F_\pi^2$$

SECOND WSR: IMPORTANT CONTRIBUTIONS FROM THE NEAR CONFORMAL REGION.

$$F_V^2 M_V^2 - F_A^2 M_A^2 = a \frac{8\pi^2}{d(\mathbb{R})} F_\pi^4, \quad a \in \mathcal{O}(1)$$

# WEINBERG SUM RULES

FIRST WSR: NO IMPORTANT CONTRIBUTIONS FROM THE NEAR CONFORMAL REGION.

$$F_V^2 - F_A^2 = F_\pi^2$$

SECOND WSR: IMPORTANT CONTRIBUTIONS FROM THE NEAR CONFORMAL REGION.

$$F_V^2 M_V^2 - F_A^2 M_A^2 = a \frac{8\pi^2}{d(\mathbb{R})} F_\pi^4, \quad a \in \mathcal{O}(1)$$

S PARAMETER, OR “ZEROTH WSR”: IMPORTANT CONTRIBUTIONS FROM THE NEAR CONFORMAL REGION.

$$S = 4\pi F_\pi^2 \left[ \frac{1}{M_V^2} + \frac{1}{M_A^2} - a \frac{8\pi^2 F_\pi^2}{d(\mathbb{R}) M_V^2 M_A^2} \right]$$

# S PARAMETER IN (M)WT

- Notice:  $S_{\text{pert}} \simeq 0.16$  both in QCD and MWT.
- $S_{\text{QCD}} \simeq 2S_{\text{pert}}$ .
- Therefore, in the worst case scenario we can expect a QCD-like (running) behavior in MWT:  $S_{\text{MWT}} \simeq 0.3$ .

# S PARAMETER IN (M)WT

- Notice:  $S_{\text{pert}} \simeq 0.16$  both in QCD and MWT.
- $S_{\text{QCD}} \simeq 2S_{\text{pert}}$ .
- Therefore, in the worst case scenario we can expect a QCD-like (running) behavior in MWT:  $S_{\text{MWT}} \simeq 0.3$ .

However in walking theories there is a negative contribution.

## S IN WALKING TECHNICOLOR

$$S = 4\pi F_\pi^2 \left[ \frac{1}{M_V^2} + \frac{1}{M_A^2} - a \frac{8\pi^2 F_\pi^2}{d(\text{R})M_V^2 M_A^2} \right], \quad a \in \mathcal{O}(1)$$

Therefore in Walking Technicolor  $S$  is expected to be smaller than in running theories. (Appelquist and Sannino 95; Kurachi, Shrock, and Yamawaki 07)

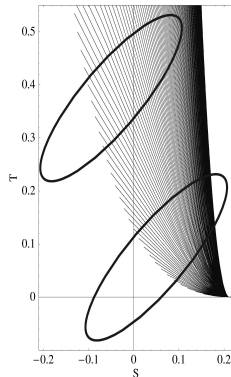
# $S$ PARAMETER IN (M)WT

- The full  $S$  is the sum of the Technicolor contribution and the contribution from the new leptons  $S = S_{\text{TC}} + S_{\text{leptons}}$ .
- A positive value of  $T$  is favored by electroweak precision data.
- The new leptons can easily give a negative contribution to  $S$  and a positive contribution to  $T$ .

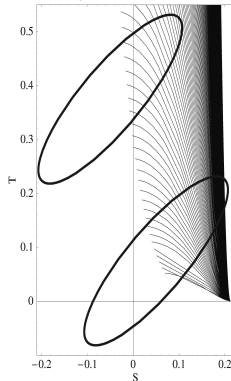


# S PARAMETER IN (M)WT

$$y = 1$$

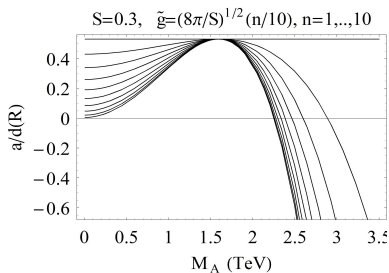
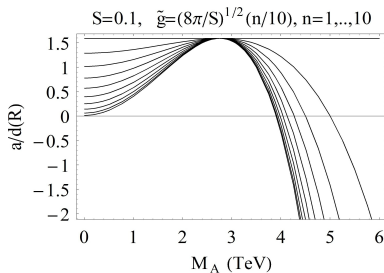


