

Walking technicolor and AdS/CFT

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Outline

- ❑ Template of walking technicolor
- ❑ AdS/CFT modeling

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Walking technicolor

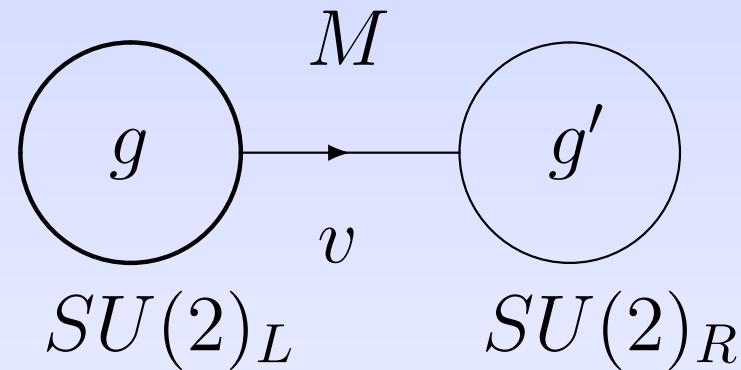
The Higgs sector of the standard model is replaced by a strongly interacting gauge theory

$$U(1)_Y \times SU(2)_L \times SU(3)_C \times SU(N_{TC})_{TC}$$

- ❑ Composite Higgs \Rightarrow no hierarchy problem
- ❑ Walking dynamics \Rightarrow flavor changing neutral currents suppressed
- ❑ Fermions in higher dimensional representations \Rightarrow passes electroweak precision tests

Template for walking technicolor

- ❑ Effective field theory with $SU(2)_L \times SU(2)_R$ chiral symmetry
- ❑ Included composites are Higgs, rho and axial vectors, and (eaten) pions



$M = \frac{1}{\sqrt{2}} (\sigma + i\Pi^a T^a)$ 2×2 linear sigma field, contains Higgs
and the pions

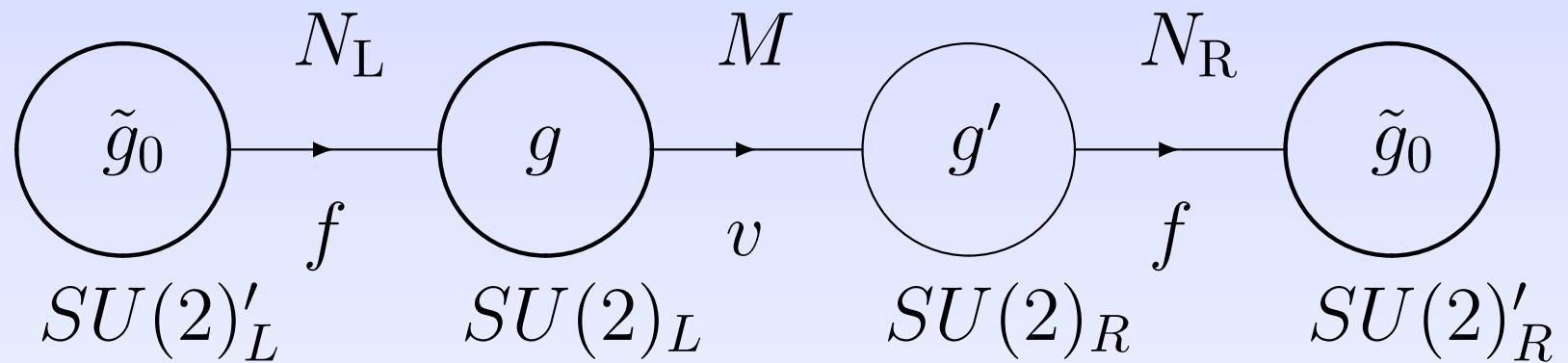
$$M \longrightarrow u_L M u_R^\dagger$$

Coupling to spin-one mesons . . .