$$y = 1/3$$



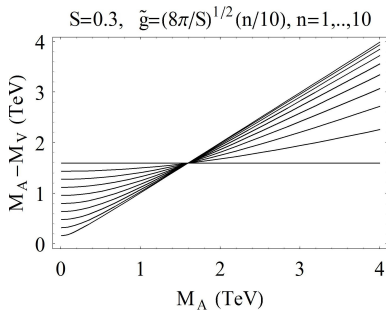
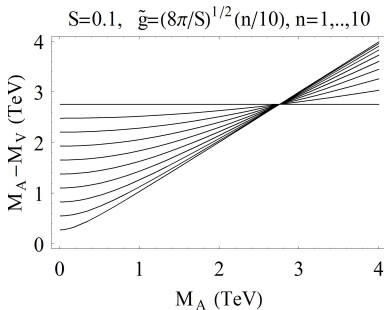
# WALKING REGIMES IN (M)WT

- 1 Impose WSR-1.
- 2 Impose WSR-0 and small  $S$ .
- 3 For different values of the vector coupling, check on the behavior of  $a$ , which is expected to be positive  $\mathcal{O}(1)$  in a walking theory.



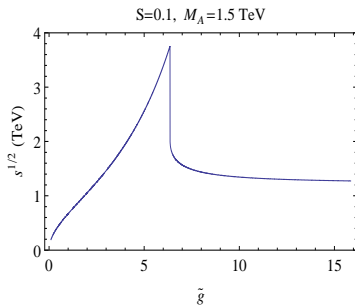
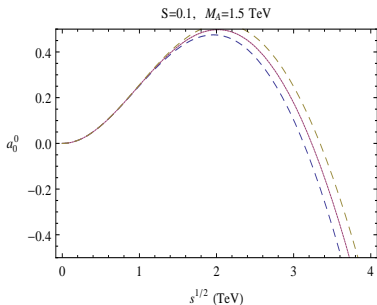
# SPECTRUM IN (M)WT

- In a walking regime the **vector can be heavier than the axial**.
- In (M)WT, if the spin-1 resonances are within the LHC range, then  **$M_V > M_A$** .



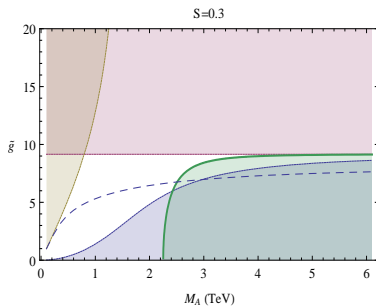
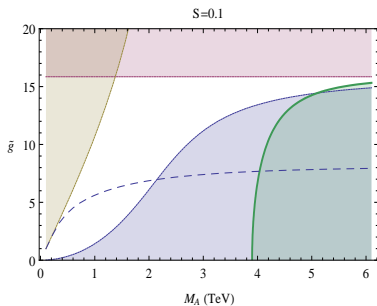
# UNITARITY WITHOUT A HIGGS

- If the Higgs boson is heavy, the technirho meson delays unitarity violation to higher energy scales.
- The unitarization drops after a threshold value of the  $\tilde{g}$  coupling is reached.



# UNITARITY WITHOUT A HIGGS

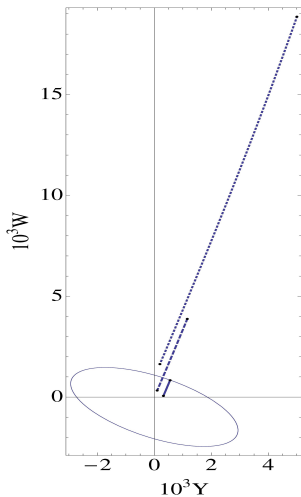
- The technirho mass needs not be within the LHC range, because of the strong coupling.
- Upper and lower bounds on  $\tilde{g}$  from the requirements  $\Gamma_{V \rightarrow \pi\pi}/M_V < 0.5$ ,  $\Gamma_{V \rightarrow AA}/M_V < 0.5$ . Upper (consistency) bound from  $\tilde{g} < \sqrt{8\pi/S}$ .



# CONSTRAINTS FROM ELECTROWEAK PARAMETERS

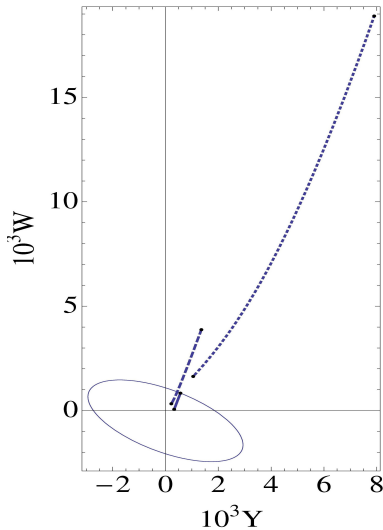
- (M)WT are universal theories, in the sense of Barbieri *et. al.*. (Barbieri, Pomarol, Rattazzi, and Strumia 04)
- $\hat{T} = \hat{U} = V = 0$ , because of custodial symmetry.
- $\hat{S}$  or  $S$  are used as input, since these are small in (M)WT.
- $W$  and  $Y$  impose constraints on the Lagrangian parameters.

# CONSTRAINTS FROM ELECTROWEAK PARAMETERS



$y = 0$ ; solid= $\tilde{g} = 8$ ,  
dashed= $\tilde{g} = 4$ ,  
dotted= $\tilde{g} = 2$ ;  
closer to the ellipse  
 $M_A = 600$  GeV,  
farther from the  
ellipse  $M_A = 150$   
GeV.

# CONSTRAINTS FROM ELECTROWEAK PARAMETERS



$y = 1$ ; solid= $\tilde{g} = 8$ ,  
dashed= $\tilde{g} = 4$ ,  
dotted= $\tilde{g} = 2$ ;  
closer to the ellipse  
 $M_A = 600$  GeV,  
farther from the  
ellipse  $M_A = 150$   
GeV.



- (Minimal) Walking Technicolor is an ordinary 4D gauge theory.
- It satisfies the electroweak constraints.
- It does not give rise to FCNC.
- It can be extended to include more models, like split technicolor or partially gauged technicolor.
- It can be written as an effective field theory with vector and axial resonances and a Higgs (NMWT), plus other resonances in models with a large chiral symmetry group (MWT).

- New leptons can always be added in the spectrum (needed in MWT), as well as matter in the adjoint.
- Dark matter candidates were discussed in previous works. (Gudnason, Kouvaris, and Sannino 06; Kainulainen, Tuominen, and Virkajarvi 06)

WORKSHOP ON  
**DYNAMICAL ELECTROWEAK SYMMETRY BREAKING**  
UNCHARTED TERRITORY - MONSTERS BE HERE!

INTERNATIONAL ADVISORY COMMITTEE

ROBERTO CASALBUONI  
SIMON CATTERALL  
SEKHAR CHIVUNDA  
LUIGI DEL DEBBIO  
NICK EVANS  
CHRISTOPHE GROJEAN  
JOSEPH SCHECHTER  
ELISABETH SIMMONS

LOCAL ORGANISING COMMITTEE

DENNIS D. DIETRICH  
ROSHAN FOADI  
MADS T. FRANDBEN  
MATTI O. JÄRVINEN  
FRANCESCO SANNINO  
MARTIN SWENSSON  
ANDREW SWANN



**UNIVERSITY OF SOUTHERN DENMARK**  
ODENSE

**SEPTEMBER 9 - 13, 2008**

<http://hep.sdu.dk/dewsb>