Template for walking technicolor

- ❑ Rho and axial vector mesons introduced by **hidden local gauge symmetry**

⇒ four-site moose



$$N_L \rightarrow u'_L N_L u_L^\dagger; \quad N_R \rightarrow u_R N_R u_R'^\dagger \text{ nonlinear sigma}$$

→ write down a general (parity & gauge invariant) effective field theory

Lagrangian

$$M = \frac{1}{\sqrt{2}} (\sigma + i\Pi^a T^a) \leftrightarrow v; \quad N_{L/R} = \exp \left(i\Pi_{L/R}^a T^a \right) \leftrightarrow f$$

$$\mathcal{L} = \mathcal{L}_{\text{kin}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{MN}} + \mathcal{L}_\gamma + \mathcal{L}_{\text{ferm}}$$

$$\begin{aligned}\mathcal{L}_{\text{kin}} &= -\frac{1}{4}\text{Tr}\left[\widetilde{W}_{\mu\nu}\widetilde{W}^{\mu\nu}\right] - \frac{1}{4}\text{Tr}[B_{\mu\nu}B^{\mu\nu}] - \frac{1}{4}\text{Tr}[F_{L\mu\nu}F_L^{\mu\nu}] - \frac{1}{4}\text{Tr}[F_{R\mu\nu}F_R^{\mu\nu}] \\ \mathcal{L}_{\text{Higgs}} &= \frac{1}{2}\text{Tr}\left[D_\mu M(D^\mu M)^\dagger\right] + \frac{v^2\lambda}{2}\text{Tr}\left[MM^\dagger\right] + \dots \\ \mathcal{L}_{\text{MN}} &= \frac{r_1}{2}\text{Tr}\left[(D^\mu N_L)^\dagger D_\mu N_L MM^\dagger + D_\mu N_R (D^\mu N_R)^\dagger M^\dagger M\right] + \dots \\ \mathcal{L}_\gamma &= -\frac{2\gamma}{v^2}\text{Tr}\left[N_L^\dagger F_{L\mu\nu} N_L M N_R F_R^{\mu\nu} N_R^\dagger M^\dagger\right]\end{aligned}$$

Bare parameters $g, g', v, \tilde{g}_0, f, \lambda, \gamma, r_1, r_2, r_3, s, \kappa$

Fix e, G_F, M_Z + constraints from gauge theory using **Weinberg sum rules**

\Rightarrow parameters $M_A, \tilde{g}, S, M_H, \gamma, \kappa, r_1 + s$

AdS/CFT modeling

- ❑ A model of walking gauge theories (with $D > 4$) → estimate for non-perturbative quantities, the S parameter
- ❑ Walking dynamics close to conformal \Rightarrow AdS/CFT expected to work well?

Model of Gürsoy, Kiritssis & Nitti of (pure glue) QCD

$$\mathcal{L} \propto \int d^5x \sqrt{-g} [R - \frac{4}{3}(\partial\Phi)^2 + V(\Phi)]$$

$$V(\Phi) \longleftrightarrow \beta_{\text{QCD}}(g)$$

Apply to walking dynamics by changing the beta function?

→ Adding fermions is tricky

Beta function mapping

In more detail (denoting $\lambda = g^2 N$)

- ❑ For $\lambda \rightarrow 0$ expect $\lambda = e^\Phi$ (and that $\log \mu$ equals the warp factor in the metric)
- ❑ $V(\Phi)$ is fixed such that the equations of motion give the renormalization equation $\frac{d\lambda}{d \log \mu} = \beta(\lambda)$

However, in general ($\lambda > 0$)

$$\lambda = f_\lambda(e^\Phi)$$

The transition function f_λ affects relation of $V(\Phi)$ and β

Beta function mapping

In general:

$$\beta_{\text{QCD}} \xleftrightarrow{f_\lambda} V(\Phi)$$

Our work in progress:

- ❑ “Exact” **all orders** QCD beta function (Ryttov, Sannino) fixes lhs

$$\beta(g) = -\frac{g^3}{(4\pi)^2} \frac{\beta_0 - \frac{2}{3} T(r) N_f \gamma(g^2)}{1 - \frac{g^2}{8\pi^2} C_2(G) \left(1 + \frac{2\beta'_0}{\beta_0}\right)}$$

- ❑ QCD Regge behavior constrains rhs

⇒ combine to reduce the uncertainty in the transition functions

Summary

- ❑ Walking technicolor passes all experimental constraints
- ❑ We study a $SU(2)_L \times SU(2)_R$ version of walking technicolor using effective field theory
- ❑ Exact QCD beta function can be used to reduce uncertainties in AdS/CFT modeling